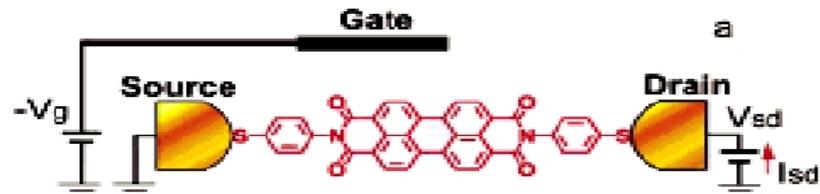


中国科学技术大学

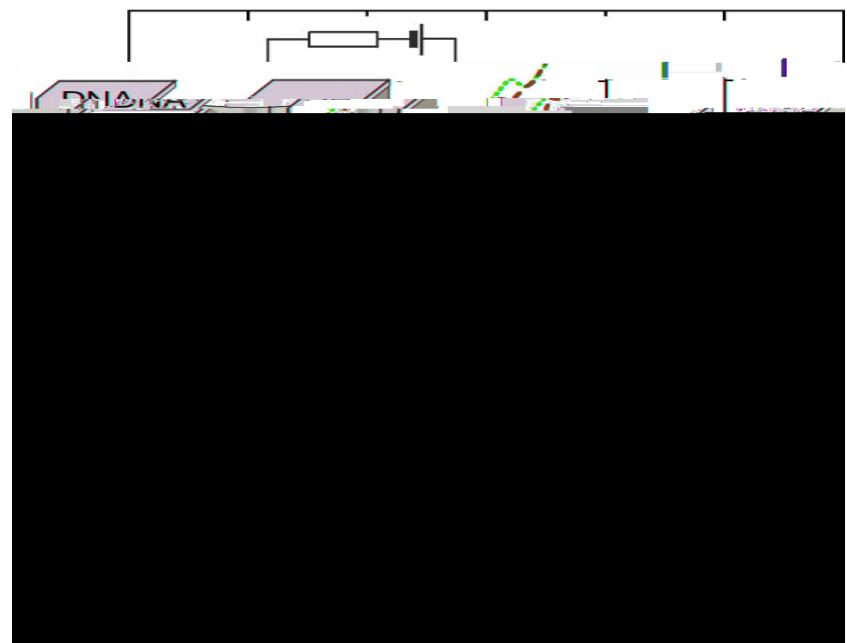




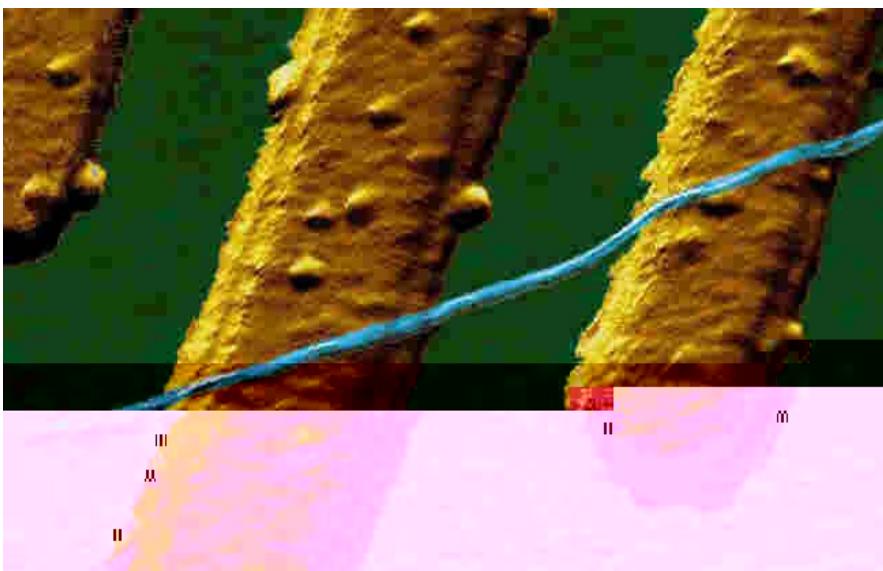
Mark A. Reed and James M. Tour, Scientific American (June '00)



FET



SWCNT across Pt Electrodes (C. Dekker, Delft)



## **Size: scale from 1 to 100 nm**

Given the same amount of space, many times faster and more powerful than silicon based devices

## **Manufacturing:**

fabricated identically and defect-free in enormous numbers

Self-assembly Manufactured  
molecular synthesis tools highly developed

## **Quantized electron energy Extended systems**

Provides thermodynamically favorable electron conduction

Conjugation acts conduction switch

## **Elegance**

Elegant solutions to non-volatile and inherently digital

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*The problems of chemistry and biology can be greatly helped if our ability to see what we are doing, and to do things on an atomic level, is ultimately developed — a development which I think cannot be avoided.*

**Richard Feynman**

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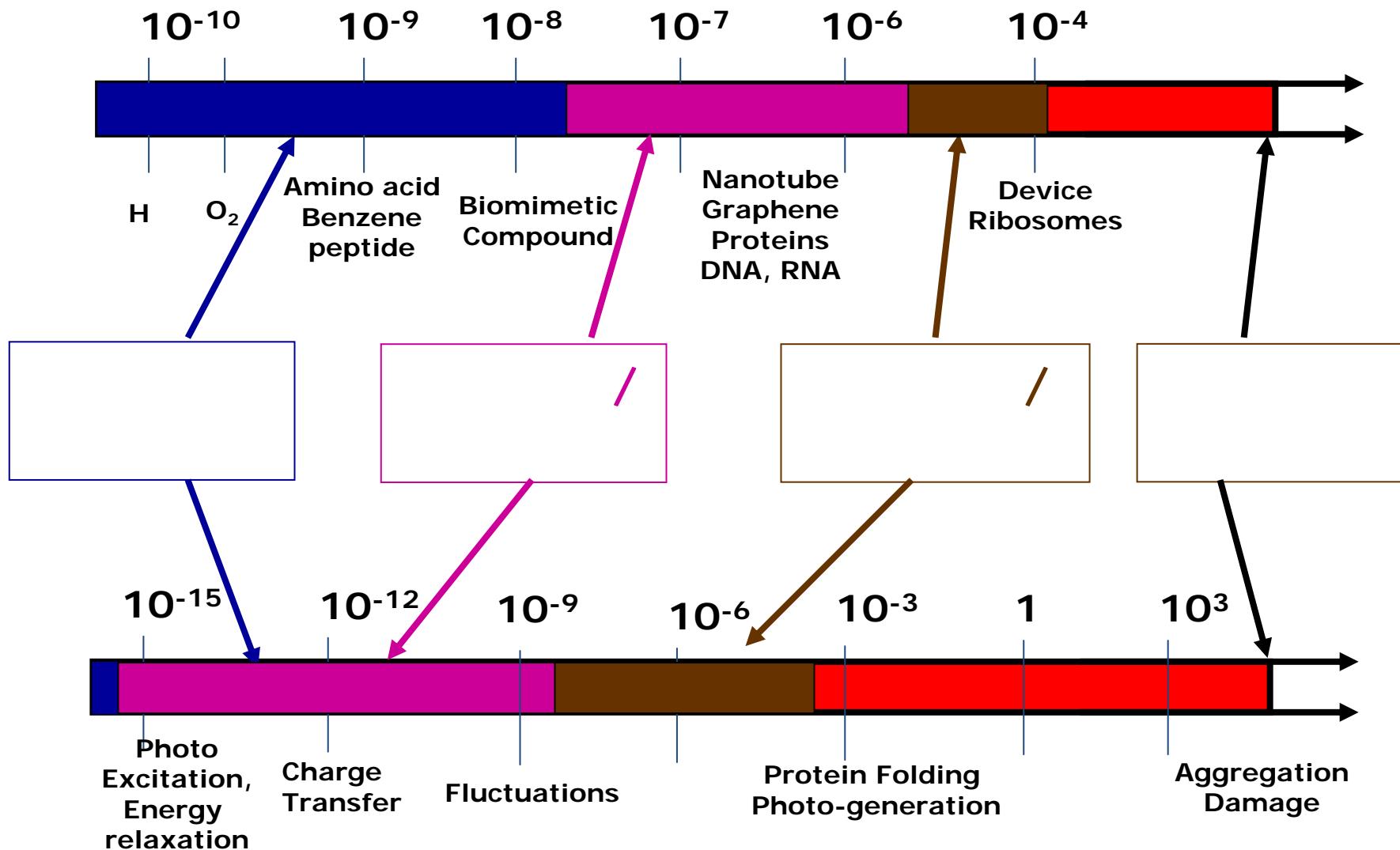
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*The fundamental laws necessary for the mathematical treatment of a large part of physics and the whole of chemistry are thus completely known, and the difficulty lies only in the fact that application of these laws leads to equations that are too complex to be solved.*

**Paul. A. M. Dirac**



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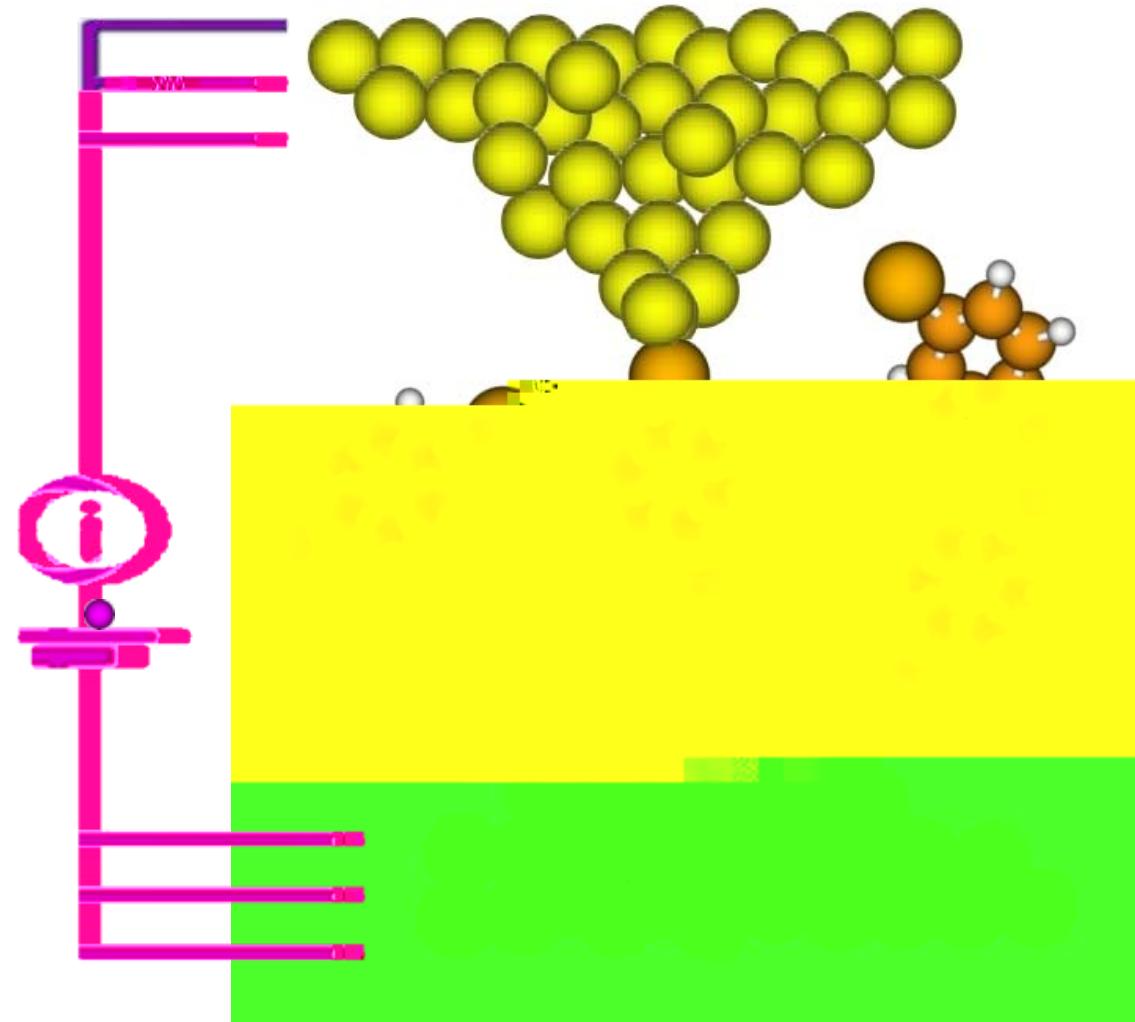
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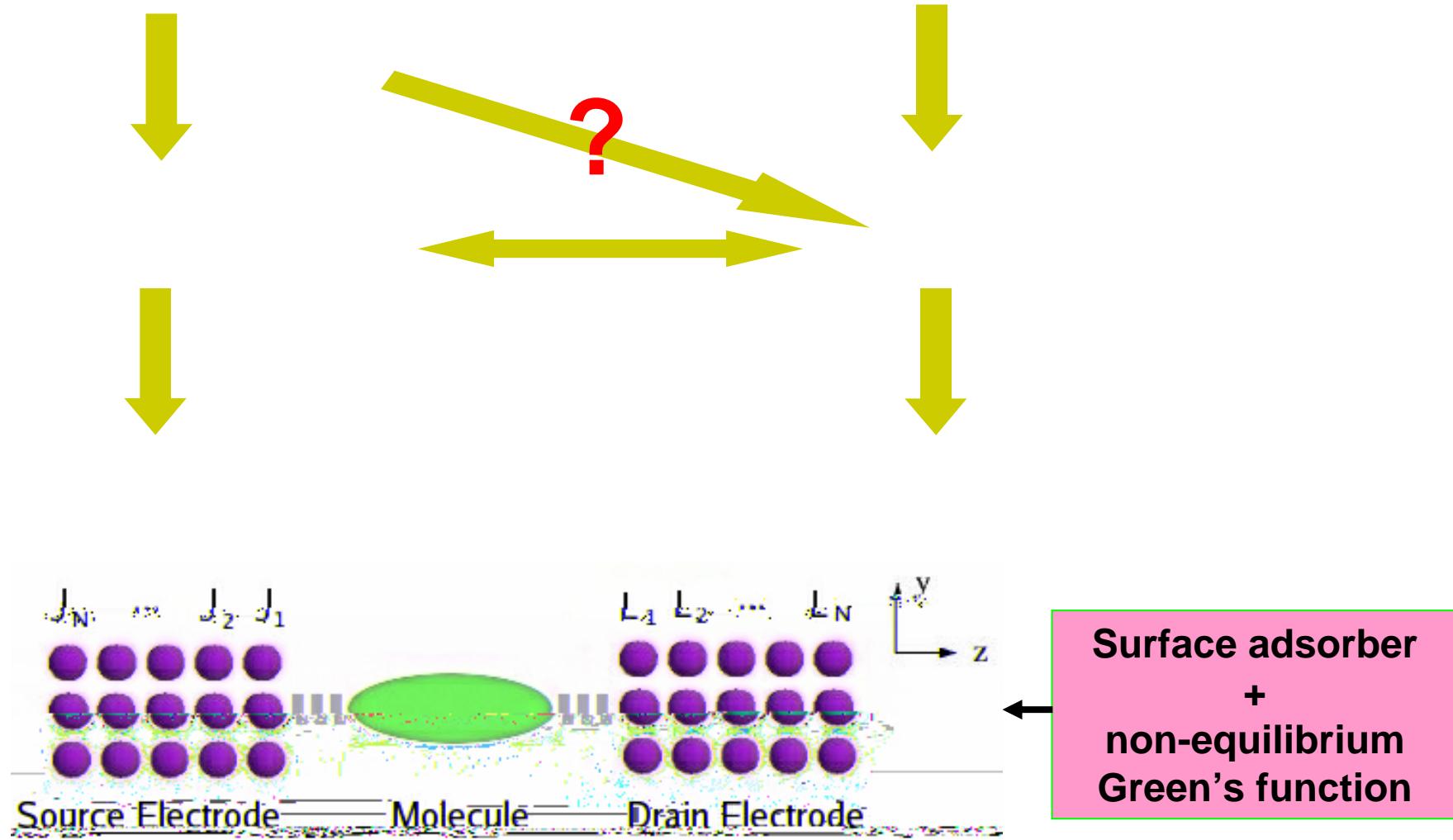
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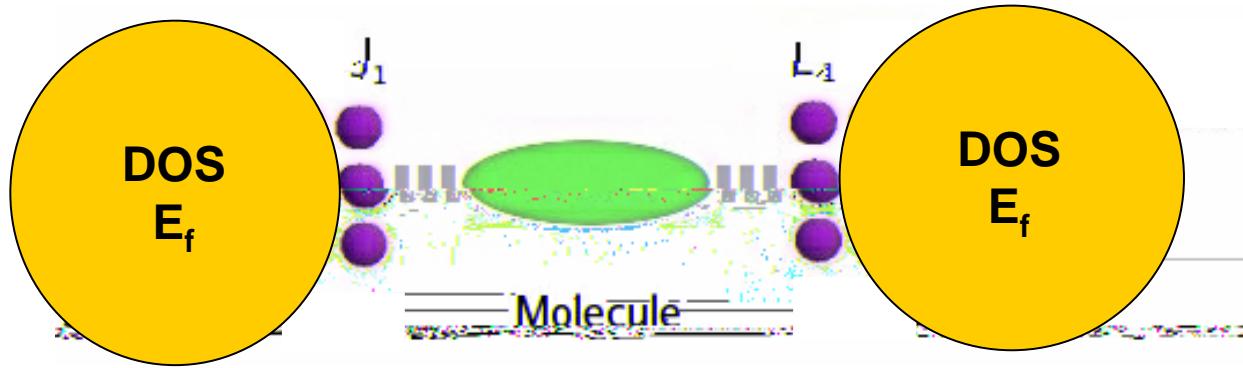
E7A9 Ei abhi a WY a] ghfm Zcf  
ac` YW` af Y` YWfcb] WŁ

u





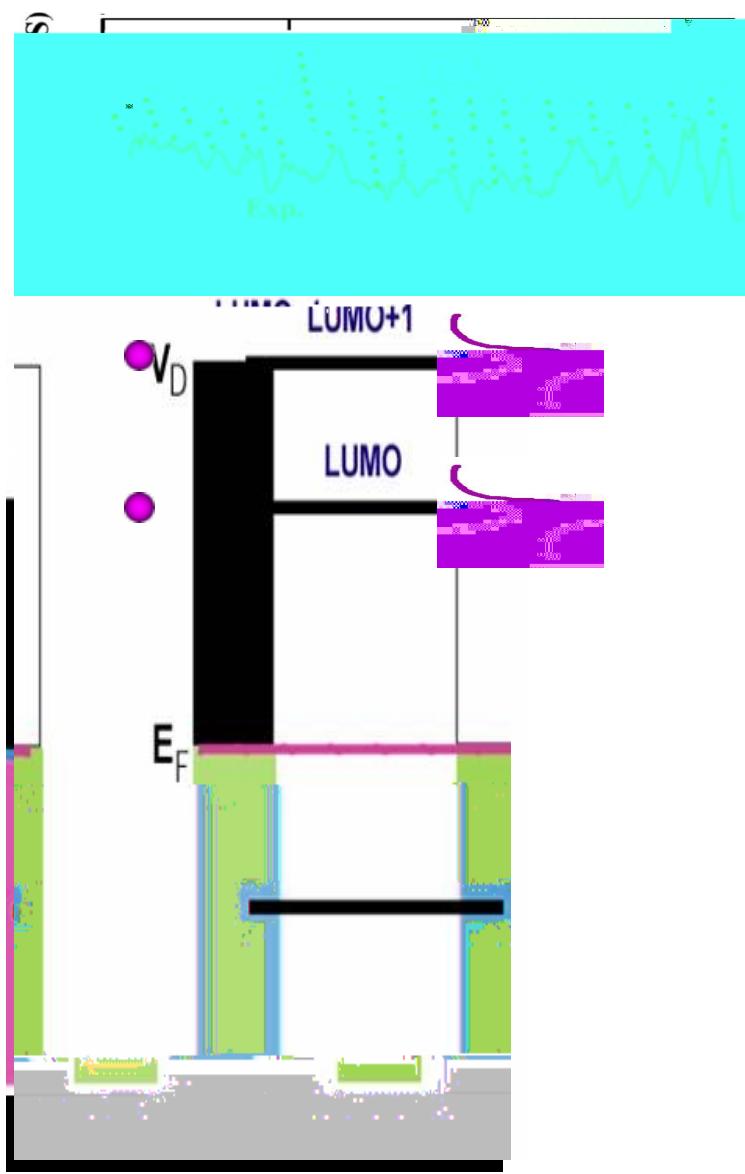
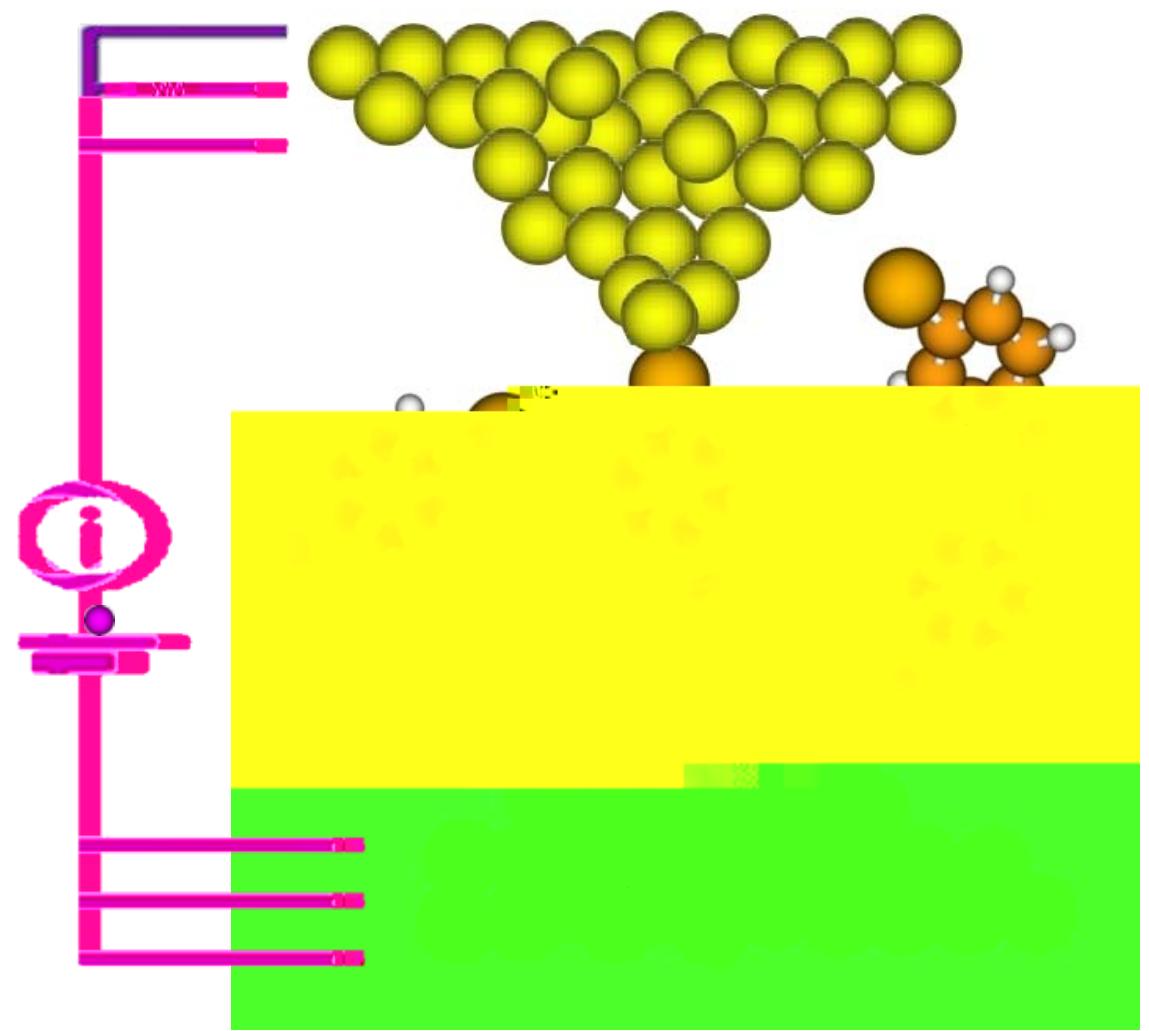
# Landauer Formula



- Treat the “extended molecule” with Quantum Chemistry
- Describe the electrode with Density of States (DOS)
- The extended molecule is in equilibrium with the source and drain by lining up the effective Fermi level

**Tunneling Current Density from Source to Drain:**

$$\frac{2\pi}{\hbar} \sum_{k,q} T \times f(E_k) \times [1 - f(E_q)] \times \delta(E_k - E_q)$$



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f|7=G 7Ybhfa` =bgYfh] cb

GWYyaYŁ

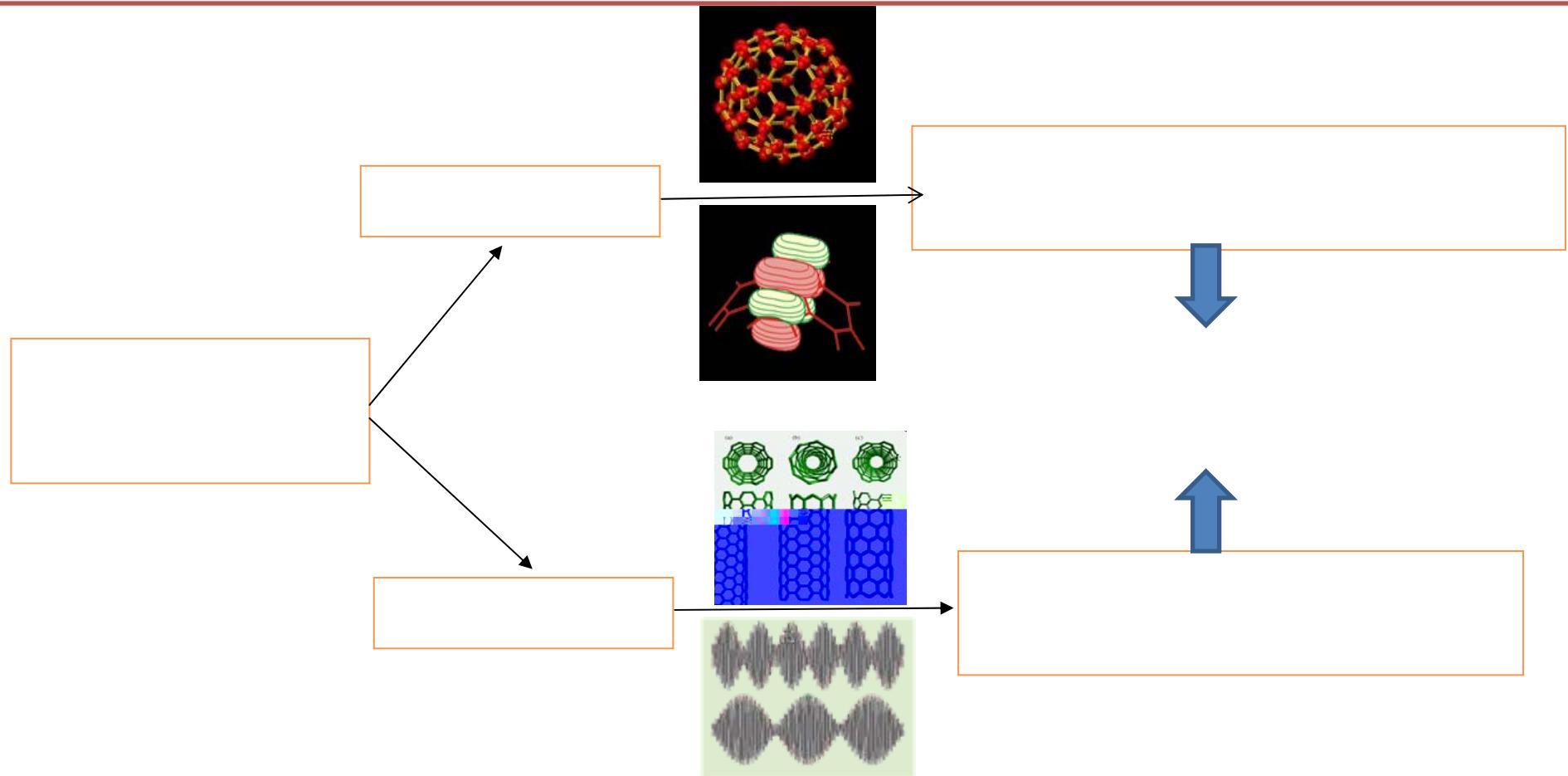
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# **Multi-scale Method ( )**

**A straightforward method can be interfaced with quantum chemistry approach**

**Allows to effectively treat very large nano-scale periodical systems without losing the accuracy**

# Electronic structure was obtained by solving the Schrödinger Equation

$$H |\Psi^\eta\rangle = \varepsilon_\eta |\Psi^\eta\rangle$$

- Energy of electrons can be written as  $\langle \Psi | \hat{H} | \Psi \rangle$
- By the *variational theorem*, this is an upper bound on the true electronic ground state energy
- We minimise the electronic energy as a function of weighting functions of molecular orbitals to obtain *self-consistent* solution (also called SCF approach)
- Use a linear combination of orbital basis functions, or sometimes Gaussian functions:  
$$\Psi_i = \sum_{\mu=1}^N C_{\mu i} \phi_\mu$$
 LCAO
- Substituting these into a 1-electron T.I. Schrödinger eq.:

$$f\Psi_i = \varepsilon_i \Psi_i \rightarrow f \sum_{\mu=1}^N C_{\mu i} \phi_\mu = \varepsilon_i \sum_{\mu=1}^N C_{\mu i} \phi_\mu$$

# Fock Matrix

- The hamitonian of the system can be described as:

$$H = \begin{pmatrix} H_{1,1} & \dots & H_{1,\frac{n}{2}} & H_{1,\frac{n}{2}+1} & \dots & H_{1,n} \\ \dots & \dots & \dots & \dots & \dots & \dots \\ H_{\frac{n}{2},1} & \dots & H_{\frac{n}{2},\frac{n}{2}} & H_{\frac{n}{2},\frac{n}{2}+1} & \dots & H_{\frac{n}{2},n} \\ H_{\frac{n}{2}+1,1} & \dots & H_{\frac{n}{2}+1,\frac{n}{2}} & H_{\frac{n}{2}+1,\frac{n}{2}+1} & \dots & H_{\frac{n}{2}+1,n} \\ \dots & \dots & \dots & \dots & \dots & \dots \\ H_{n,1} & \dots & H_{n,\frac{n}{2}} & H_{n,\frac{n}{2}+1} & \dots & H_{n,n} \end{pmatrix}$$

where  $F$  is the *Fock n* matrix and  $S$  is the *overlap matrix*

- The equation can be written compactly in matrix form :

$$FC = SC\varepsilon, \text{ or,}$$

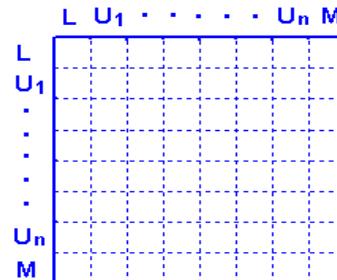
$$(F - SC\varepsilon)C = 0$$

L	U <sub>1</sub>	...	...	U <sub>n</sub>	M
U <sub>1</sub>					
.					
.					
U <sub>n</sub>					
M					

all equations

m)

# The most resource consuming part in Quantum Chemistry calculation is to obtain Fock Matrix through SCF solution



- molecular orbital coefficients  $C_i$  that minimise the energy and give us the best single determinant
- In the secular equation, the Fock matrix itself depends on the coefficients  $C_i$ , and therefore it is necessary to use an iterative procedure when solving the equation
  - We guess values for the  $C_i$  initially and the variational principle (which states that  $E_{\text{guess}}$  is always  $> E_{\text{true}}$ ) allows us to optimise the  $C_i$  until we reach self-consistency
  - However, the Fock matrix contains many two-electron integrals which are very time-consuming to calculate

## **(Central Insertion Scheme)**

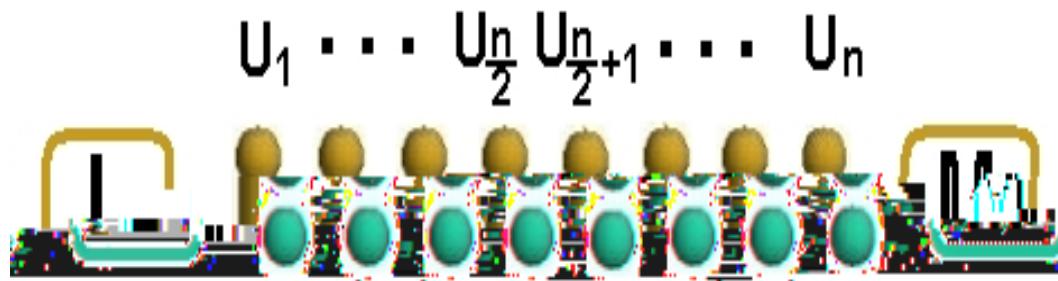
**Deduce out the Fock matrix of large nano-scale systems from the smaller scale system**

**Greatly reduce computation time and resource by avoiding the SCF process**

# (Central Insertion Scheme)

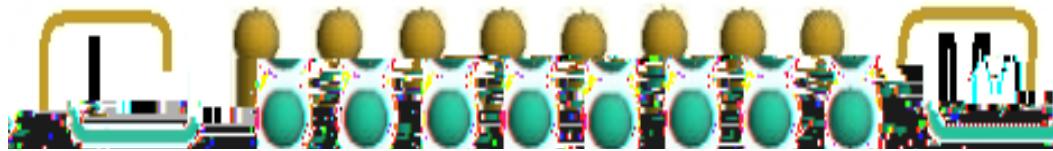
Approximations:

- 1) periodical short-nano-scale systems are long enough to have the electronic structures converged in the central parts
- 2) long distance interaction ( $> 20 \text{ \AA}$ ) between subsystems dies out in the short-nano-scale systems



## 1) Electronic structures converged in the core parts

$$U_1 \cdots U_{\frac{n}{2}} U_{\frac{n}{2}+1} \cdots U_n$$



- The hamiltonian of the system can be described as:

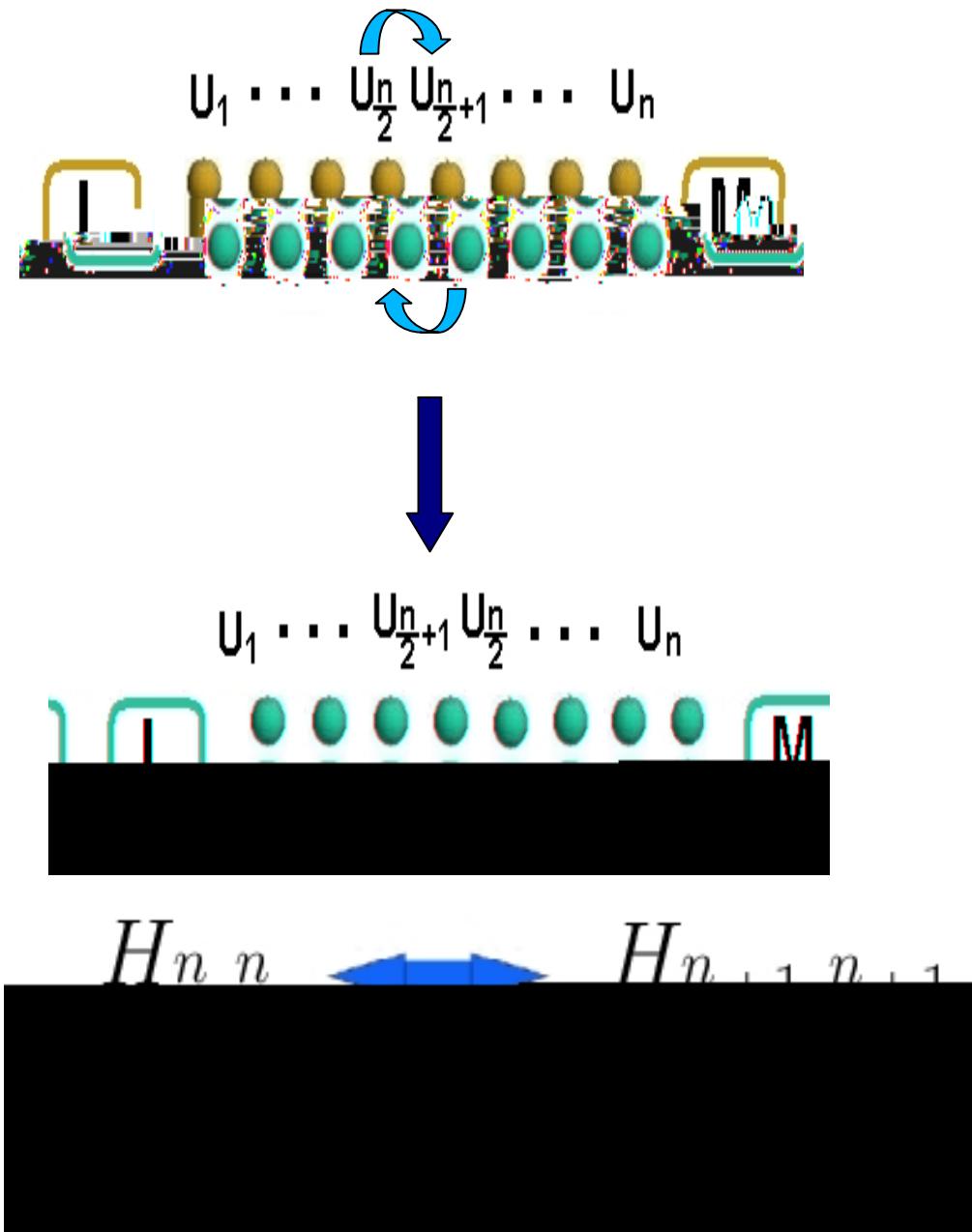
$$H = \begin{pmatrix} H_{1,1} & \dots & H_{1,\frac{n}{2}} & H_{1,\frac{n}{2}+1} & \dots & H_{1,n} \\ \dots & \dots & \dots & \dots & \dots & \dots \\ H_{\frac{n}{2},1} & \dots & H_{\frac{n}{2},\frac{n}{2}} & H_{\frac{n}{2},\frac{n}{2}+1} & \dots & H_{\frac{n}{2},n} \\ H_{\frac{n}{2}+1,1} & \dots & H_{\frac{n}{2}+1,\frac{n}{2}} & H_{\frac{n}{2}+1,\frac{n}{2}+1} & \dots & H_{\frac{n}{2}+1,n} \\ \dots & \dots & \dots & \dots & \dots & \dots \\ H_{n,1} & \dots & H_{n,\frac{n}{2}} & H_{n,\frac{n}{2}+1} & \dots & H_{n,n} \end{pmatrix}$$

- The convergence of interaction

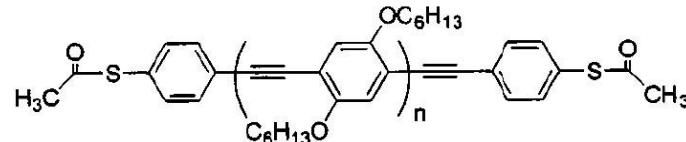
$$H_{\frac{n}{2},\frac{n}{2}} = H_{\frac{n}{2}+1,\frac{n}{2}+1} = \dots$$

$$H_{\frac{n}{2},\frac{n}{2}+1} = H_{\frac{n}{2}+1,\frac{n}{2}+2}$$

# 1) Electronic structures converged in the core parts

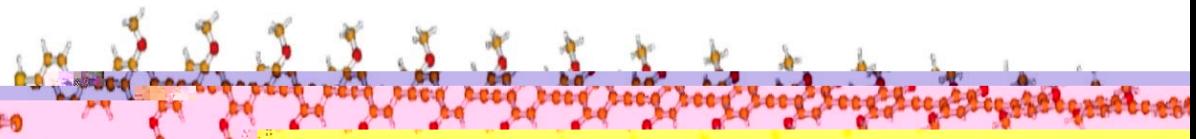


Take polymer ( $n=10$ ) as an example:



$H_n \ n$

$H_{n+1} \ n+1$



Gaussian03 computed HOMO(highest occupied orbital)-LUMO(lowest unoccupied orbital) gap: 2.784eV

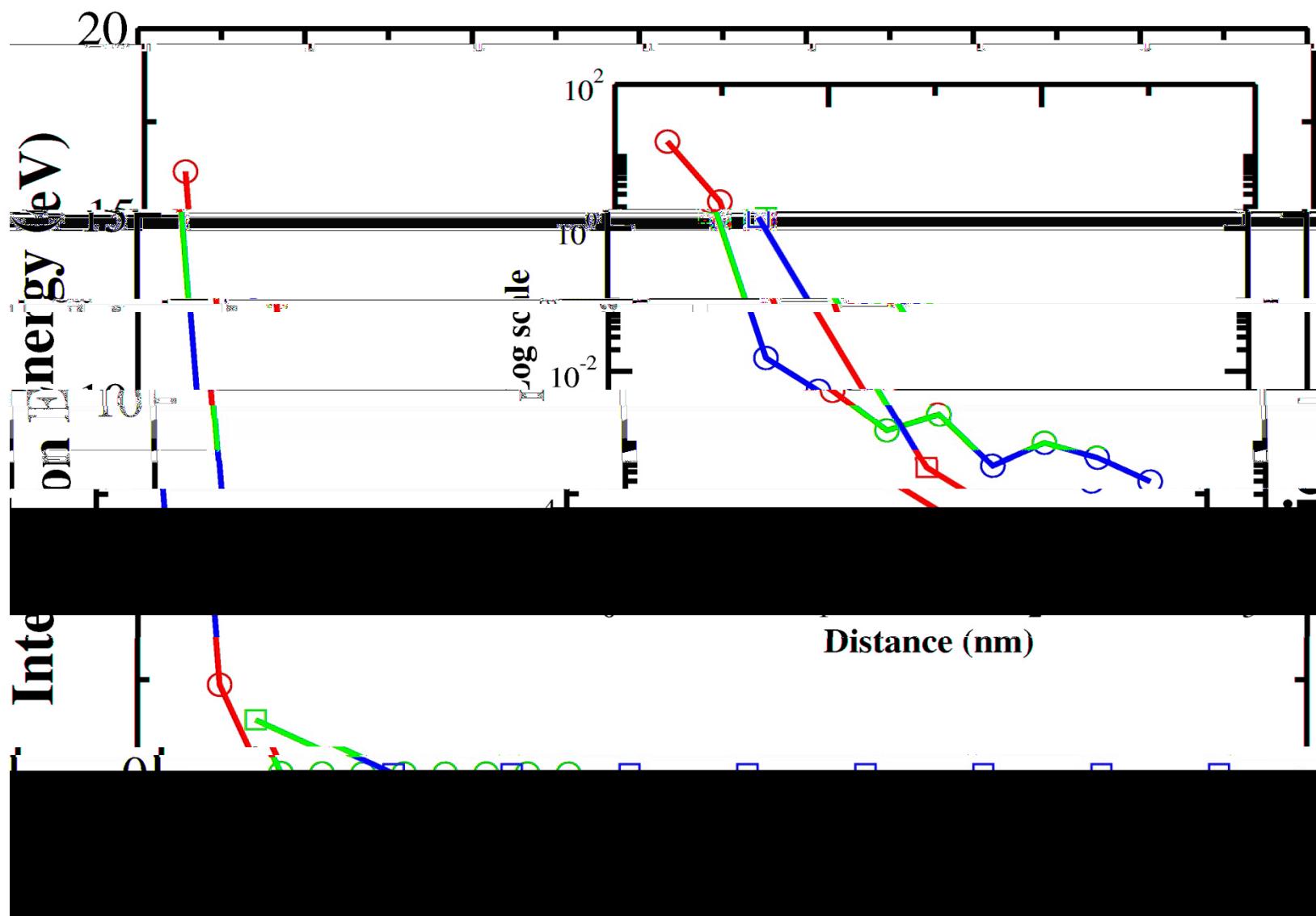
Reconstructing computed HOMO-LUMO gap: 2.786eV, error 2 meV

Orbital energy error in the range of (HOMO-20, LUMO+20):

Biggest error 8 meV      smallest error 0.0005meV

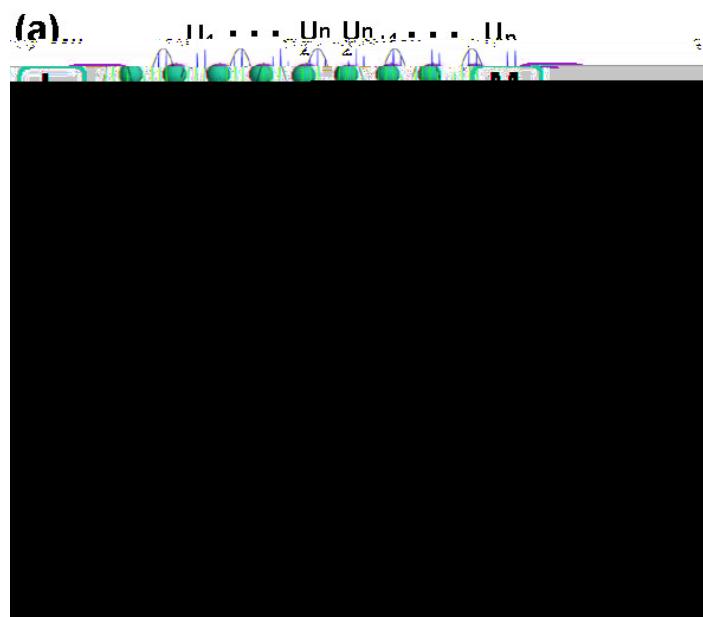
electron density	Reconstructing results	GAUSSIAN03 Results	error
$U_1$	0.0351	0.0357	0.0006
$U_2$	0.0766	0.0771	0.0006
$U_3$	0.1259	0.1251	0.0008
$U_4$	0.1684	0.1657	0.0027
$U_5$	0.1901	0.1847	0.0053
$U_6$	0.1835	0.1864	0.0028
$U_7$	0.1512	0.1550	0.0038
$U_8$	0.1041	0.1086	0.0045
$\equiv U_9 \equiv$	0.0571	0.0603	0.0032
$U_{10}$	0.0227	0.0242	0.0015

## 2) long distance interaction dies out



Interaction energy between U1 and other units as a function of distance between them.

Blue Circle dots: CNT21; Red Square dots: POI Y10

**(a)****(b)**

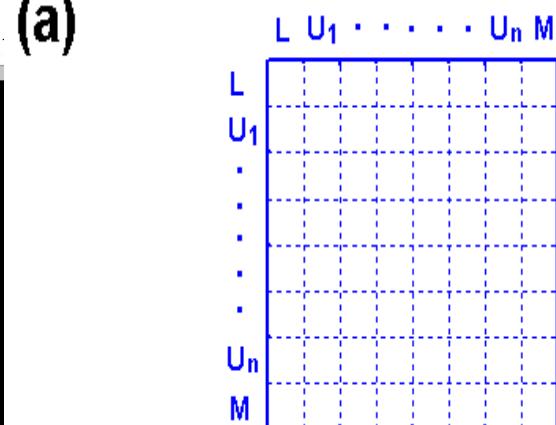
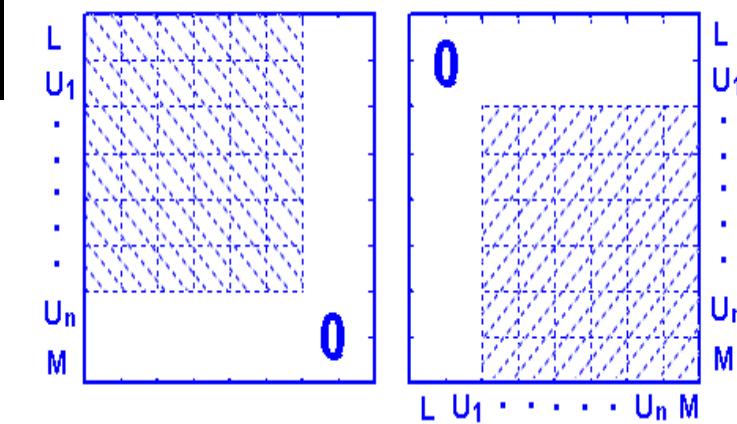
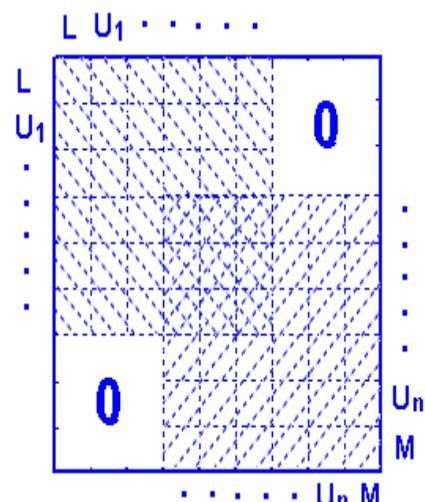
- The hamiltonian of the system can be described as:

$$H = \begin{pmatrix} H_{1,1} & \dots & H_{1,\frac{n}{2}} & H_{1,\frac{n}{2}+1} & \dots & H_{1,n} \\ \dots & \dots & \dots & \dots & \dots & \dots \\ H_{\frac{n}{2},1} & \dots & H_{\frac{n}{2},\frac{n}{2}} & H_{\frac{n}{2},\frac{n}{2}+1} & \dots & H_{\frac{n}{2},n} \\ H_{\frac{n}{2}+1,1} & \dots & H_{\frac{n}{2}+1,\frac{n}{2}} & H_{\frac{n}{2}+1,\frac{n}{2}+1} & \dots & H_{\frac{n}{2}+1,n} \\ \dots & \dots & \dots & \dots & \dots & \dots \\ H_{n,1} & \dots & H_{n,\frac{n}{2}} & H_{n,\frac{n}{2}+1} & \dots & H_{n,n} \end{pmatrix}$$

- The convergence of interaction

$$H_{\frac{n}{2},\frac{n}{2}} = \frac{H_{\frac{n}{2},\frac{n}{2}}}{2} + \frac{H_{\frac{n}{2}+1,\frac{n}{2}+1}}{2} + \dots$$

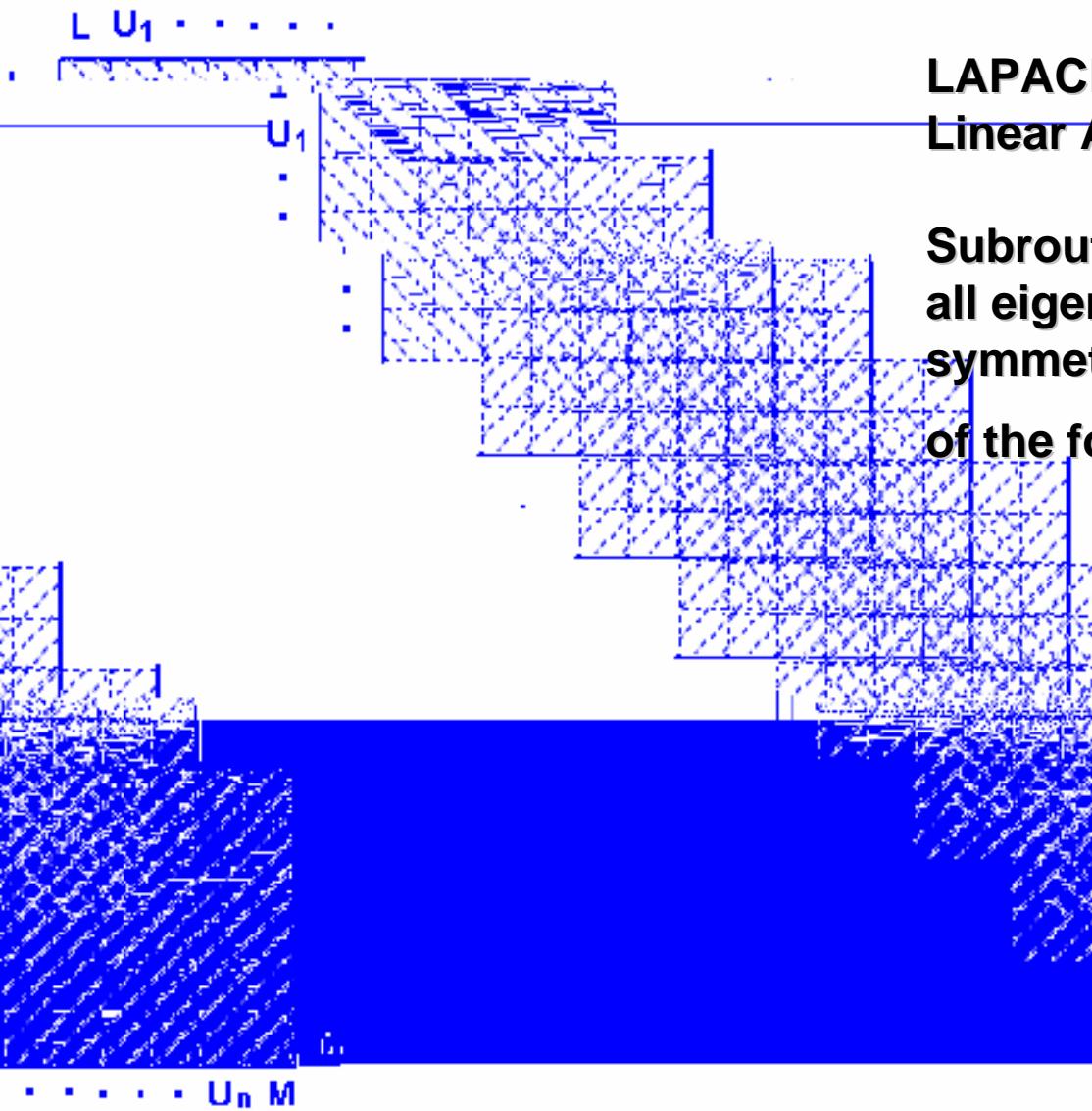
$$H_{\frac{n}{2},\frac{n}{2}+1} = H_{\frac{n}{2}+1,\frac{n}{2}+2}$$

**(b)****(c)**

6] c! Babc@Y[ c

# Eigen value solution

For very large scale

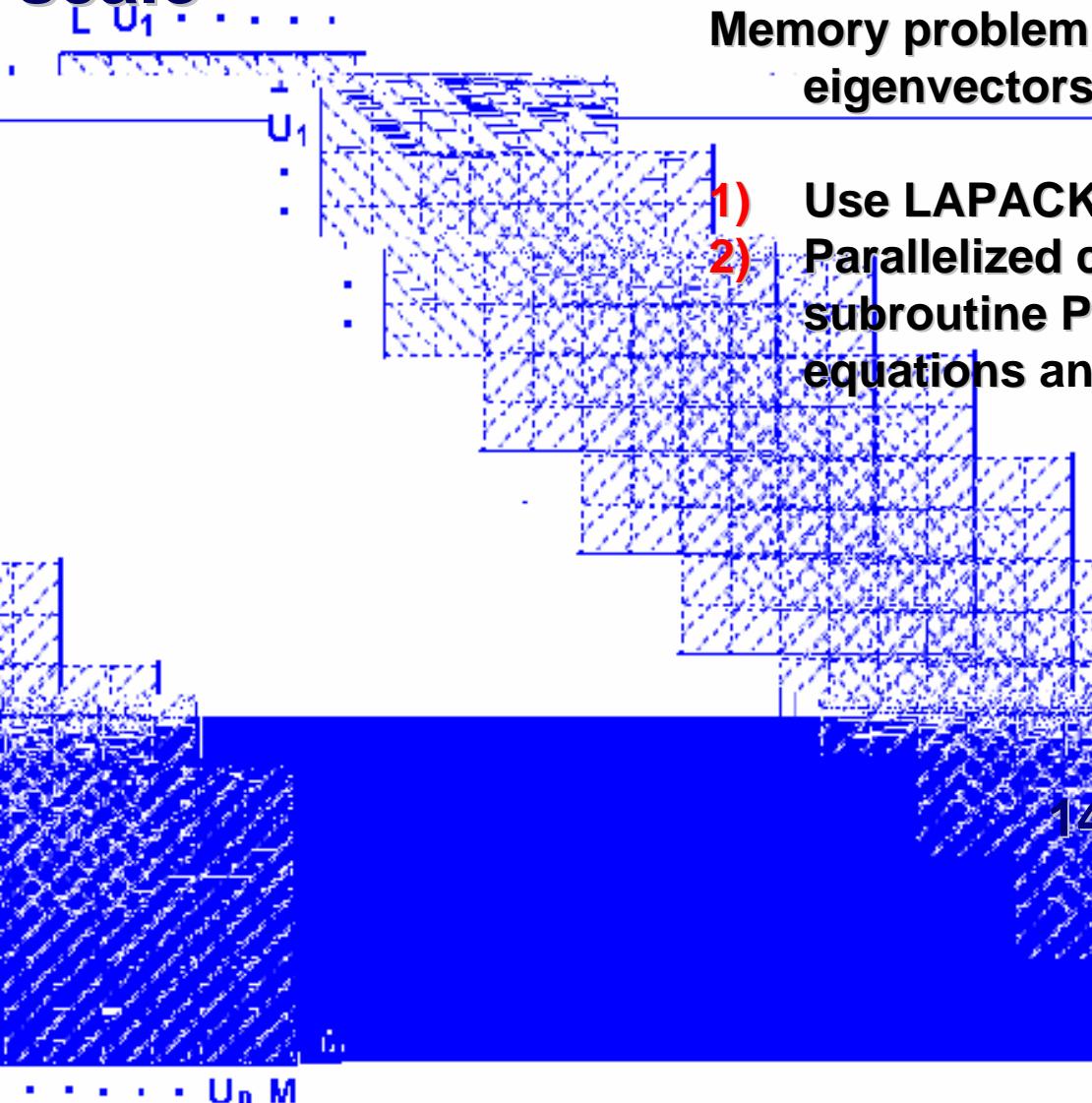


LAPACK –  
Linear Algebra PACKage

**Subroutine ssbgv() or dsbgv()** Compute  
all eigenvalues, and eigenvectors of a  
symmetric-definite banded eigenproblem,  
of the form  $A^*x=(\lambda)B^*x$

# Eigen value solution

For even larger  
scale



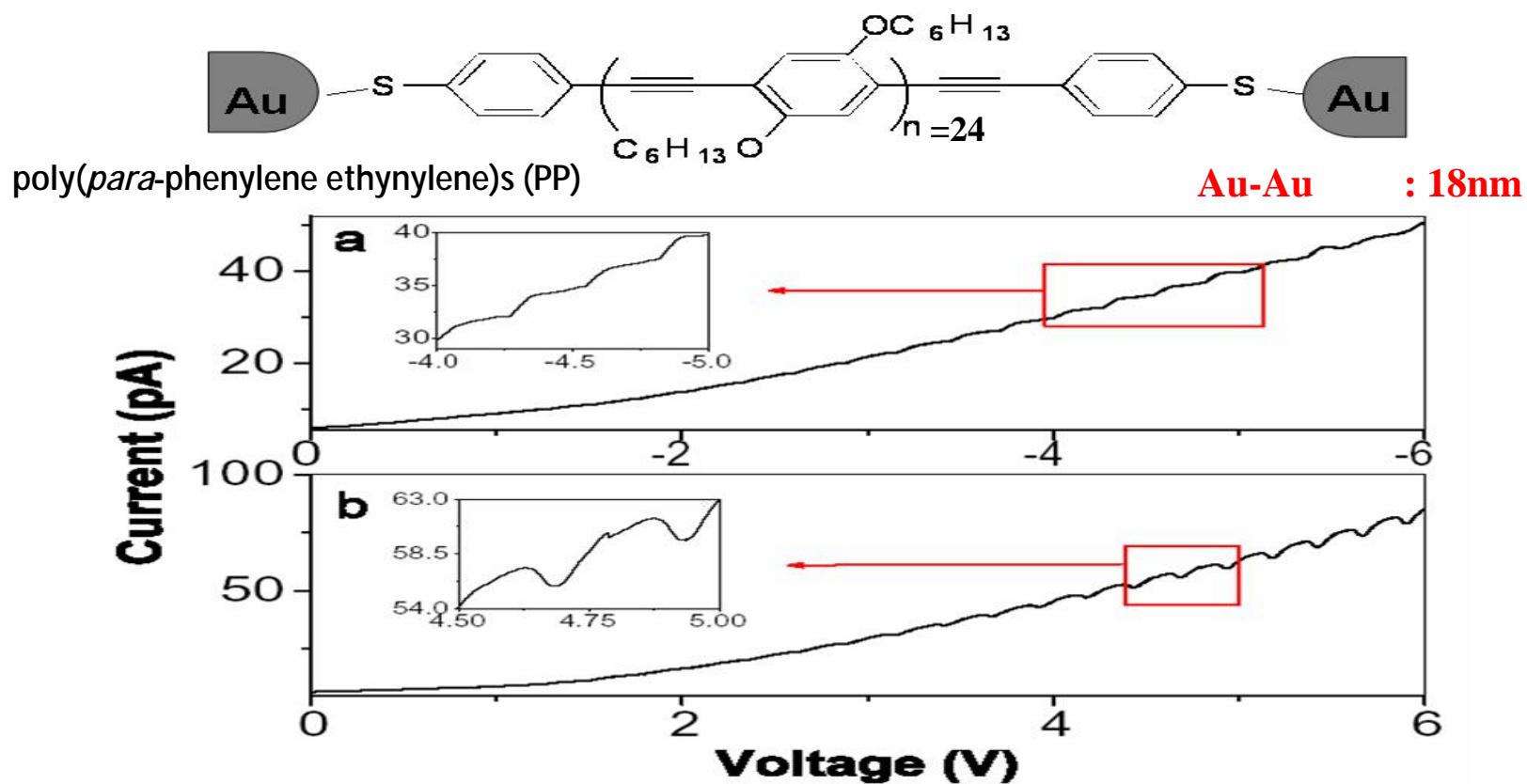
Memory problem, can not find eigenvalues and eigenvectors at the same time

- 1) Use LAPACK Compute all eigenvalues,
- 2) Parallelized calculation with scalapack subroutine PDDBSV to solve the linear equations and find eigenvectors

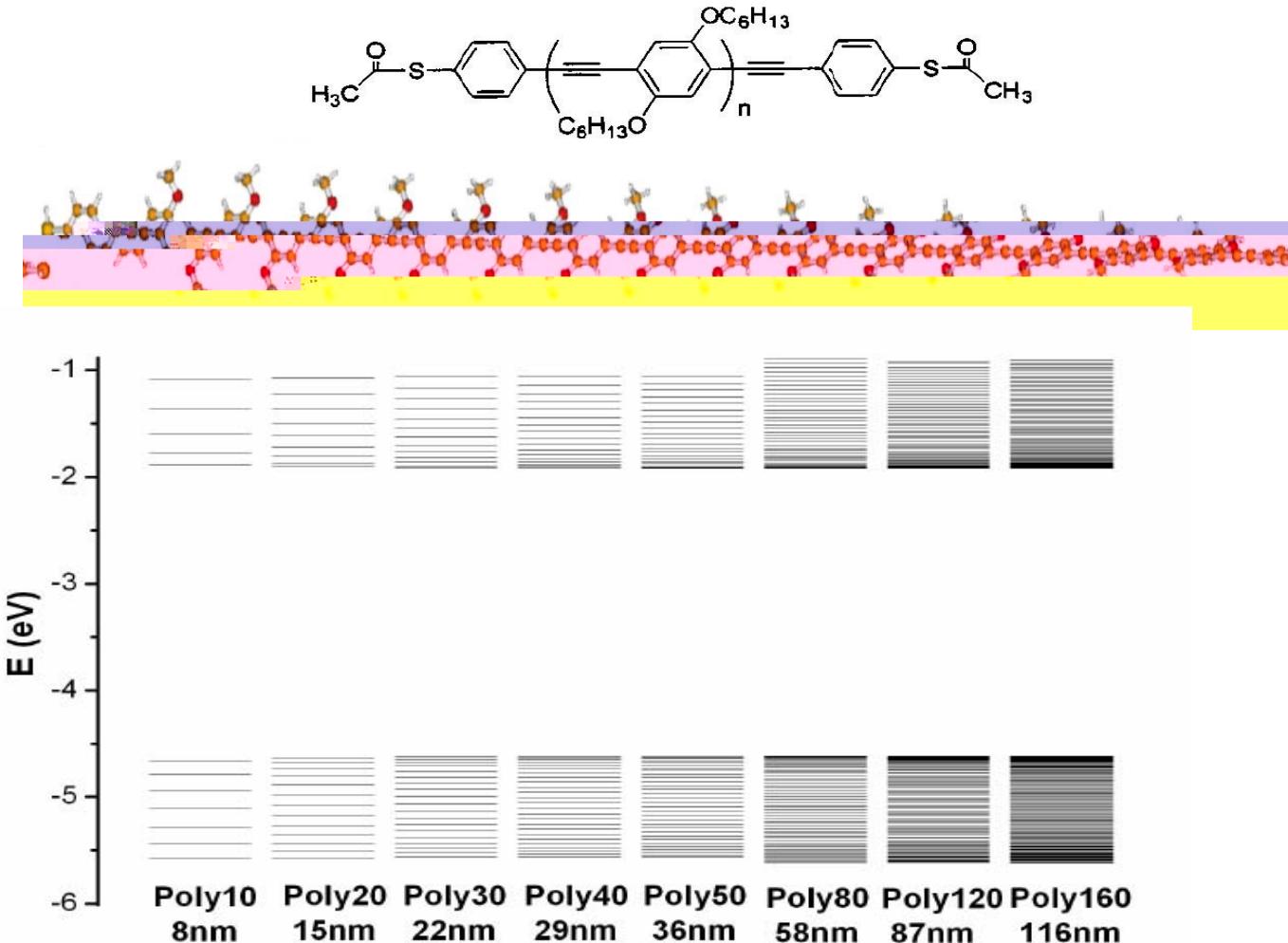
Biggest system:  
200 nm (5 5) carbon nanotube  
97, 020 electrons,  
145, 570 gaussian basis functions  
Maxtrix 145, 570\*145,570

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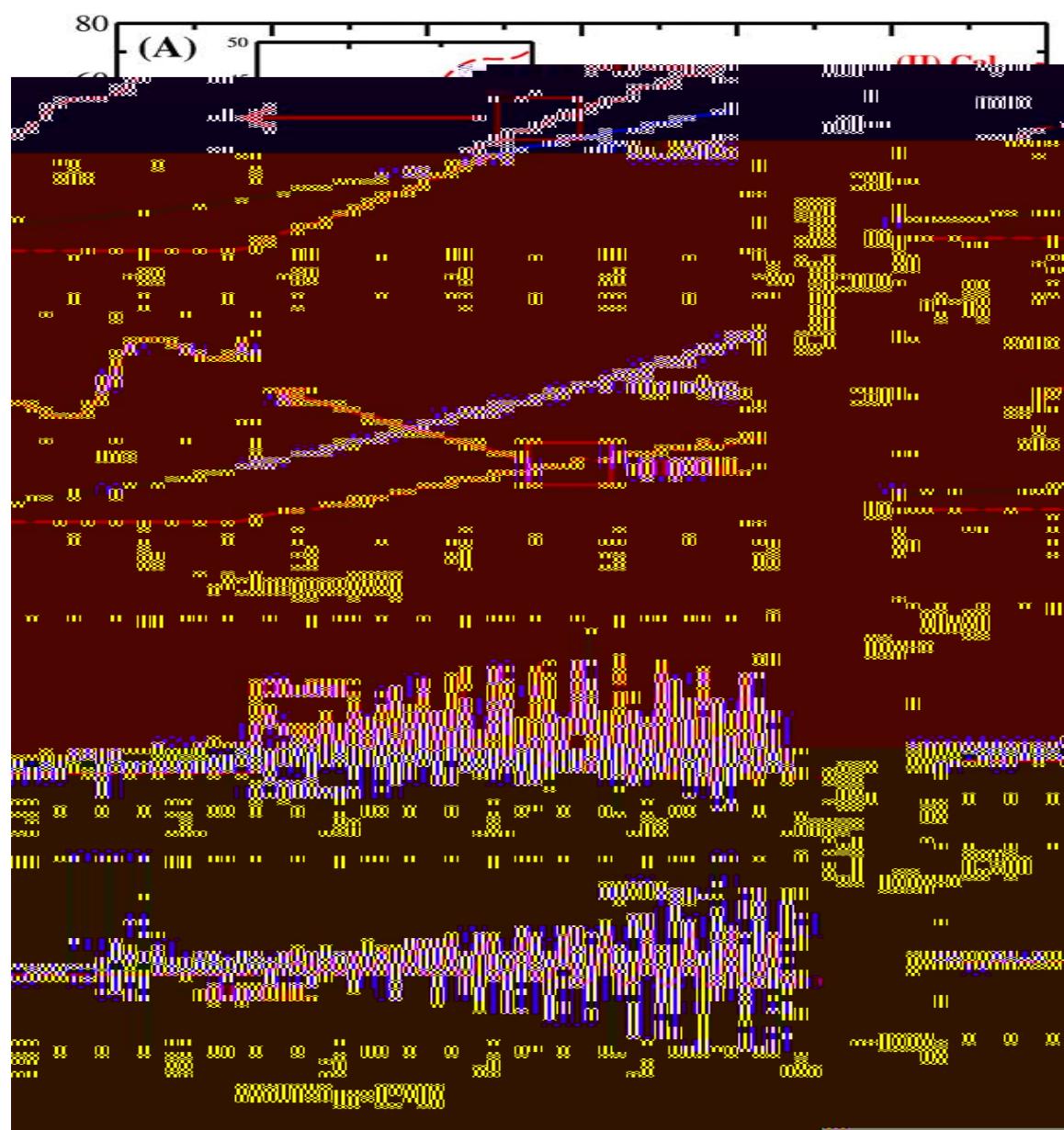
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fID9L



# Conjugated Polymer

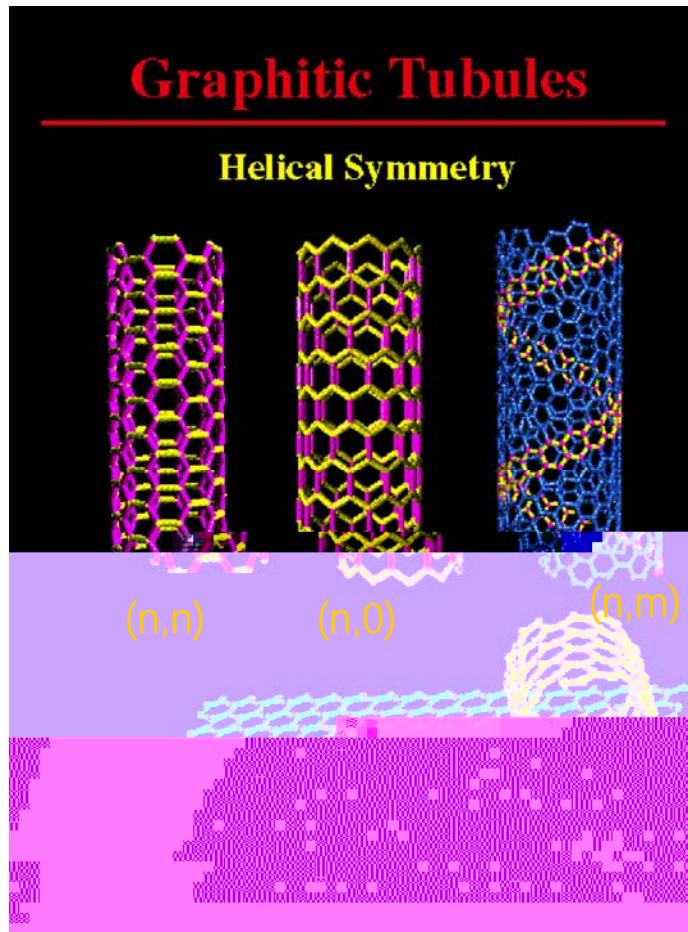


**Increasing of size, Density of State(DOS) increases  
Discrete molecular orbital to continuous band  
HOMO-LUMO gap remains constant as 2.78eV  
Energy structure analogous to inorganic Semiconductor**

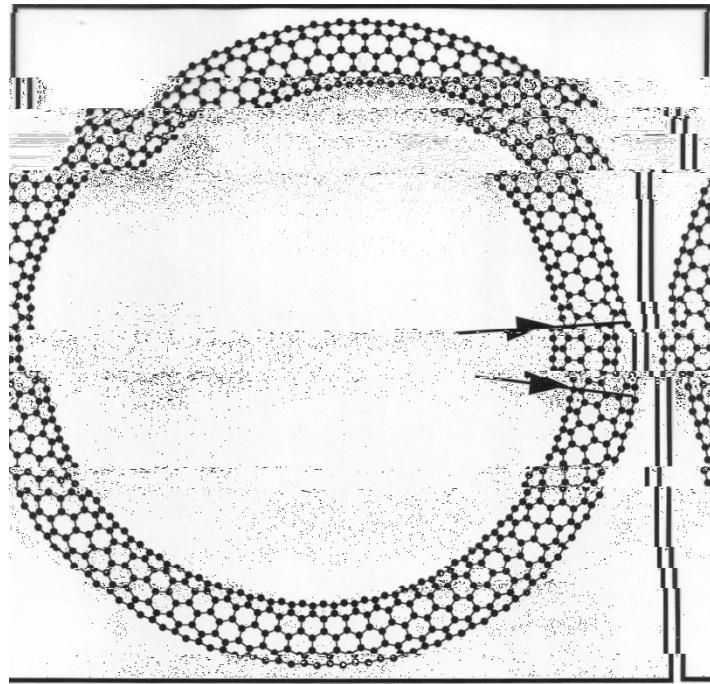


- CIS calculation provide accurate electronic structures for very long conjugated polymers

# Carbon Nanotube-based Structures



A Toroidal Carbon Nanotube

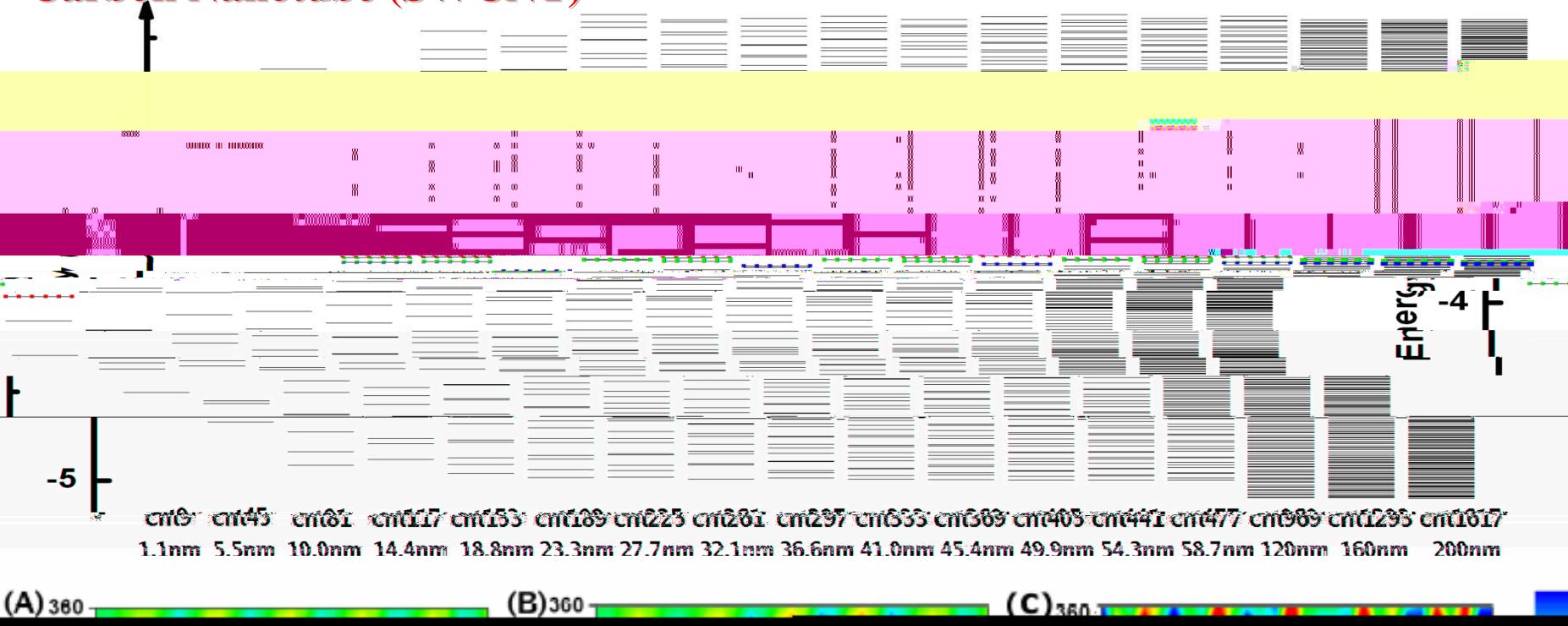


Carbon Nanotubes (CNTs) and toroidal carbon nanotubes (TCNs) are special class of low-dimensional systems which are expected to exhibit novel physical properties due to their special geometries (cylindrical/toroidal) as well as nanoscale sizes

**(5, 5) Metallic Single Walled Carbon Nanotube (SWCNT)**

B3LYP/6-31G

Max: 16170 C atoms described by 150, 000 Gaussian Basis

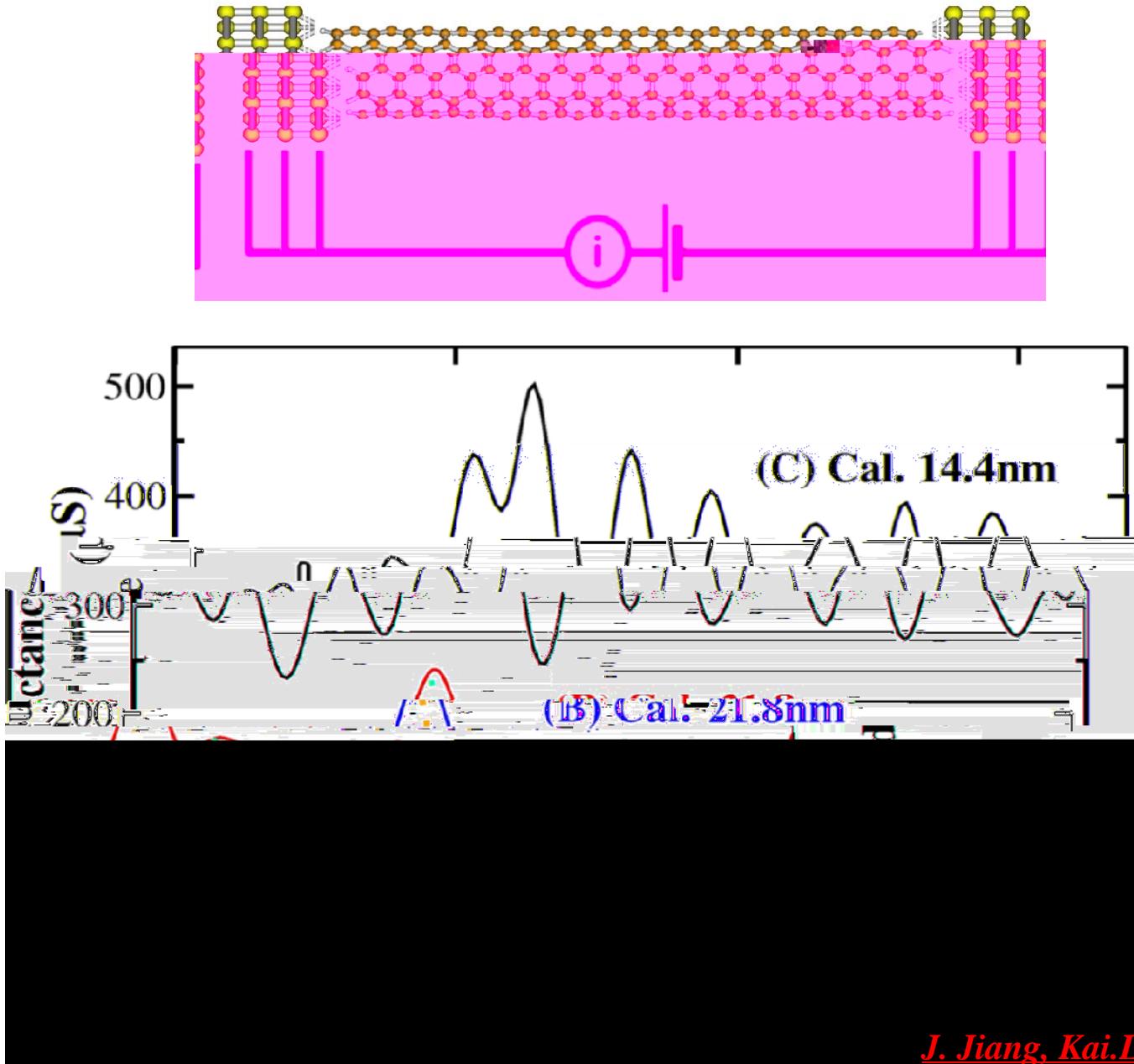


G03 LUMO of 4.1 nm SWCNT

CIS LUMO of 4.1 nm SWCNT

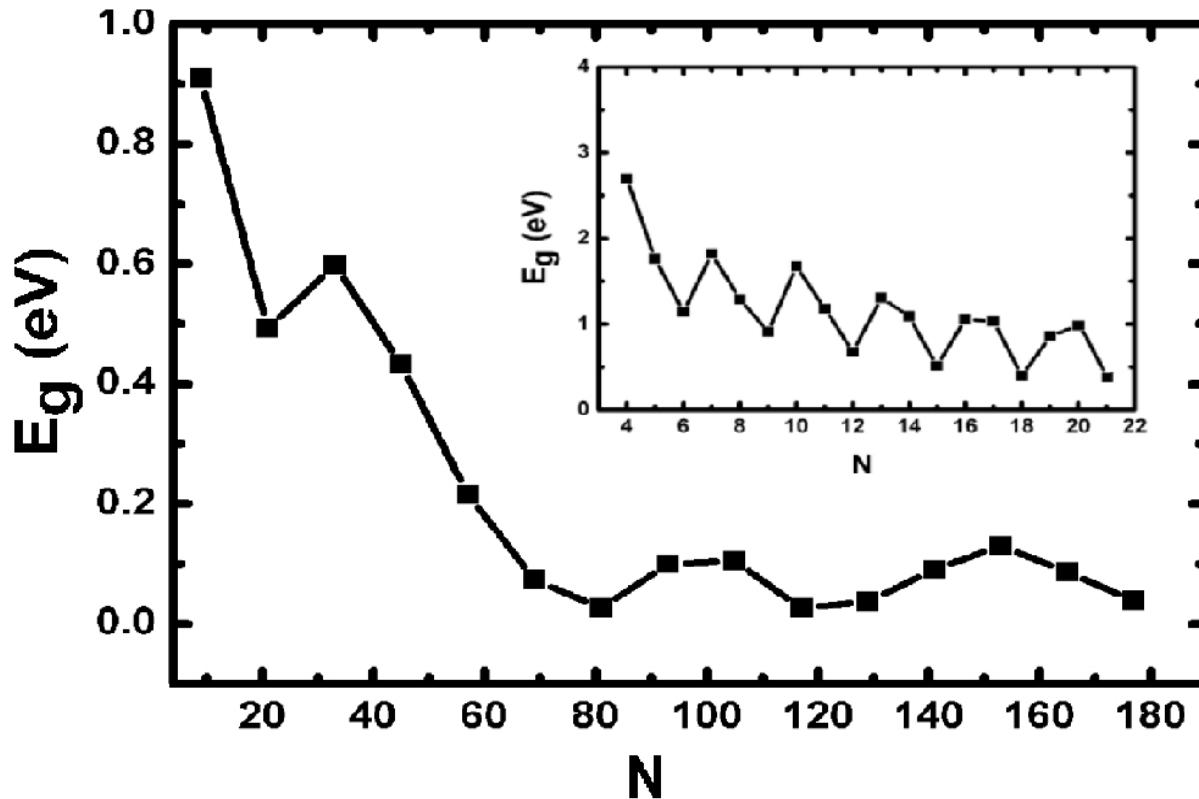
CIS LUMO of 200 nm SWCNT

# Electron transport in (5, 5) CNT



J. Jiang, Kai Liu, Y. Luo,  
J. Chem. Phys. 124, 214711. 2006

## (5, 5) metallic CNT

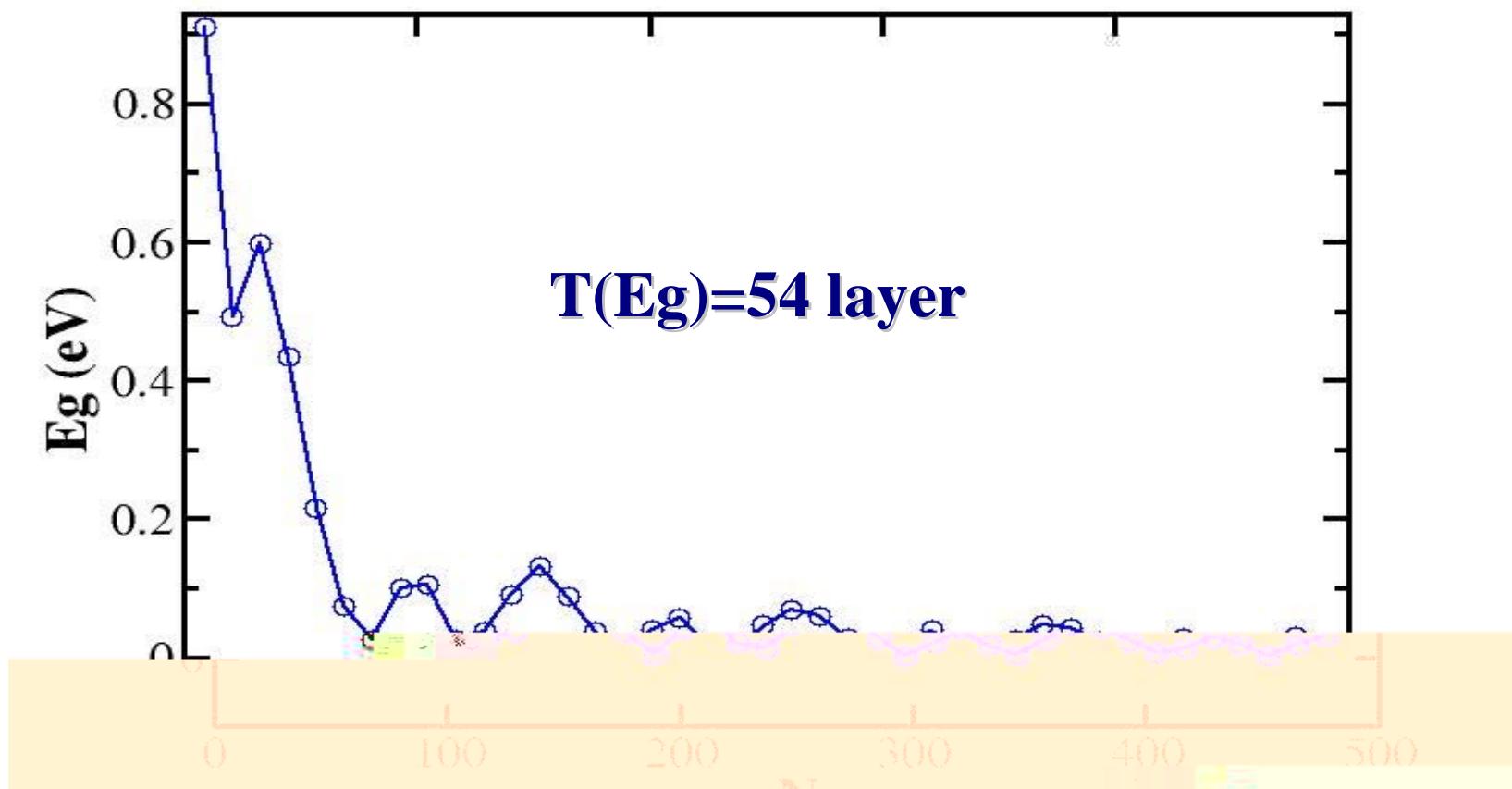
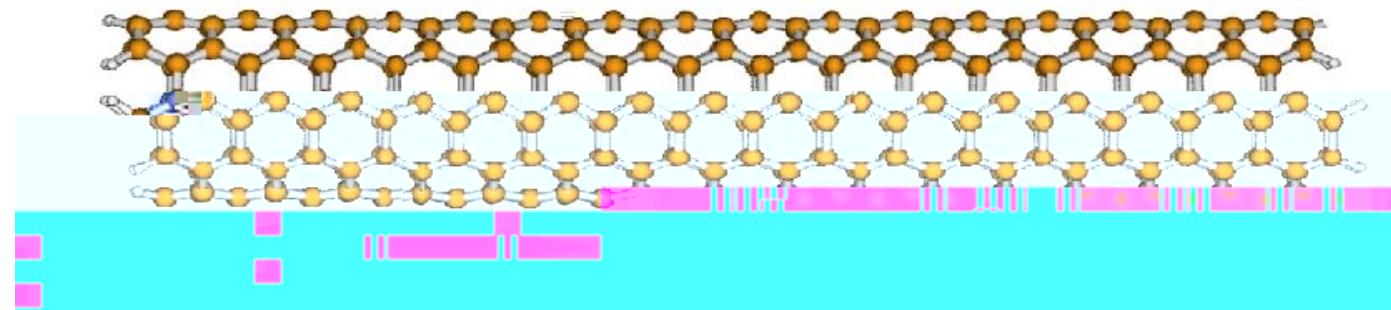


Bigger-period oscillation superimposed on the 3 units period oscillation:

Interference effect between the one dimensional quantum confinement along the tube direction and the two dimensional quantum confinement along the tube ring direction.

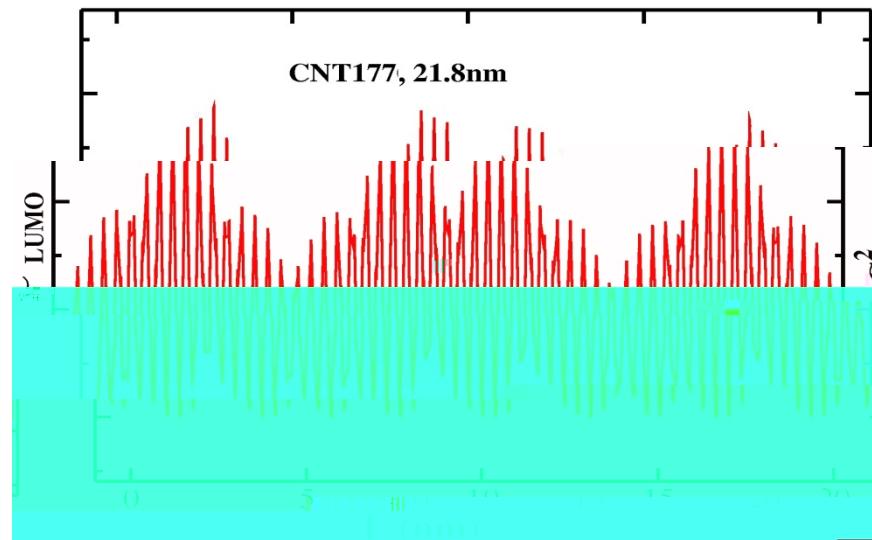
# (5, 5) metallic CNT

DFT  
B3LYP/6-31G

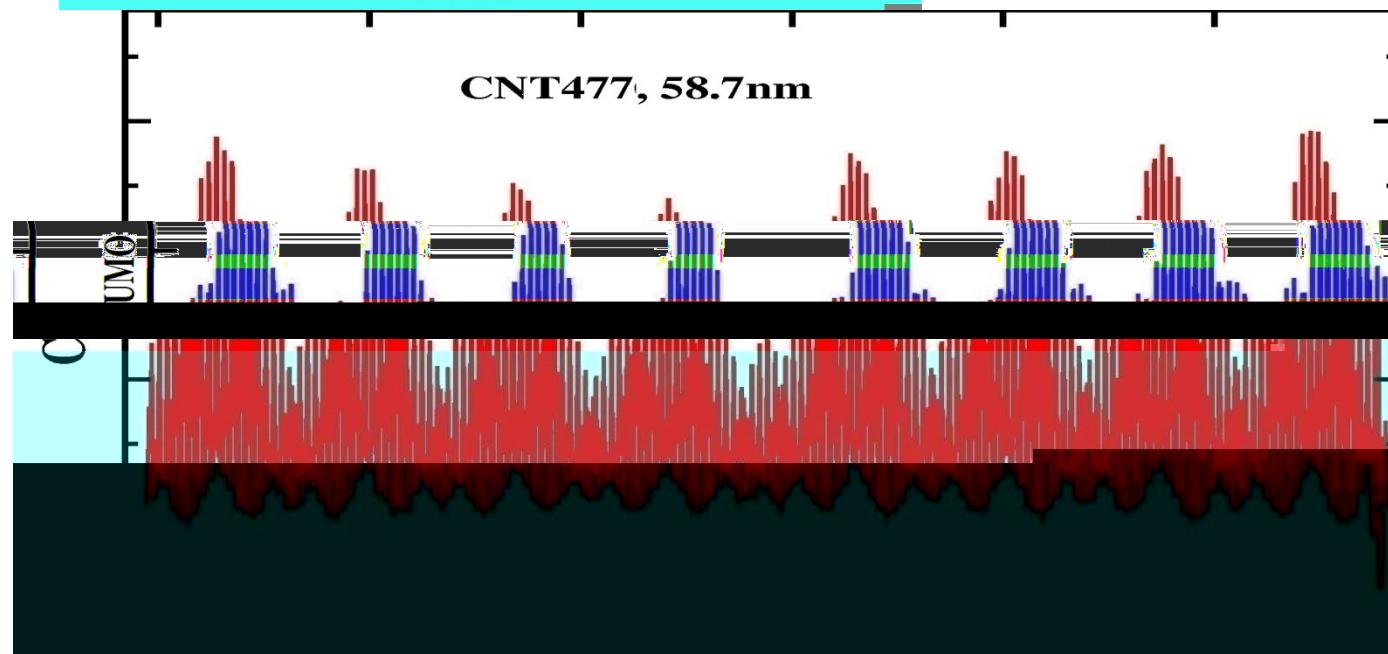
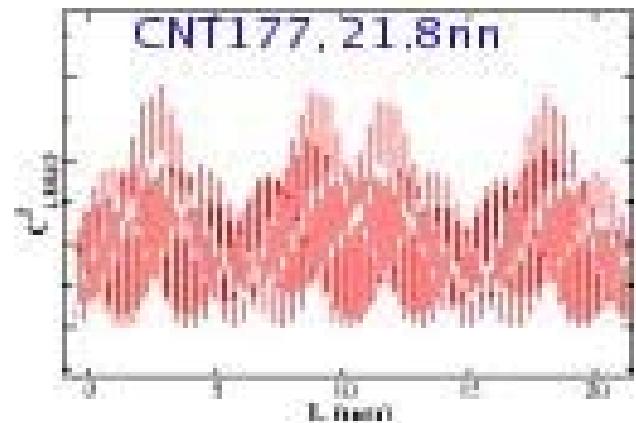


# (5, 5) metallic CNT

## 1D Wavefunction along tube direction



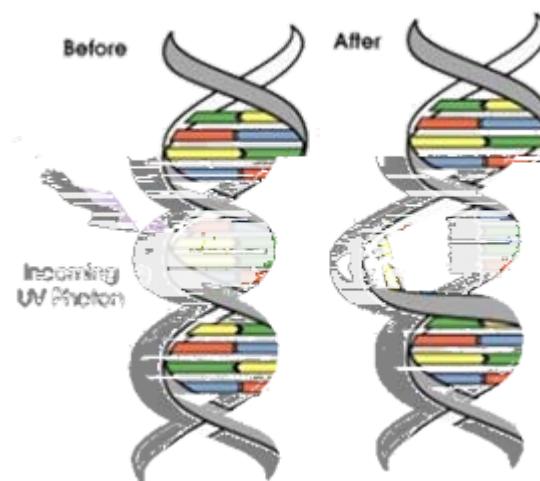
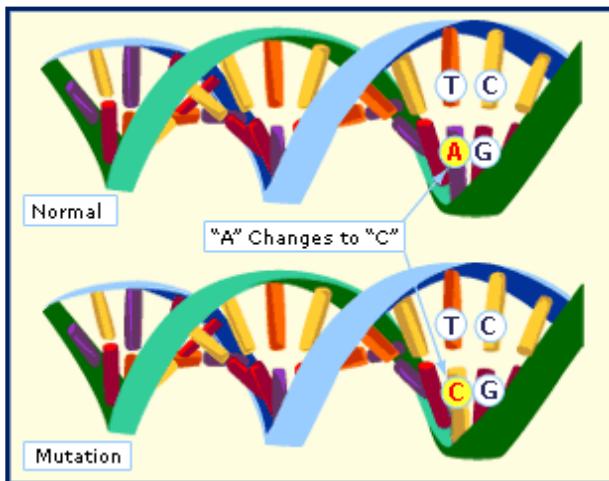
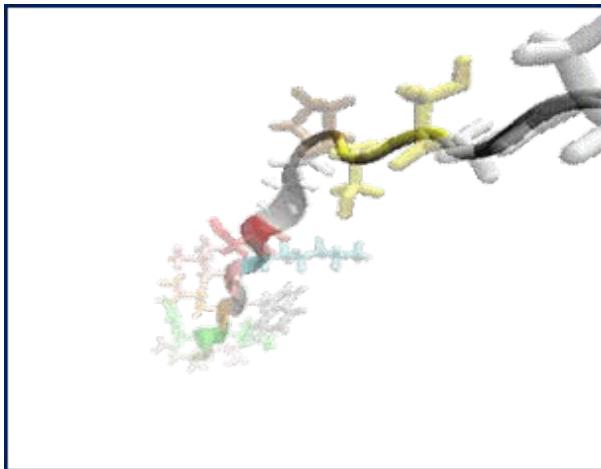
Blur effect



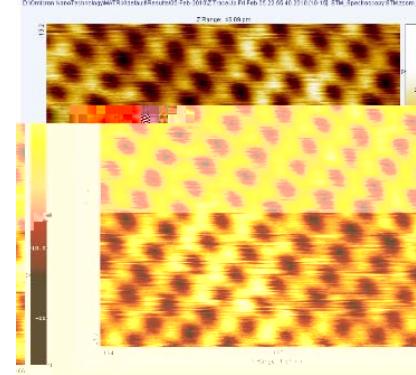
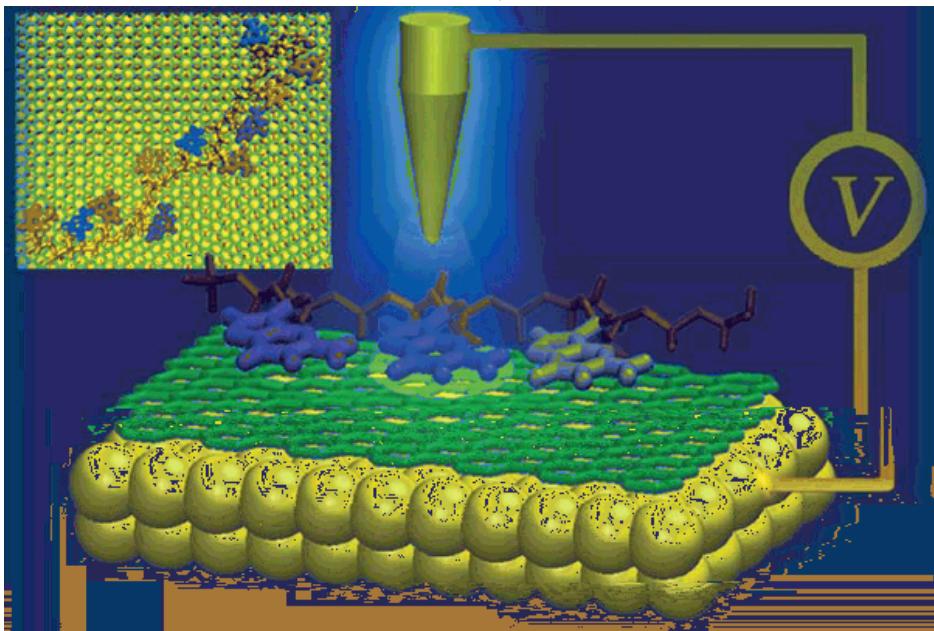
$T(w)=36$  layer

- CIS calculation provide accurate electronic structures for very long CNT
- CIS found bloch-alike wavefunction in CNT without assuming any periodical boundary conditions.

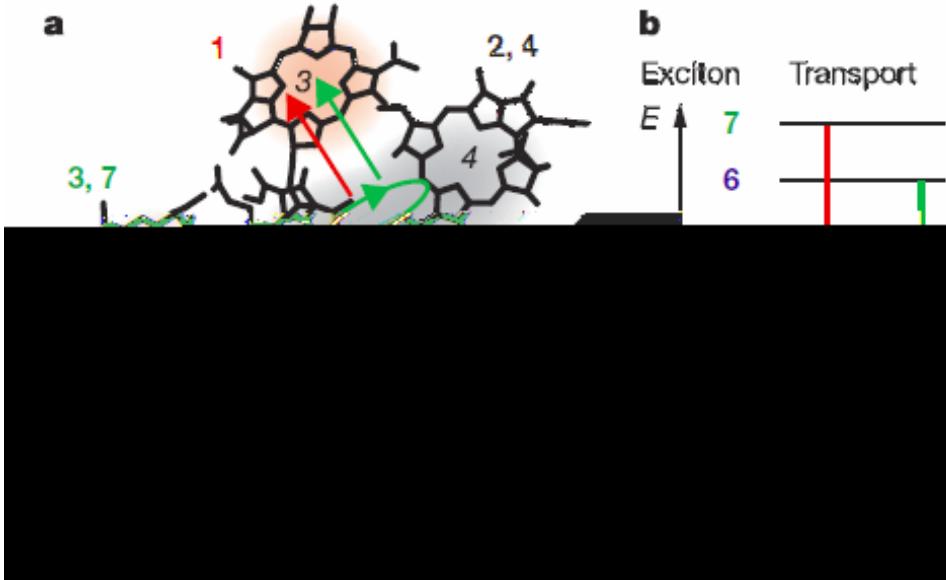
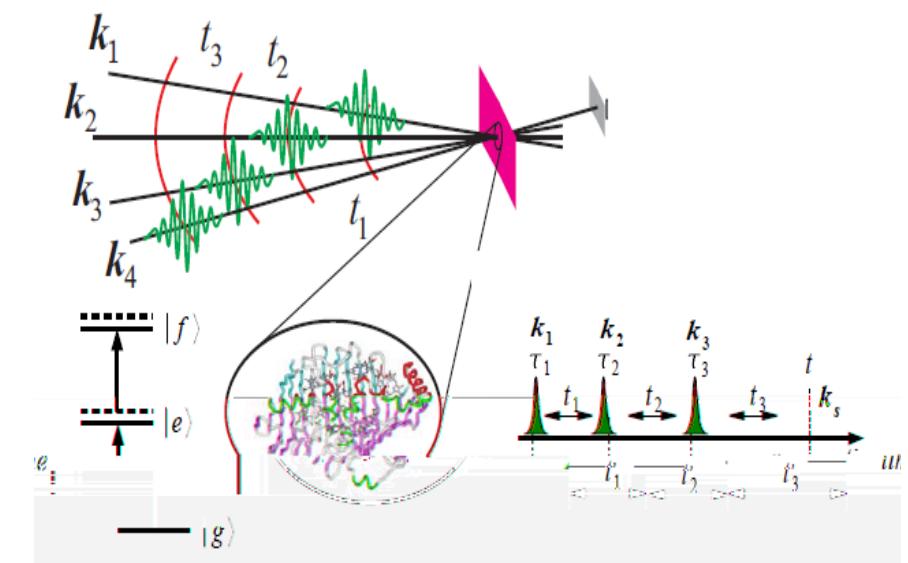




GFM



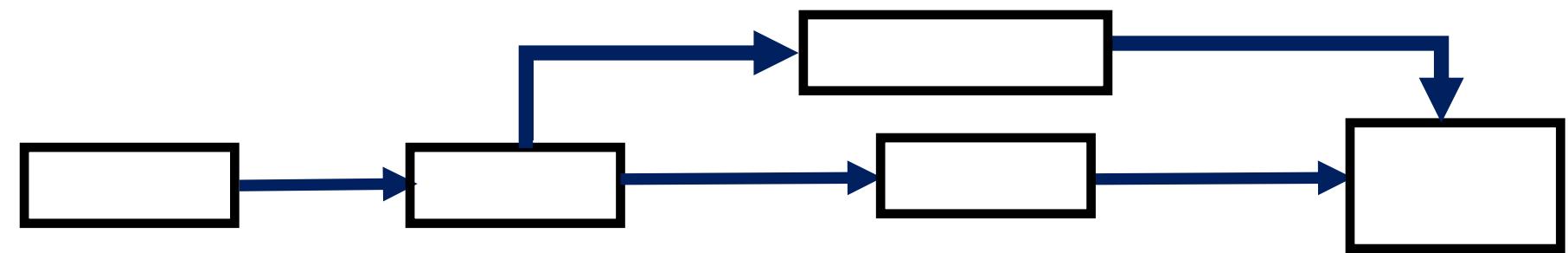
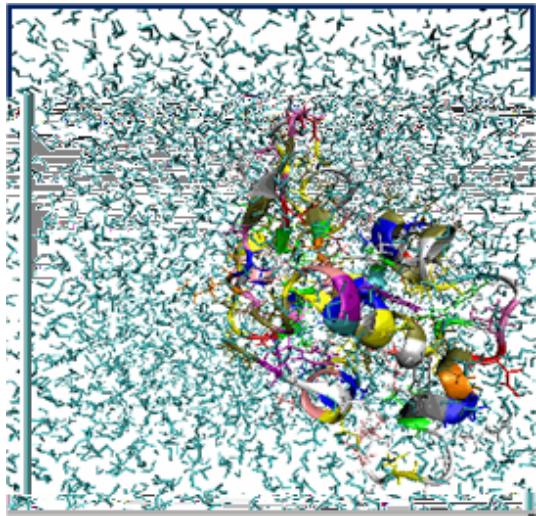
STM image of a single-walled carbon nanotube



Graham Fleming, Tobias Jungen, Berkeley

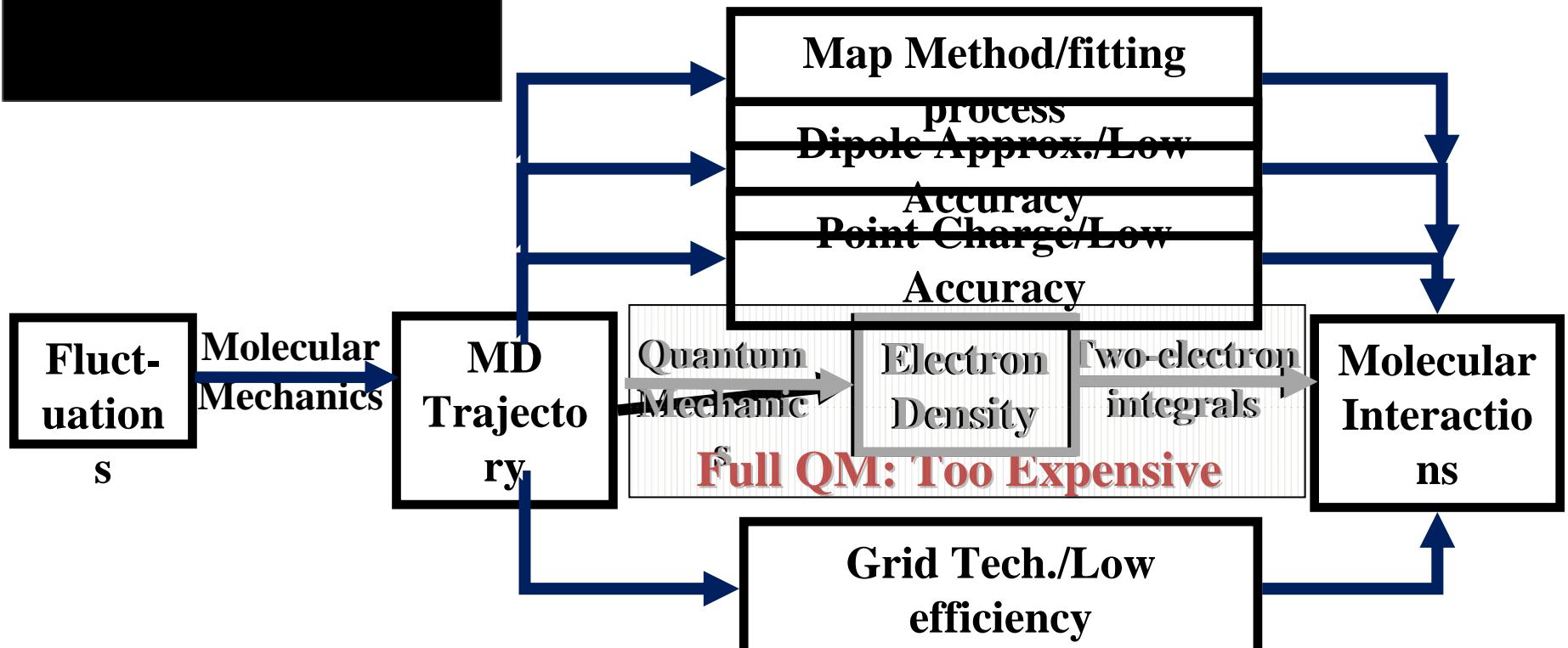
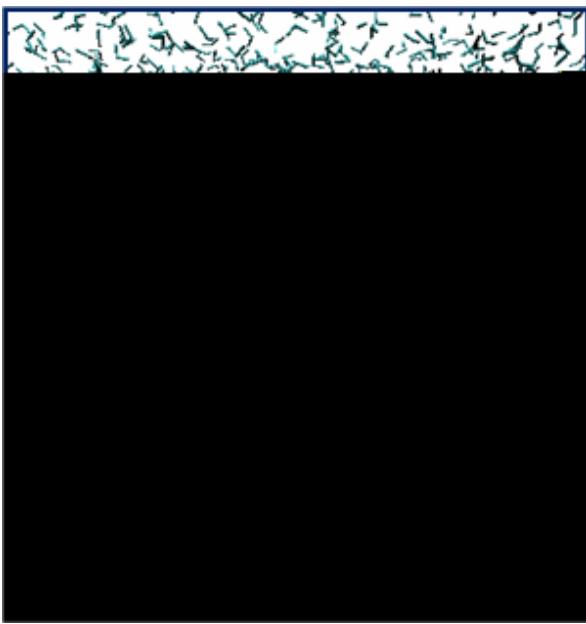
9<9. fl 9 Whcb <aa] ` hcb] ab k] h\ 9` Whf cghah] W  
: ` i Whi ah] cbg Ł fl Ł

GD97HFCB 9<9.



# Challenge for theoretical study

Fluctuations of proteins and environment: Change spectral features



# Map method for environmental fluctuations

Energy of a state can be expressed as

$$E = E_0 - \mu \cdot F$$

Map for excitation energy are constructed by

$$\varepsilon_m^F = \varepsilon_m + \sum_i \alpha_i F_i + \sum_i \beta_i \frac{dF_i}{dx} + \dots$$

or

$$\varepsilon_m^F = \varepsilon_m + \sum_i \alpha'_i K_i + \sum_i \beta'_i K_i^2 + \dots$$

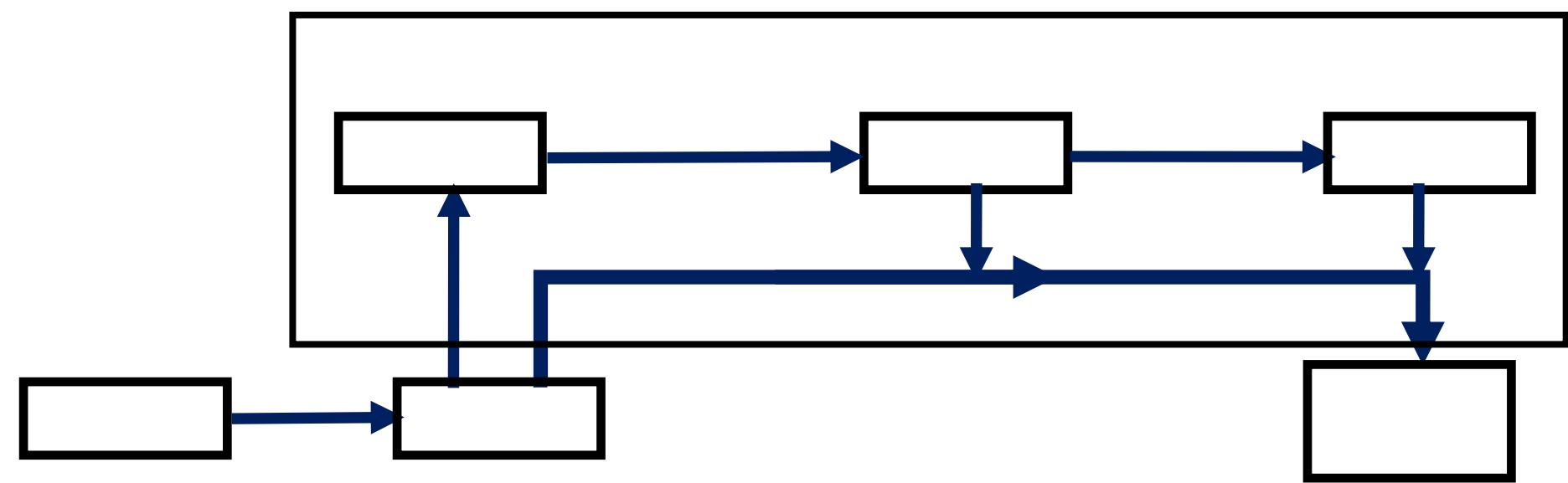
Advantages: Cheap computation costs & Easy to test  
--successful utilizations in 2DIR studies

Disadvantages:

Rely on many QM calculations and heavy analysis

Fitted parameters are not transferable

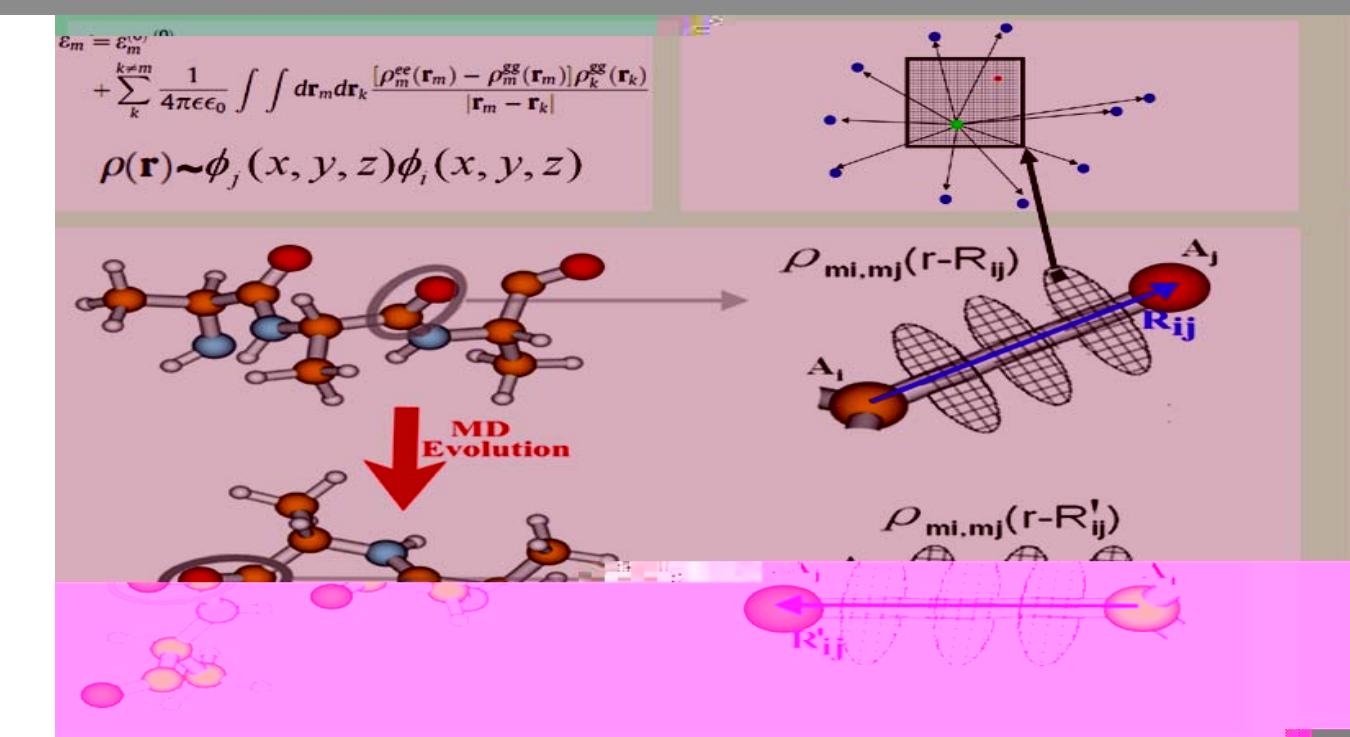
Lack of predict power



# EHEF (Exciton Hamiltonian with Electrostatic Fluctuations) algorithm

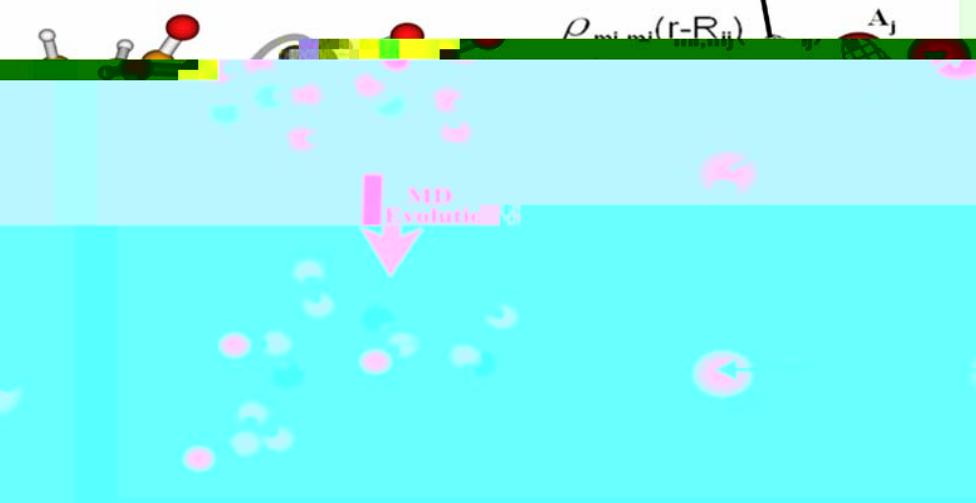
EHEF, uses a limited number of charge grids to produce the same local electrostatic potential as that induced by charge density, and converts the integration of charge density to the product of two charge grids. It avoid repeated QM studies and calculate interactions with low cost, so that efficiently construct QM-based exciton Hamiltonian for thousands of MM snapshots

$$\frac{1}{4\pi\epsilon\epsilon_0} \int \int d\mathbf{r}_m d\mathbf{r}_k \frac{[\rho_m^{ee}(\mathbf{r}_m) - \rho_m^{gg}(\mathbf{r}_m)]\rho_k^{gg}(\mathbf{r}_k)}{|\mathbf{r}_m - \mathbf{r}_k|} \cong \sum_{i,j} \frac{M_m^{eg}(i) M_k^{gg}(j)}{4\pi\epsilon\epsilon_0 r_{ij}}$$



# EHEF (Exciton Hamiltonian with Electrostatic Fluctuations) algorithm

$$\begin{aligned}\epsilon_m &= \epsilon_m^{(0)} \\ &+ \sum_k^{\neq m} \frac{1}{4\pi\epsilon_0} \int \int d\mathbf{r}_m d\mathbf{r}_k \frac{[\rho_m^{ee}(\mathbf{r}_m) - \rho_m^{gg}(\mathbf{r}_m)]\rho_k^{gg}(\mathbf{r}_k)}{|\mathbf{r}_m - \mathbf{r}_k|} \\ \rho(\mathbf{r}) &\sim \phi_j(x, y, z)\phi_i(x, y, z)\end{aligned}$$



1. Define atomic frame for each pair of atoms, ( $\mathbf{R}_i$ ,  $\mathbf{r}_j$ ,  $\mathbf{r}_k$ )
2. Calculations involve charge densities (As expensive as calculating 2-e integrals)
3. To obtain low computation cost, use a limited number of large grids ( $10^3$  for an amino acid).
4. Each Grid is divided into a large number of sub-grids ( $10^6$  for an amino acid), to maintain the same precision as 2-e integrals.
5. Electrostatic potentials at random sample positions are calculated from charges in  $10^3$  sub-grids, and are used to fit one point charge for one grid.
6. Charge distributions are saved in the atomic frame.
7. Assume fixed charge distribution in the atomic frame
8. Update the new atomic frame during MD evolutions ( $\mathbf{R}'_i$ ,  $\mathbf{r}'_j$ ,  $\mathbf{r}'_k$ )
9. Update the charge density in the new atomic frame.
10. Calculate energy shift due to fluctuations.

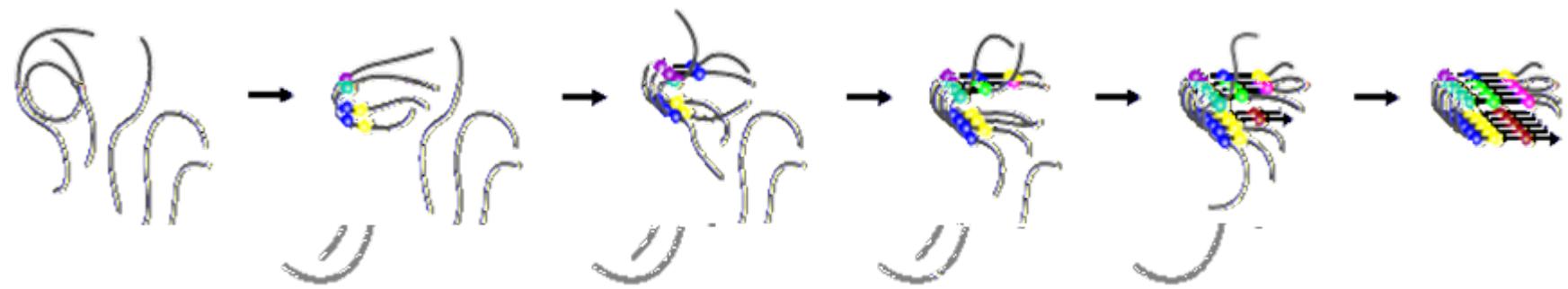
# EHEF (Exciton Hamiltonian with Electrostatic Fluctuations) algorithm

**Accurate -- calculations without fitting parameters**

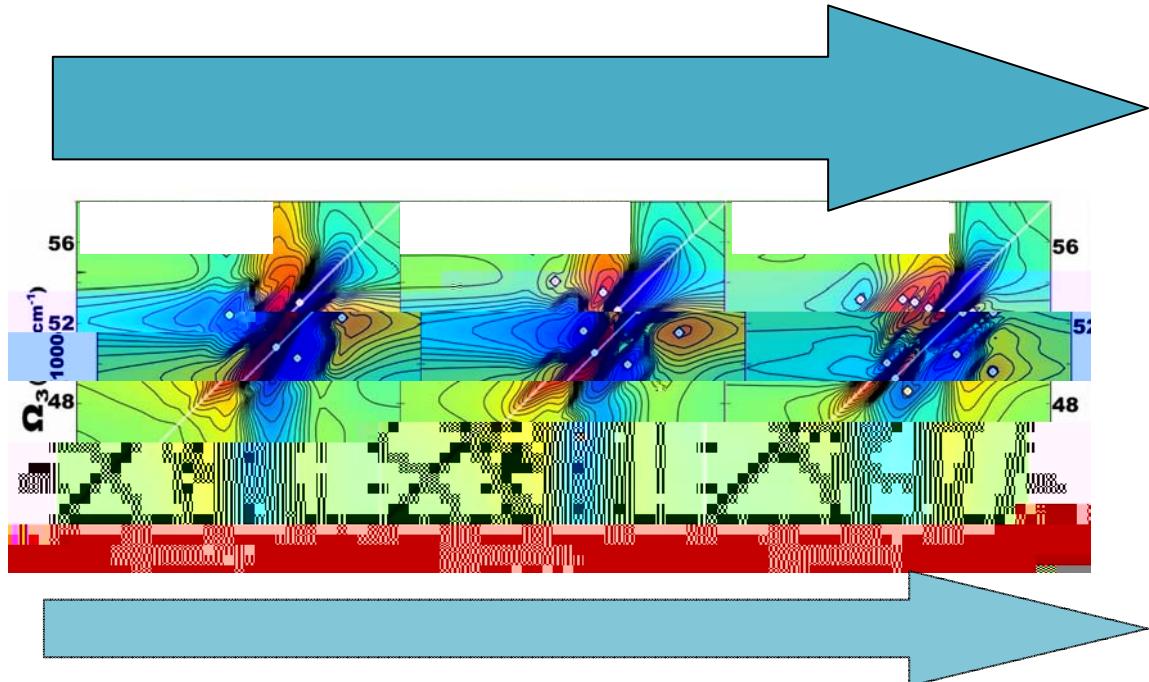
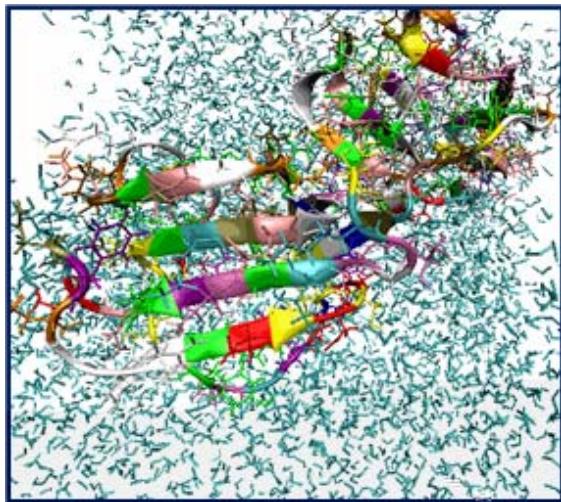
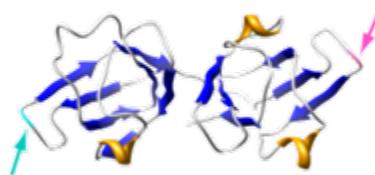
**Efficient -- Atomic frame to avoid repeated QM studies, Grid charge fitted from electrostatic potential avoid integrations**

**Transferable for different systems and applications**

**Able to explain experiments and predict new information**



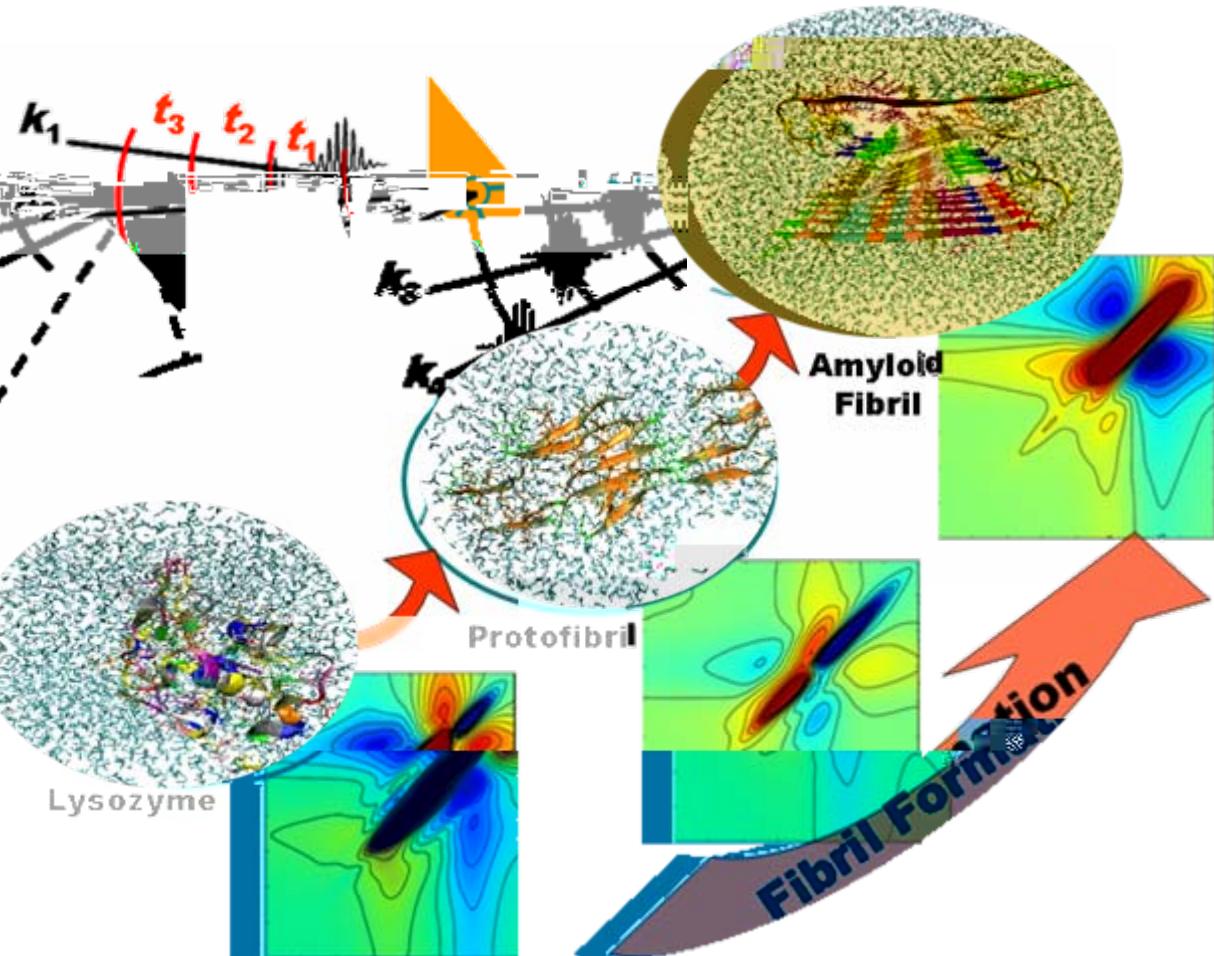
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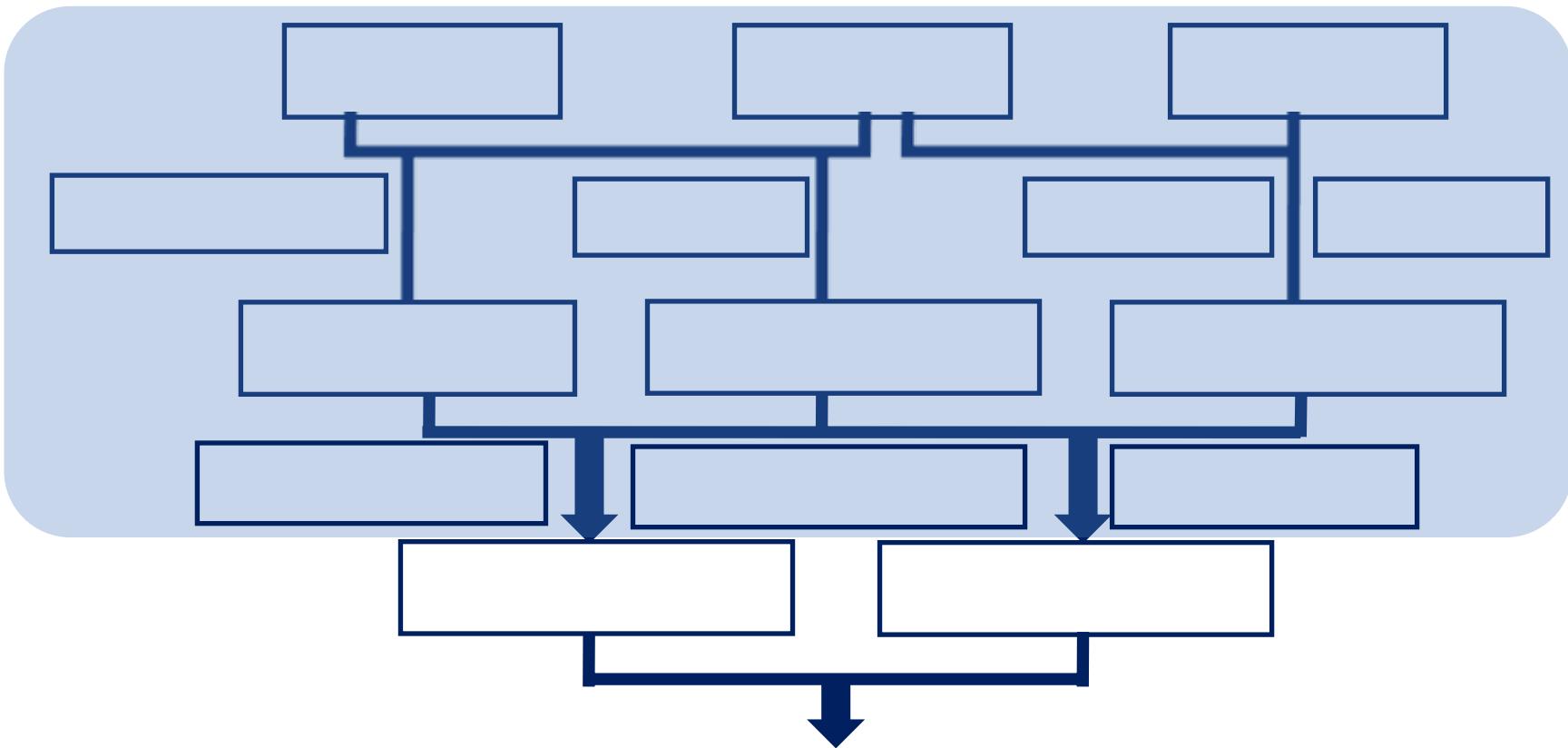
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X-ray NMR CD



Jiang, Mukamel Angew. Chem. Int. Ed 49, 9666(2010); J. Am. Chem. Soc., 132, 7769(2010);  
J. Phys. Chem. B, 114, 12150(2010);



- (Goddard:Caltech, Martinez:Stanford,  
Steinhauser:Fraunhofer )

- Gaussian70 98

- G] XbYm Y] d : "

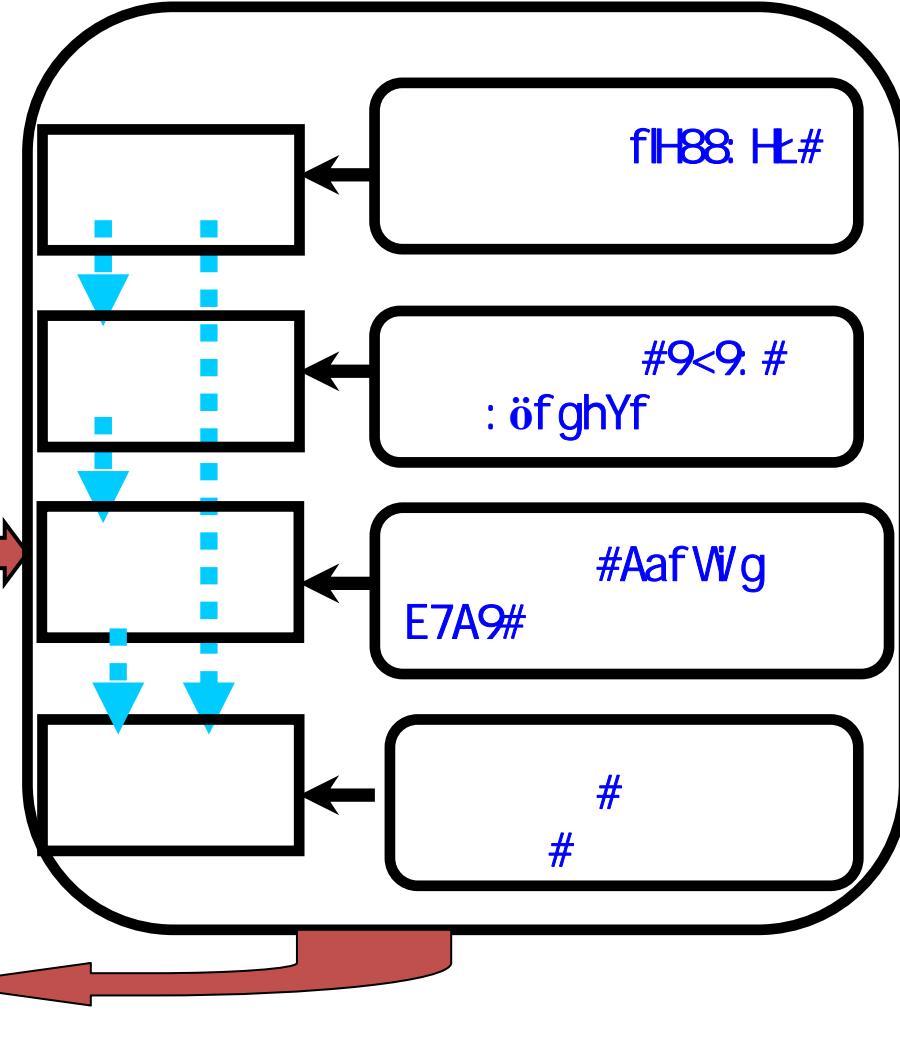
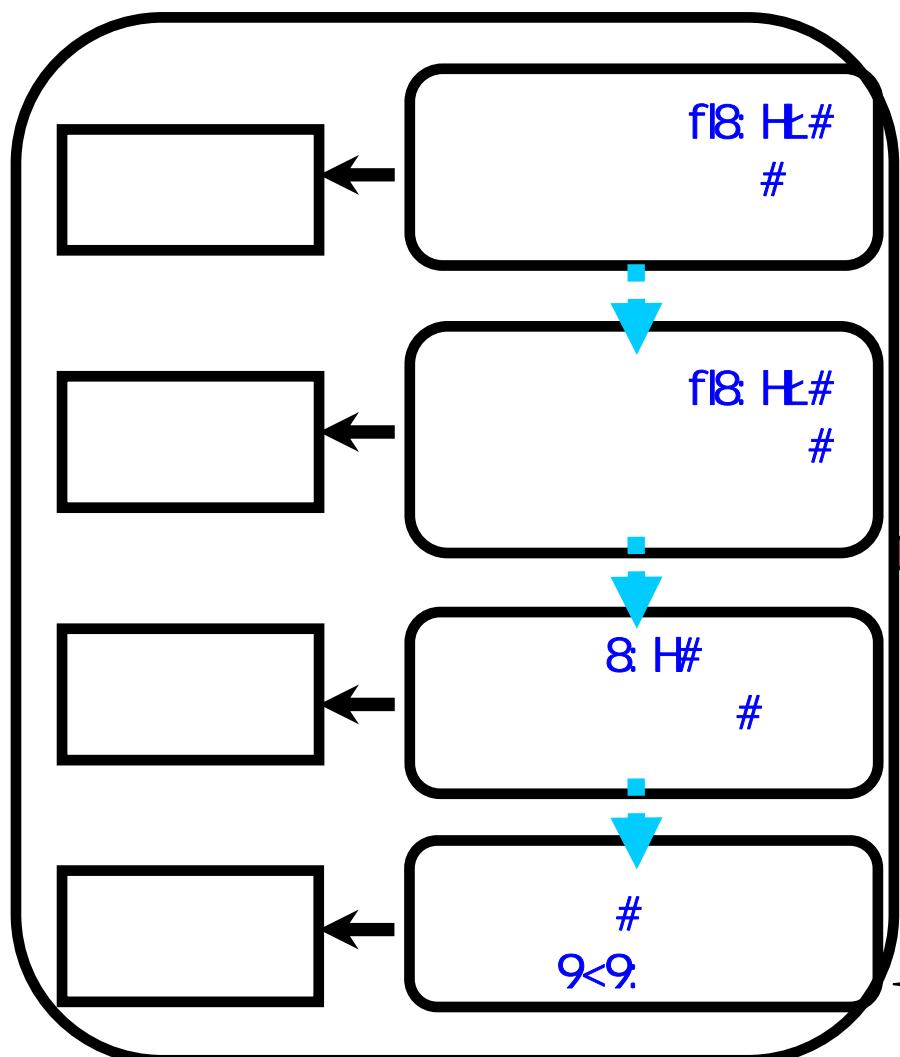
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Dr. Daniel Healion

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