

<http://power.itp.ac.cn/~suncp/quantum.htm>





What is Life ?, Cambridge, 1946.

“

”



Erwin Schrödinger
1887-1961

DNA

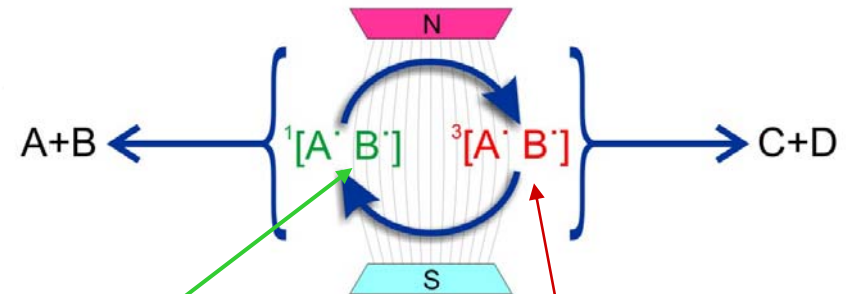
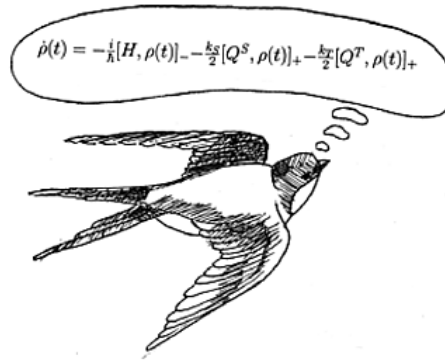


free radical pair

Biophysical Journal Volume 78 February 2000 707-718

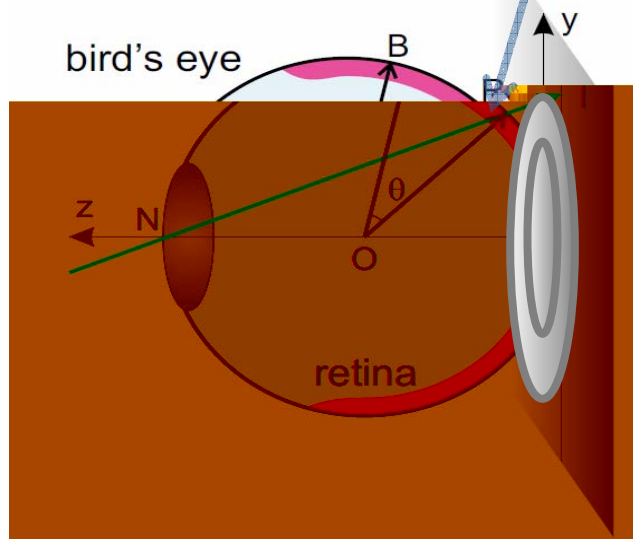
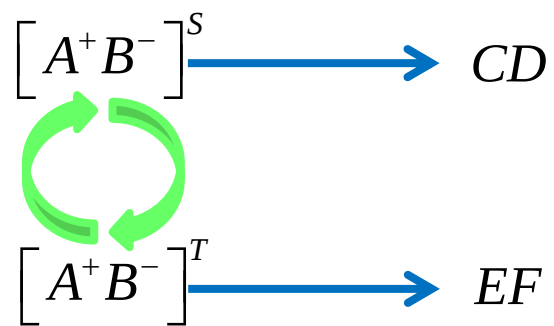
A Model for Photoreceptor-Based Magnetoreception in Birds

Thorsten Ritz, Salih Adem, and Klaus Schulten



$$|0,0\rangle = (\uparrow\downarrow - \downarrow\uparrow)/\sqrt{2} \quad s = 0 \quad (\text{singlet})$$

$$\begin{aligned} |1,1\rangle &= \uparrow\uparrow \\ |1,0\rangle &= (\uparrow\downarrow + \downarrow\uparrow)/\sqrt{2} \\ |1,-1\rangle &= \downarrow\downarrow \end{aligned} \quad s = 1 \quad (\text{triplet})$$



Chemical compass model of avian magnetoreception

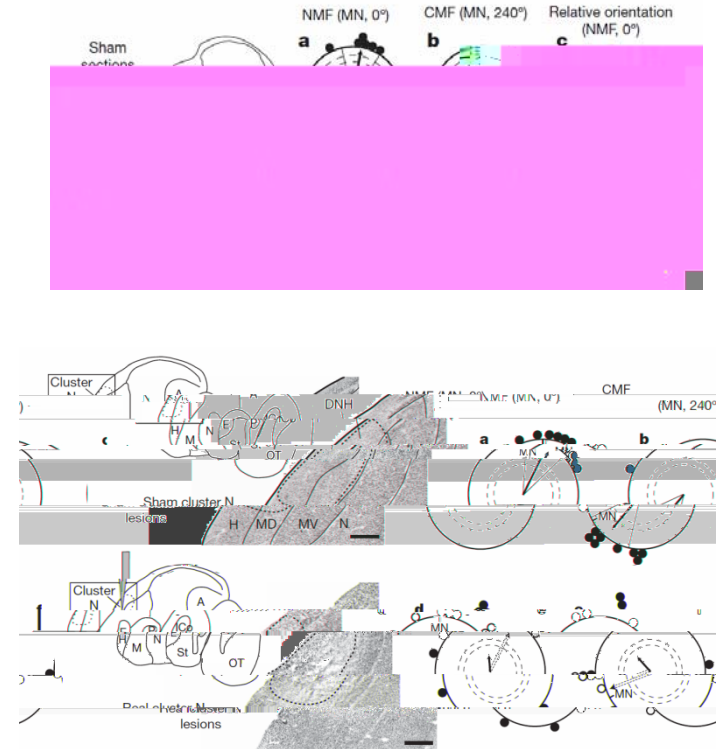
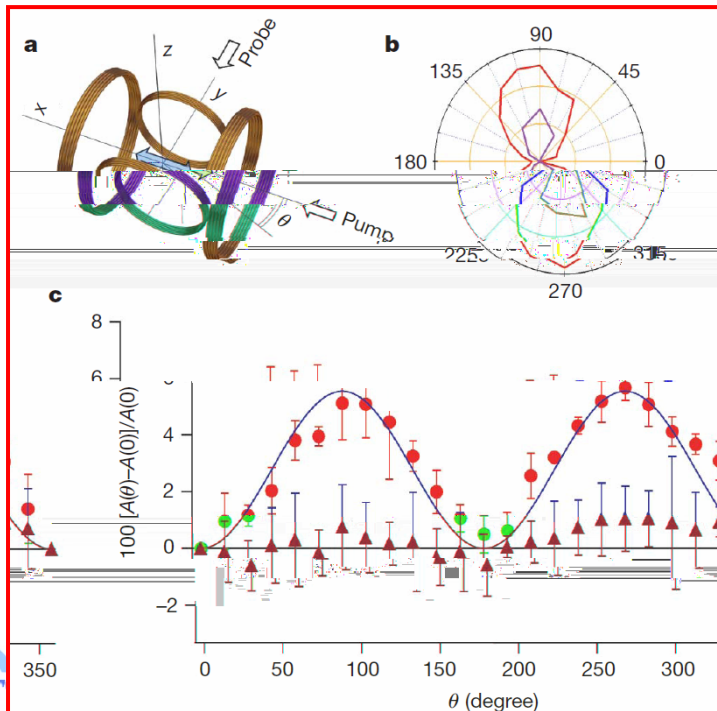
Liddell³, Devens Gust³, Christiane R. Timmel¹ & P. J. Hore²

Nature, Vol 453| 15 May 2008

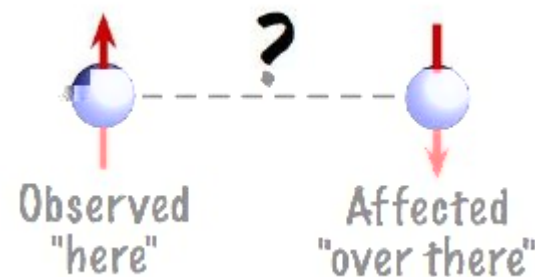
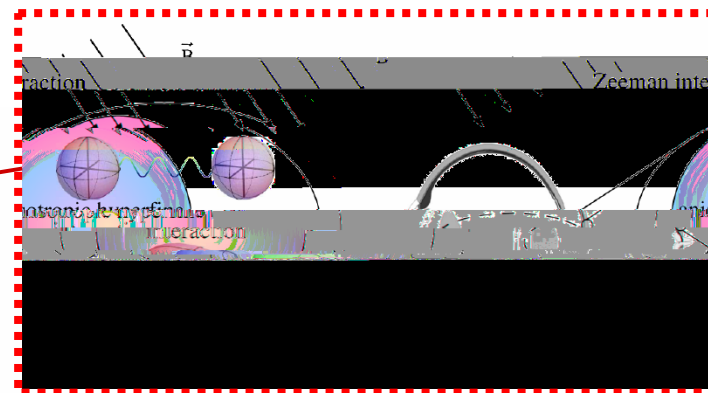
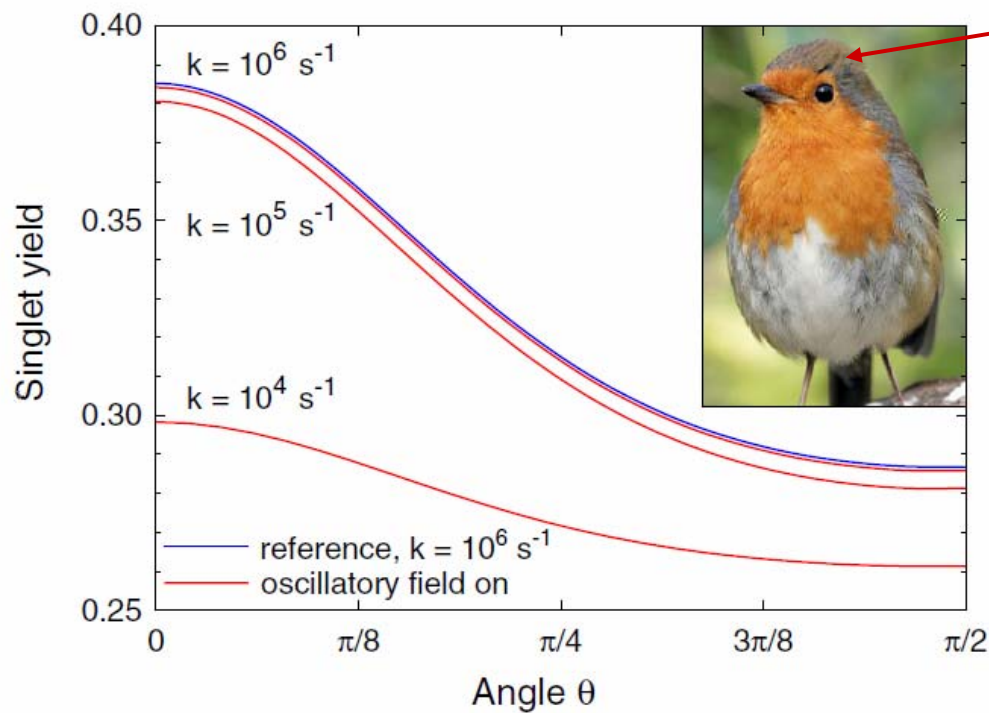
Visual but not trigeminal mediation of magnetic compass information in a migratory bird

Manuela Zapka¹, Dominik Heyers¹, Christine M. Hein¹, Svenja Engels¹, Nils-Lasse Schneider¹, Jörg Hans¹, Simon Weller¹, David Dwyer¹, Dmitry Kisilov¹, J. Martin Wild² & Henrik Mouritsen¹

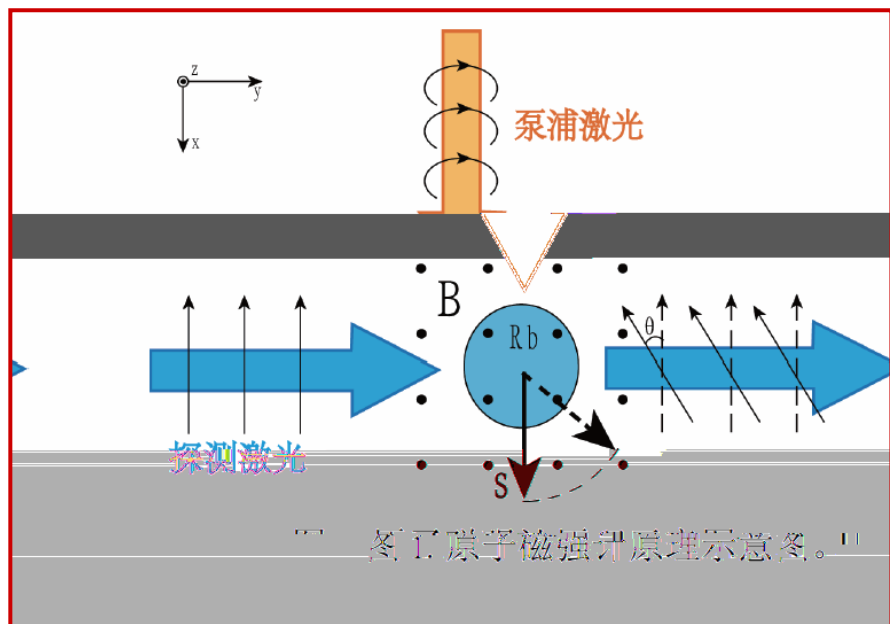
Nature, Vol 461| 29 October 2009



Guager et al , Phys. Rev. Lett. 106, 040503 (2011)

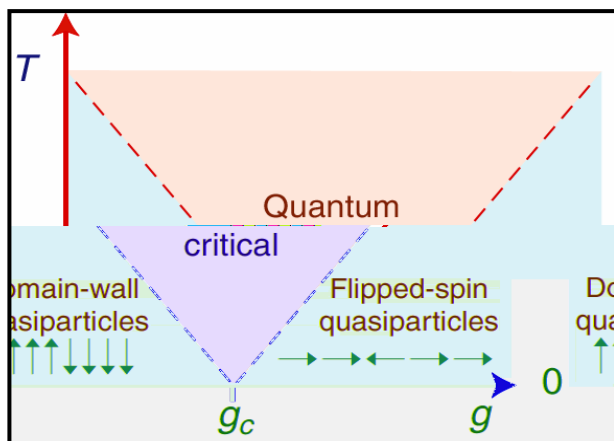


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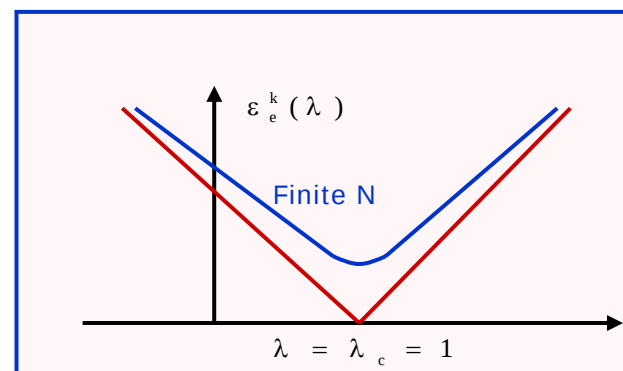


Romalis
 $10 \text{ fT} / \sqrt{\text{Hz}}$
SQUID

2002
 $2 \text{ aT} / \sqrt{\text{Hz}}$



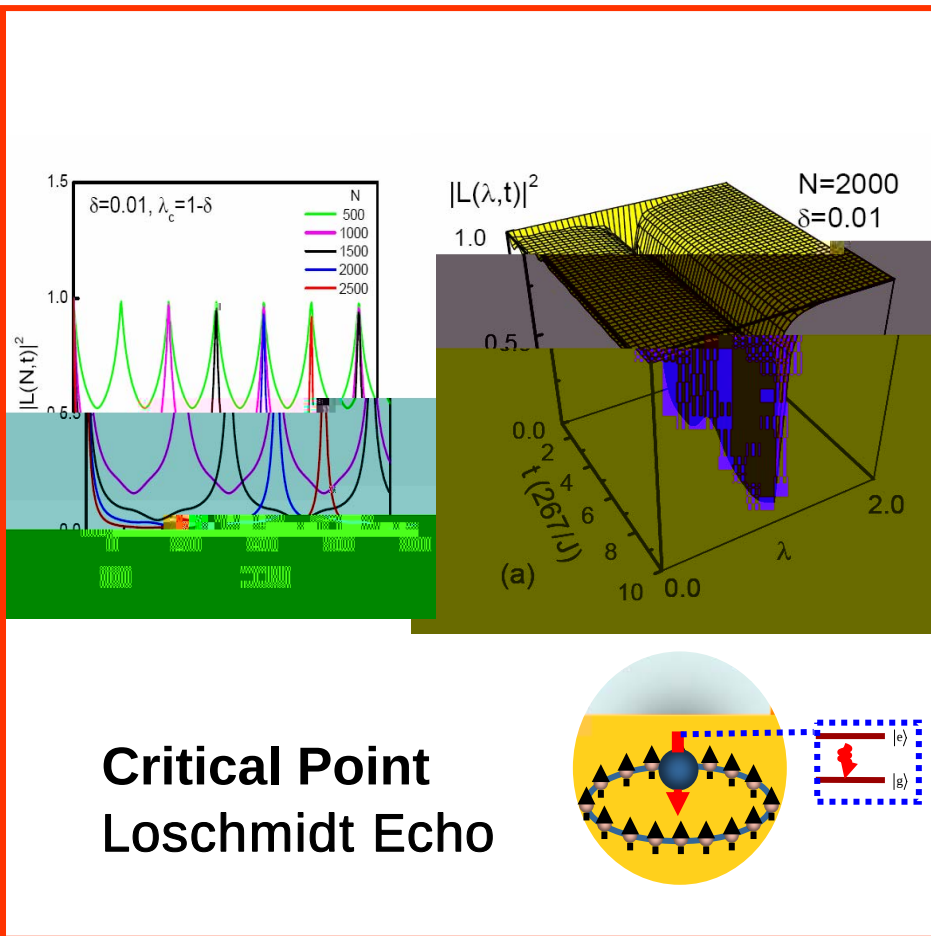
S Sachdev, (2000)
Quantum Phase Transitions.
Cambridge University Press



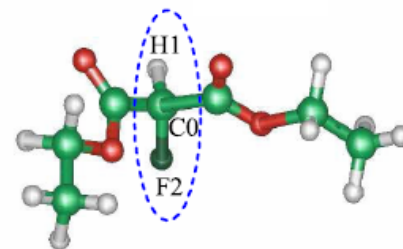
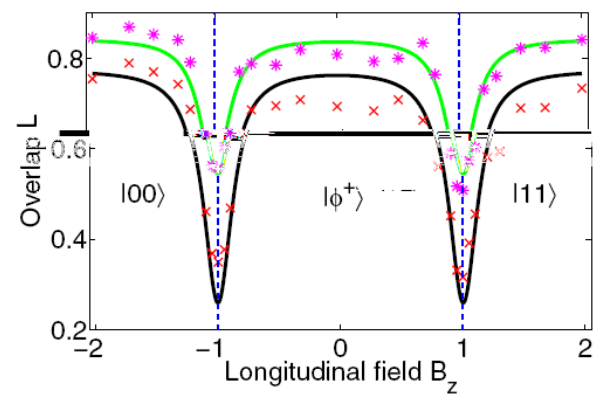
$$\varepsilon_e^k(\lambda) = 2J\sqrt{1 + \lambda^2 - 2\lambda \cos(ka)}$$

Quan, Song, Liu, Zanardi, Sun Phys. Rev. Lett. 96, 140604 (2006)

Suter



PRL 100, 100501 (2008)



Quan et al , Phys. Rev. Lett. 96, 140604 (2006)-

[6]

2006

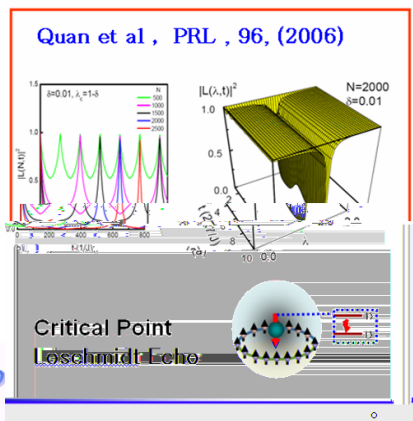
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Laflamme

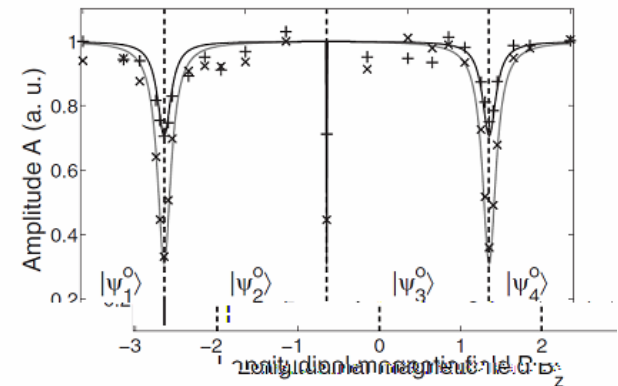
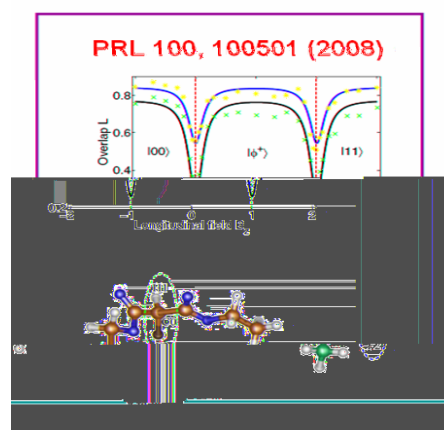
Phys.Rev. A 79,012305 (2009)

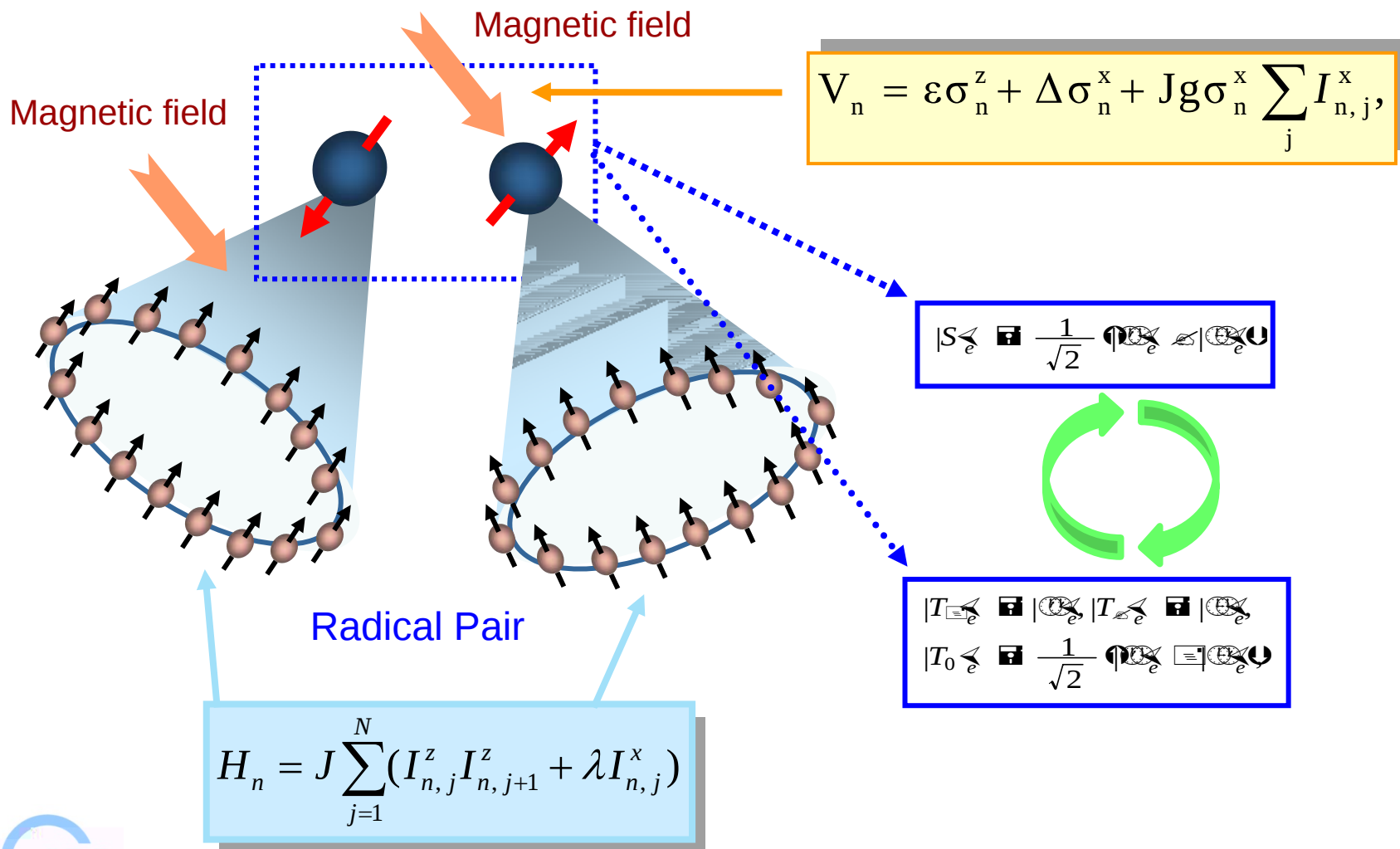
There has been a recent flurry of activity following the observation [6] that the proximity to a quantum critical point enhances the sensitivity of a system to external perturbations, as measured by quantum-information-theoretical quantities.

我们的理论



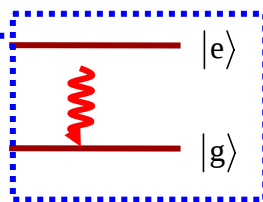
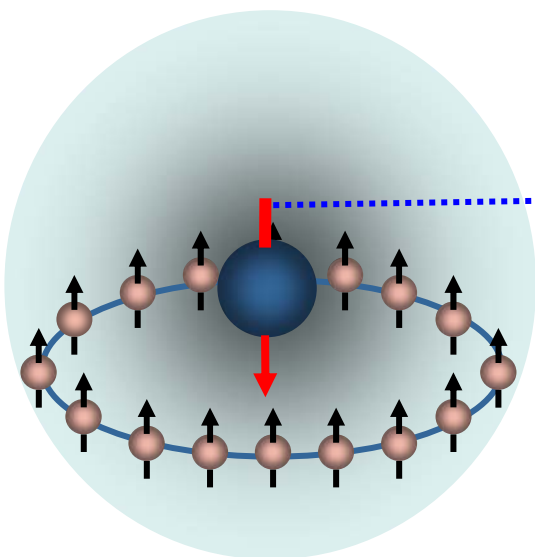
德国Suter小组的实验证实





Hepp-Coleman

$$H = H_0 + H_1 = -J \sum_j (\sigma_j^z \sigma_{j+1}^z + g \sigma_j^x + \delta |e\rangle \langle e| \sigma_j^x)$$



$$H_e = H_e(\lambda) = H_g + V_e$$

$$V_e = -J\delta \sum_j \sigma_j^x$$

$$H_g = H_0 = -J \sum_j (\sigma_j^z \sigma_{j+1}^z + g \sigma_j^x)$$

$$|\psi(0)\rangle = (c_g |g\rangle + c_e |e\rangle) \otimes |G\rangle_g \longrightarrow |\psi(t)\rangle = c_g |g\rangle \otimes |\varphi_g(t)\rangle + c_e |e\rangle \otimes |\varphi_e(t)\rangle$$

$$\rho_s(t) = c_g c_g^* |g\rangle \langle g| + c_e c_e^* |e\rangle \langle e| + \langle \phi_e | \phi_g \rangle c_g c_e^* |g\rangle \langle e| + \langle \phi_g | \phi_e \rangle c_e c_g^* |e\rangle \langle g|$$

Loschmidt echo

$$[\rho_s(t)]_{\text{eg}} = c_g c_e^* D(t)$$

$$D(t) = \langle \varphi_g(t) | \varphi_e(t) \rangle \quad |\varphi_\alpha(t)\rangle = \exp(-iH_\alpha t) |G\rangle_g$$

(Loschmidt echo)

$$L(\lambda, t) = \left| \langle \phi_g(t) | \phi_e(t) \rangle \right|^2 = \left| \langle G |_g \exp(iH_g t) \exp(-iH_e t) | G \rangle_g \right|^2$$

$$L(\lambda, t) = \prod_{k>0} [1 - \sin^2(2\alpha_k) \sin^2(\epsilon_e^k t)].$$

$$H_e = -J \sum_j [\sigma_j^z \sigma_{j+1}^z + (g + \delta) \sigma_j^x]$$

$$= \sum_k \varepsilon_e^k (A_k^\dagger A_k - 1/2)$$

$$H_g = -J \sum_j (\sigma_j^z \sigma_{j+1}^z + g \sigma_j^x)$$

$$= \sum_k \varepsilon_g^k (B_k^\dagger B_k - 1/2)$$

$$A_k |G\rangle_e = 0$$

$$B_k |G\rangle_g = 0$$



$$B_{\pm k} = \cos(\alpha_k) A_{\pm k} - i \sin(\alpha_k) (A_{\mp k})^\dagger,$$

BCS-like ground state:

$$|G\rangle_g = \prod_{k>0} [i \cos(\alpha_k) + \sin(\alpha_k) A_k^\dagger A_{-k}^\dagger] |G\rangle_e$$

H. T. Quan, Z. Song, X. F. Liu, and C. P. Sun, PRL 96, 140604 (2006)

Born-Oppenheimer

$$\begin{aligned} & \left| \begin{array}{c} \text{Hand} \\ \text{Eye} \end{array} \right\rangle \cos \frac{\theta}{2} \left| \begin{array}{c} \text{Clock} \\ \text{Hand} \end{array} \right\rangle \sin \frac{\theta}{2} \left| \begin{array}{c} \text{Hand} \\ \text{Eye} \end{array} \right\rangle \\ & \left| \begin{array}{c} \text{Hand} \\ \text{Eye} \end{array} \right\rangle \sin \frac{\theta}{2} \left| \begin{array}{c} \text{Clock} \\ \text{Hand} \end{array} \right\rangle \cos \frac{\theta}{2} \left| \begin{array}{c} \text{Hand} \\ \text{Eye} \end{array} \right\rangle \end{aligned}$$

$$\frac{\gamma}{2} \tan^{-1} \left(\frac{F_2}{F_1} \right).$$

$| \rangle$

Loschmidt Echo (I)

$$\rho_S(t) = \frac{1}{2} [|+-\rangle \langle +-| + |-+\rangle \langle -+| - D(t) |+-\rangle \langle -+| - D^*(t) |-+\rangle \langle +-|]$$

$$\rho_s(t) = \frac{1}{2} [|+-\rangle \langle +-| + |-+\rangle \langle -+| - D(t) |+-\rangle \langle -+| - D^*(t) |-+\rangle \langle +-|]$$

Loschmidt Echo

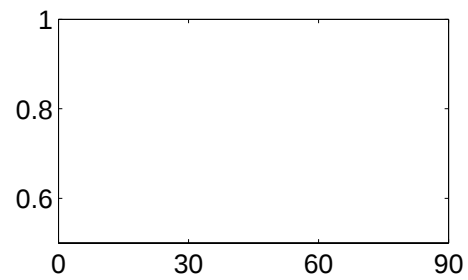
$$L(t) = \text{tr} |U(t) \rho_S U^\dagger(t)|^2$$

$$\dot{\rho}_S(t) = -\frac{i}{\hbar} [H_S, \rho_S(t)] + \sum_c \int_0^t r_c(t, \tau) \rho_S(\tau) d\tau,$$

$$f_S(t) = \langle \rho_S(t) | \rho_S(0) \rangle$$

$$r_c(t) = k_S \exp(-k_S t)$$

$$\Lambda(\theta) = \frac{\partial \Phi_s}{\partial \theta}.$$



Loschmidt Echo ()

$$f_S(t) = \frac{1}{2} \left(1 + \frac{1}{k_S} \int_0^t L(t-s) ds \right)$$

$$f_S(t) = \frac{1}{2} \left(1 + \frac{1}{k_S} \int_0^t L(t-s) ds \right)$$

$$L(t) = \prod_{k>0} \frac{A_k^2 + B_k^2}{\left[1 + \cosh(\beta \varepsilon_k^-)\right]^2},$$

$$A_k = \left[\cos^2 \alpha_k \cos(\varepsilon_k^+ - \varepsilon_k^-)t + \sin^2 \alpha_k \cos(\varepsilon_k^+ + \varepsilon_k^-)t \right] \cosh(\beta \varepsilon_k^-) + 1,$$

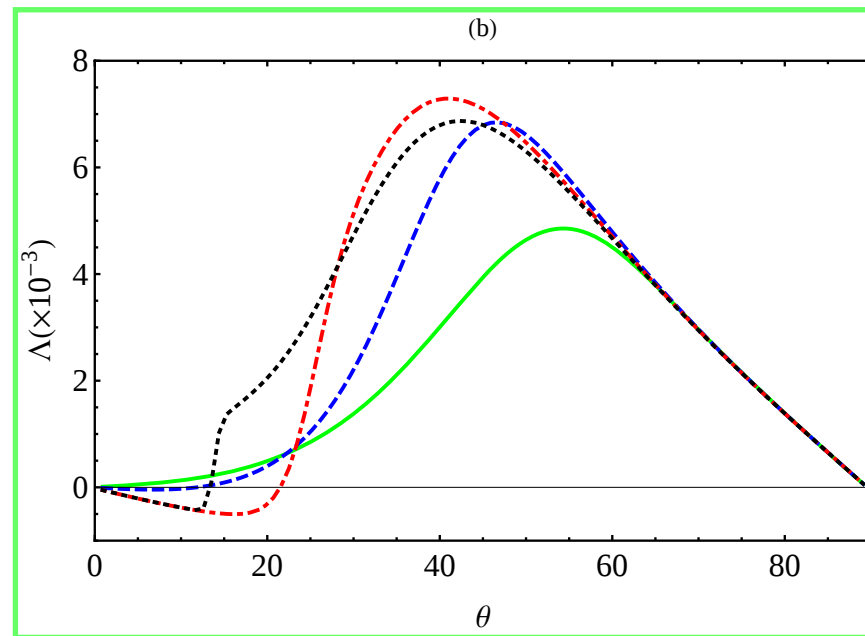
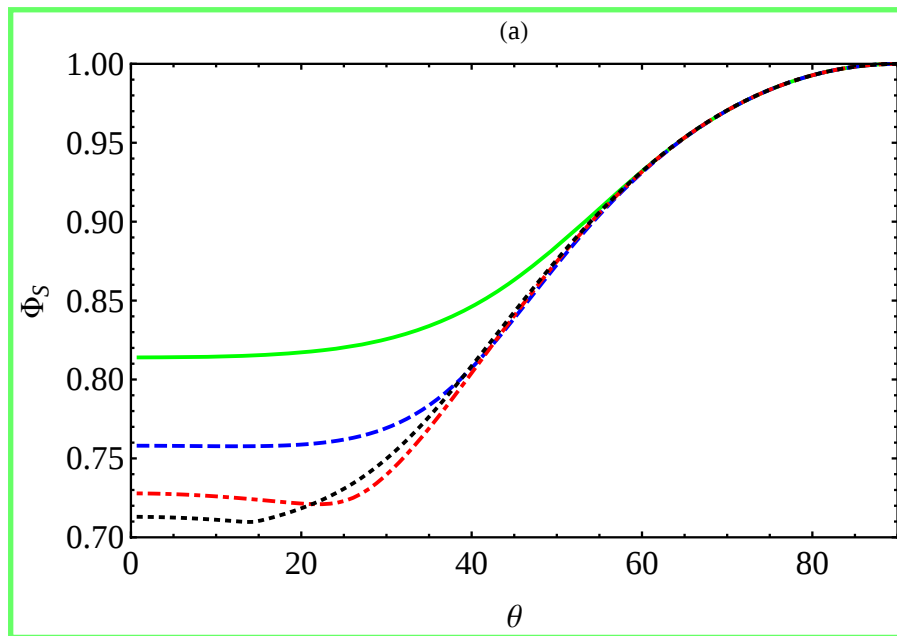
$$B_k = \left[\cos^2 \alpha_k \sin(\varepsilon_k^+ - \varepsilon_k^-)t - \sin^2 \alpha_k \sin(\varepsilon_k^+ + \varepsilon_k^-)t \right] \sinh(\beta \varepsilon_k^-),$$

$$\alpha_k = \frac{1}{2}(\theta_k^- - \theta_k^+),$$

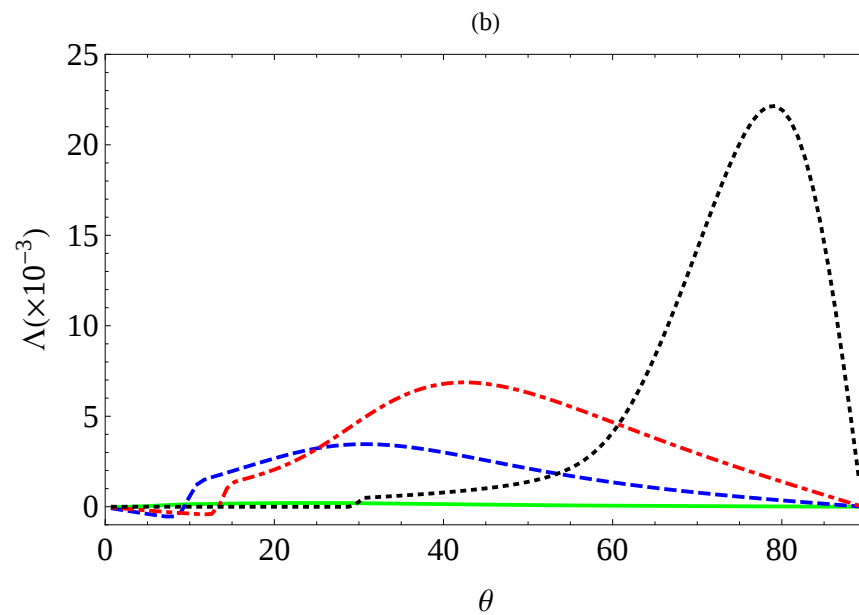
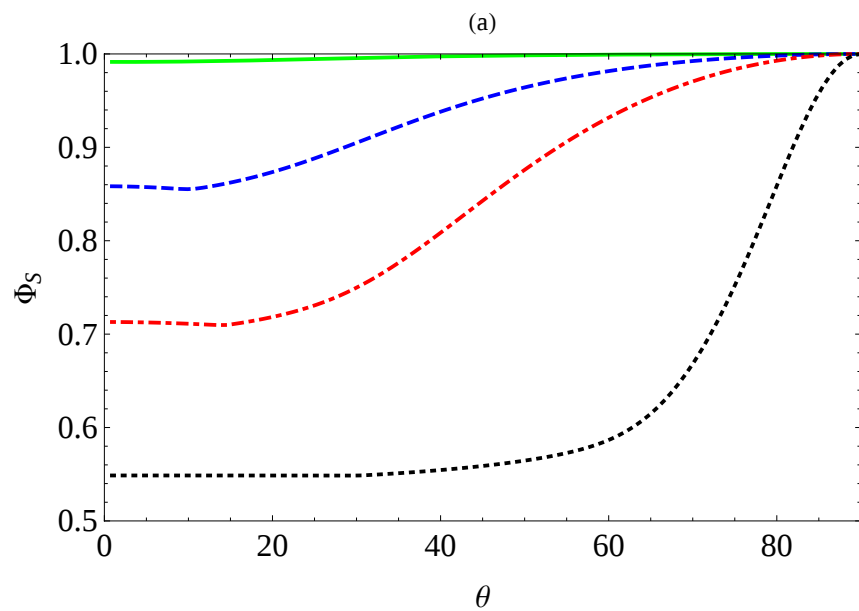
$$\theta_k^\pm = \tan^{-1}(\Delta_k / \eta_k^\pm),$$

$$\varepsilon_k^\pm = \sqrt{(\eta_k^\pm)^2 + \Delta_k^2},$$

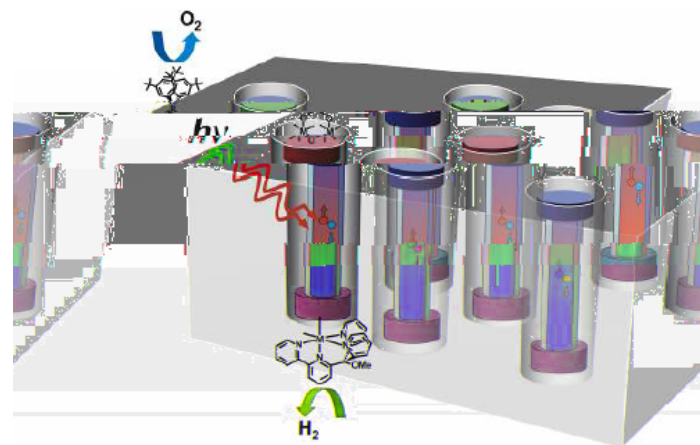
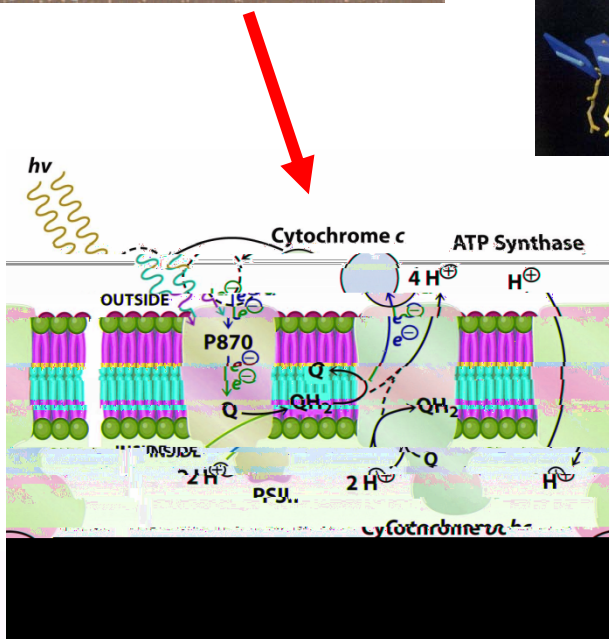
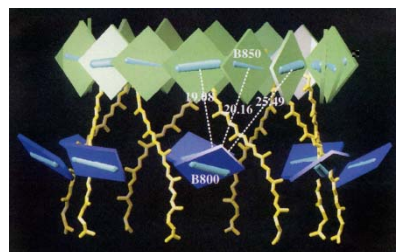
$$\eta_k^\pm = \eta_k(\lambda_\pm(\theta)).$$



- $N = 11$
- - - $N = 21$
- . - $N = 101$
- . . . $N = 1001$



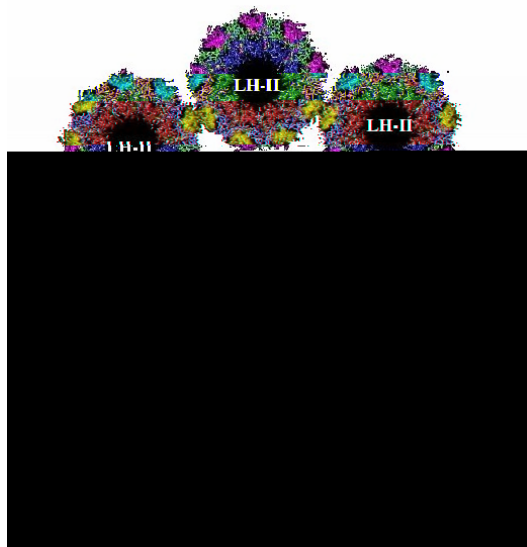
— $g_0 = 0.1$
- - $g_0 = 0.5$
- . - $g_0 = 1$
... $g_0 = 5$



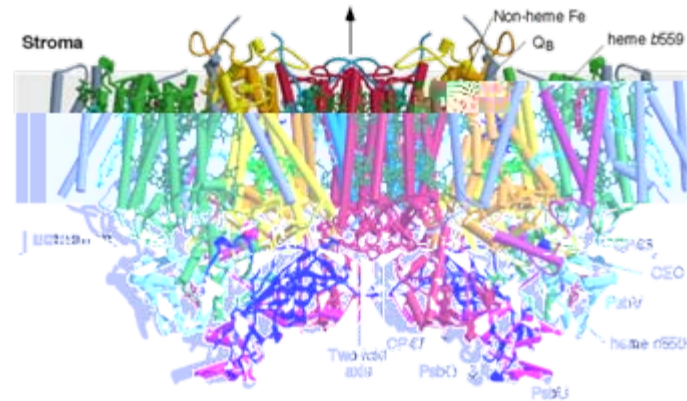
2009

LH II LH I

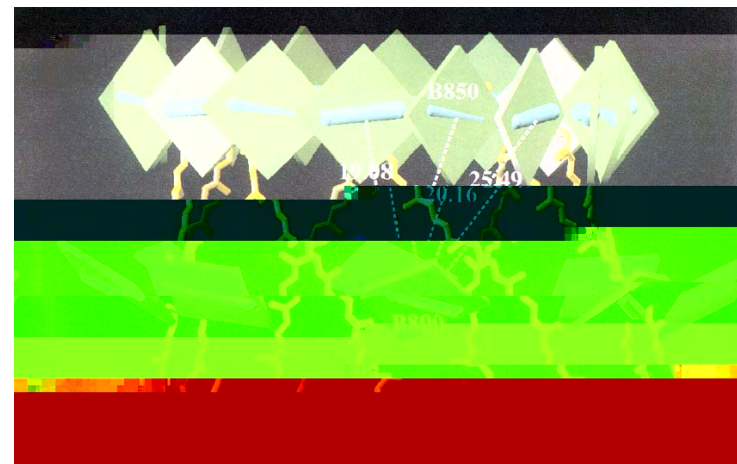
LHI I

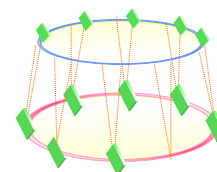


purple bacteria
(photosystem II)



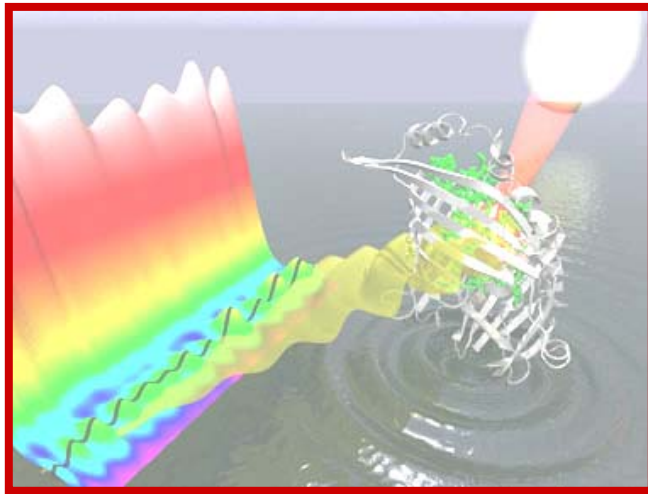
LH 11





Coherence Dynamics in Photosynthesis: Protein Protection of Excitonic Coherence

Science 8 June 2007:Vol. 316. no. 5830, pp. 1462

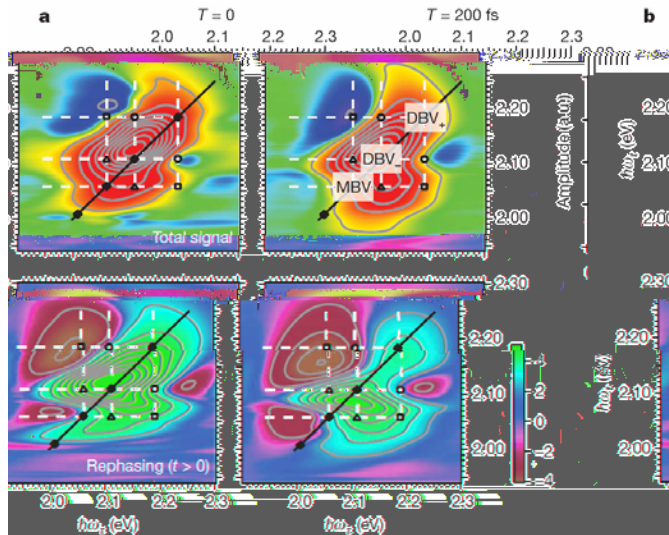


Fleming 2007
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(Fenna-
Matthew-Olson(FMO)
bacteri-ochlorophyll
complex) 660fs

Coherently wired light-harvesting in photosynthetic marine algae at ambient temperature

NATURE | Vol 463 | 644 | 2010



(marine cryptophyte algae)
(two Dimensional photon echo spectroscopy),

300fs

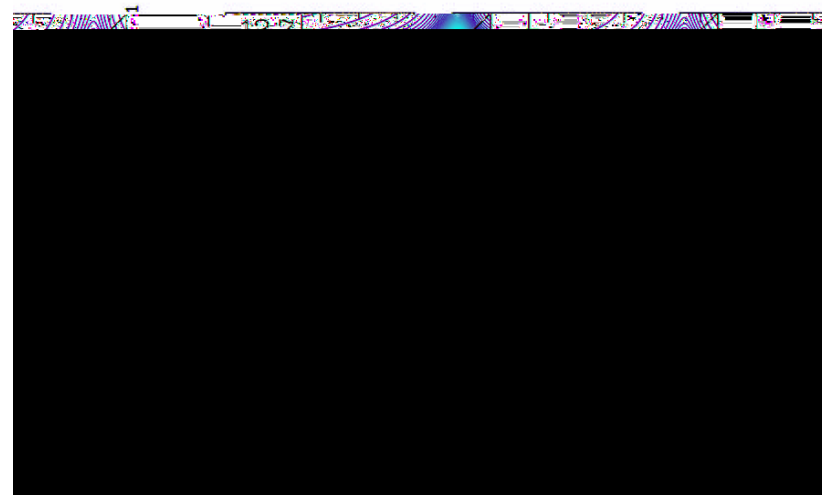
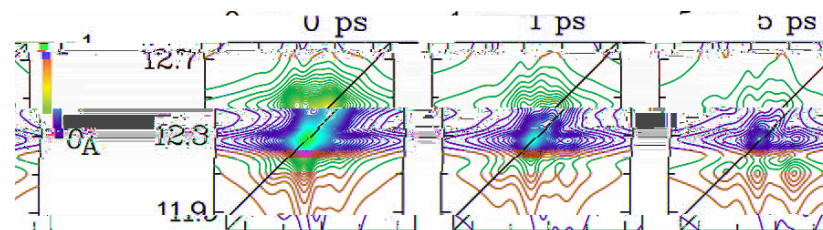
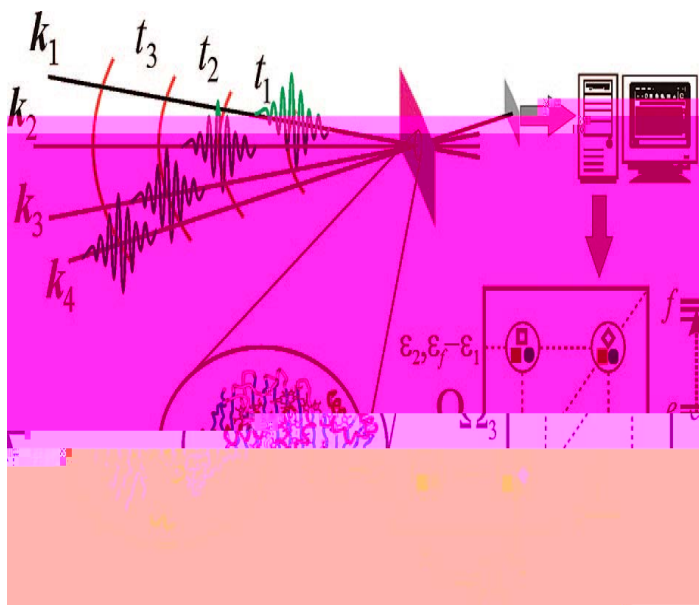
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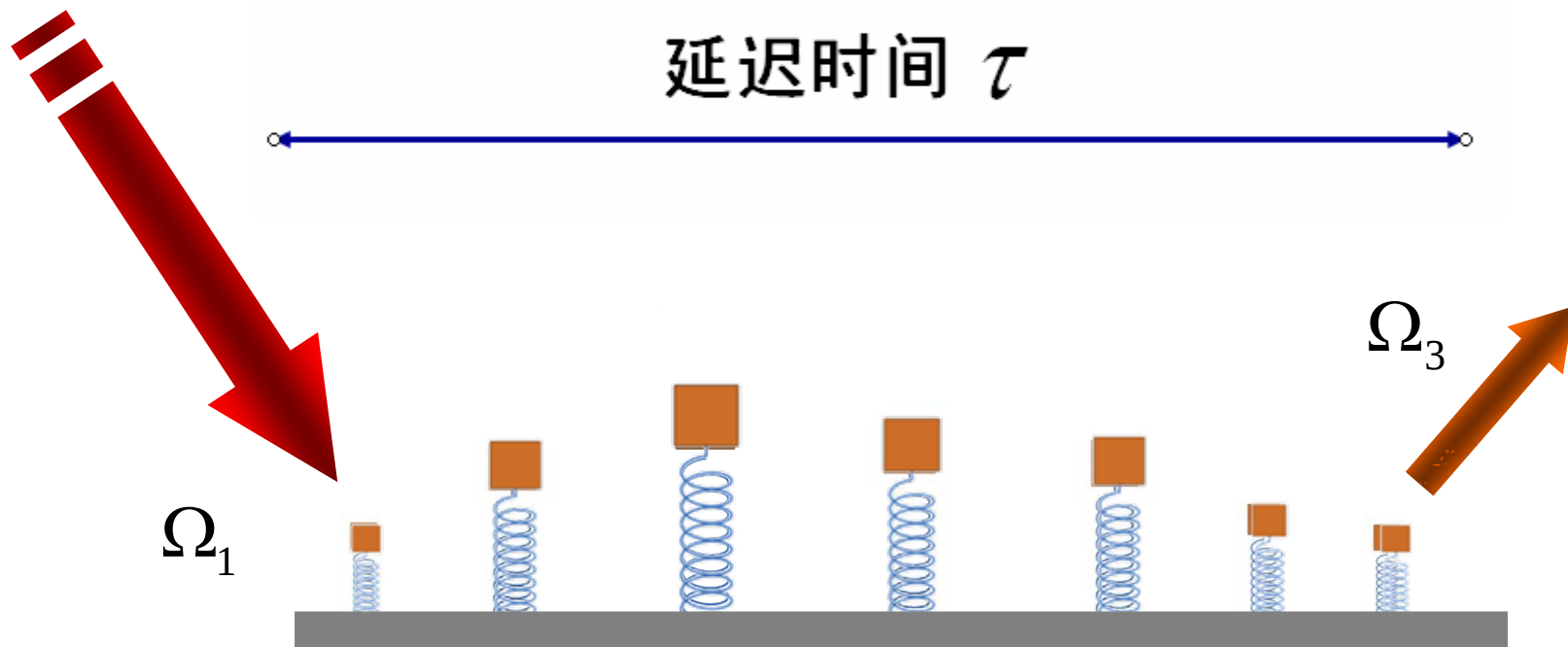
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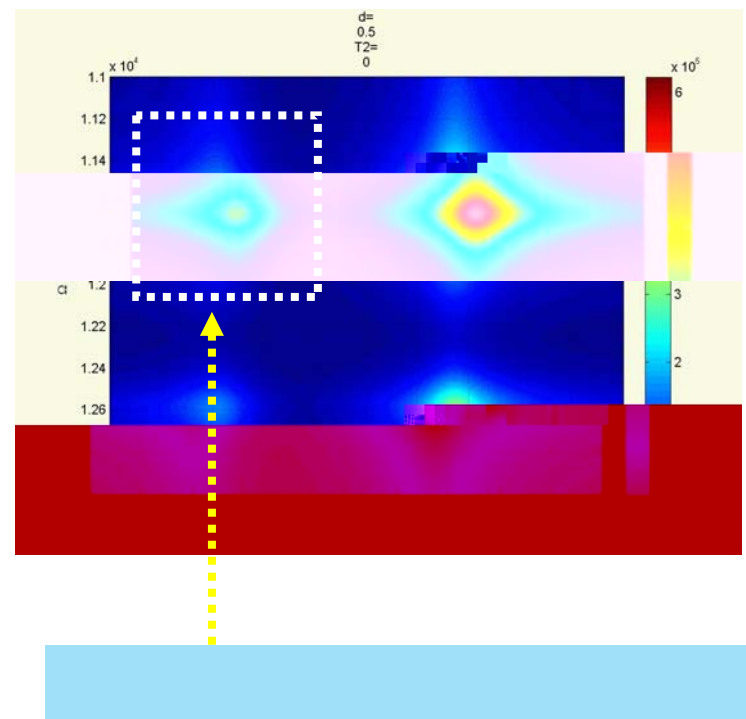
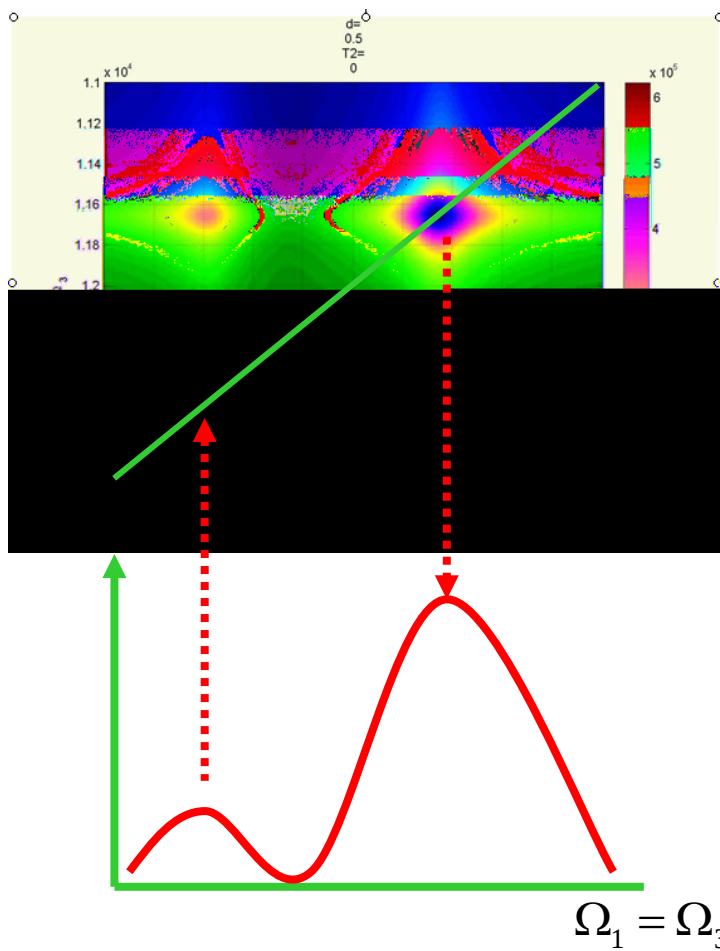
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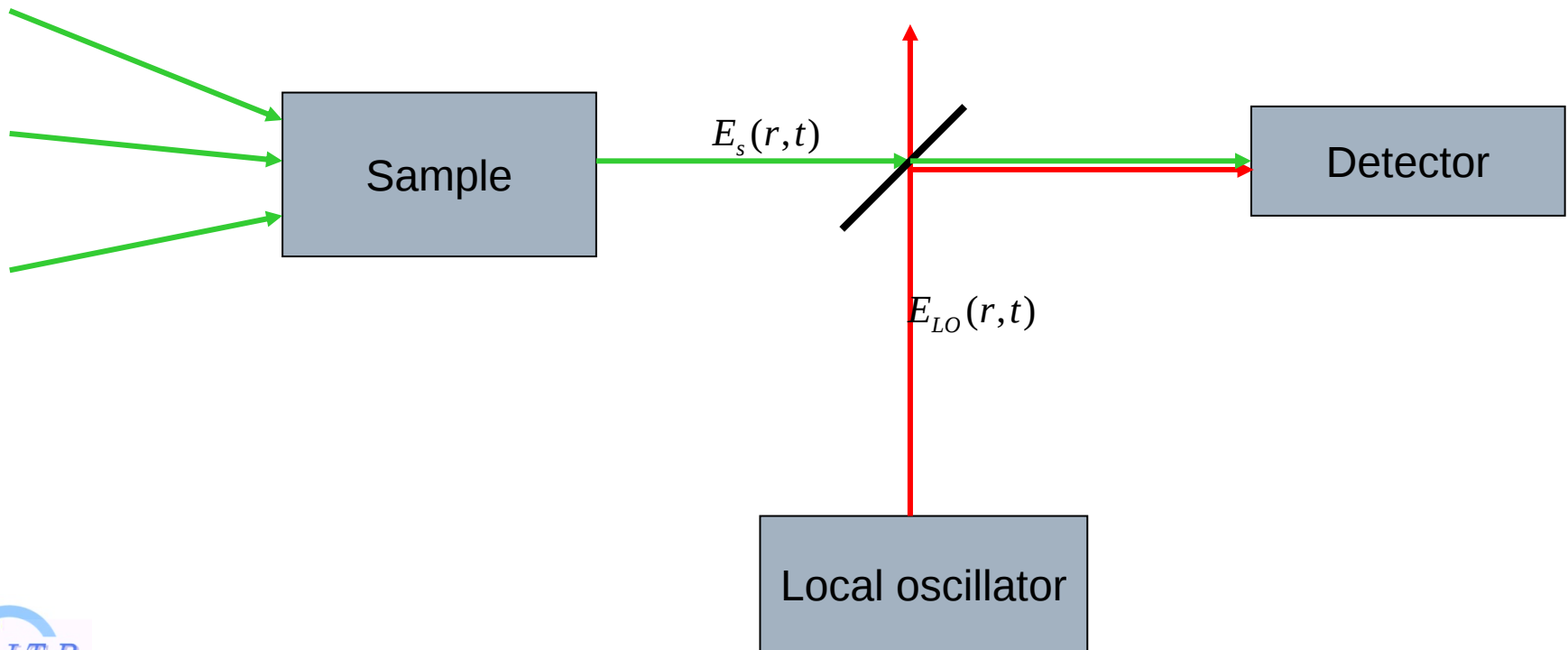


$$I_{k_s}(\Omega_3, \tau, \Omega_1)$$

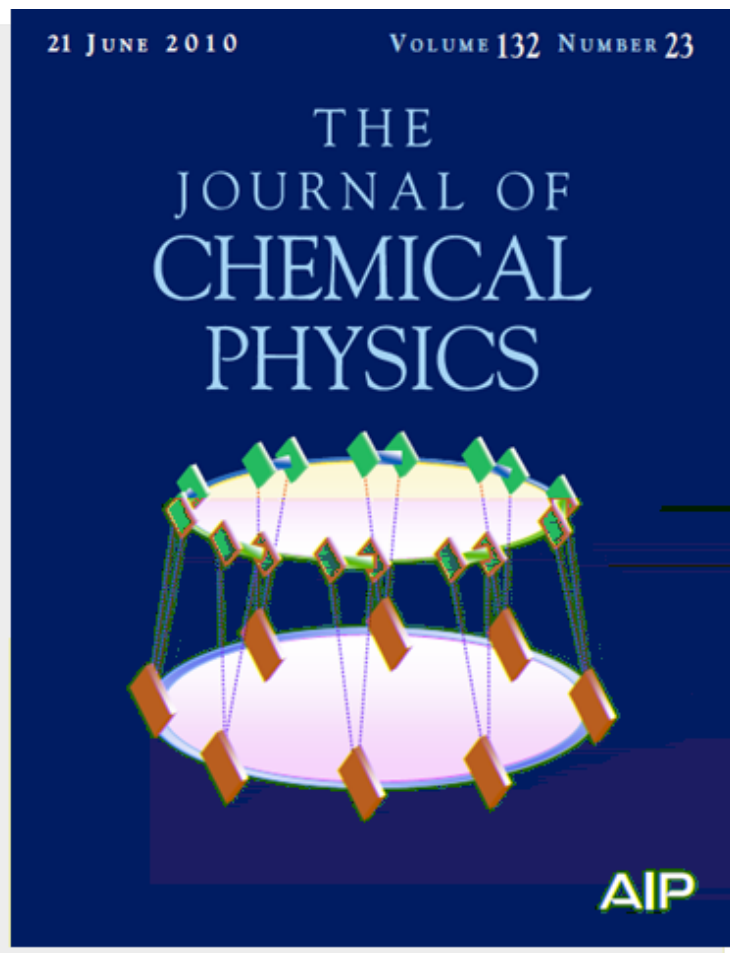


Heterodyne detected signal

□ Experimental setup for heterodyne detection



Dimerization-assisted energy transport in light-harvesting complexes



THE JOURNAL OF CHEMICAL PHYSICS 132, 234501 (2010)

Dimerization-assisted energy transport in light-harvesting complexes

S. Yang,¹ D. Z. Xu,¹ Z. Song,² and C. P. Sun^{1,a)}

¹Institute of Theoretical Physics, Chinese Academy of Sciences, Beijing 100190, China

²School of Physics, Beijing Normal University, Beijing 100875, China

20 February 2010; accepted 5 May 2010; published online 15 June 2010

(Received

he role of the dimer structure of light-harvesting complex II (LH2) in excitation transfer H2 [without a reaction center (RC)] to the LH1 (surrounding the RC) or from the LH2 LH2. The excited and unexcited states of a bacteriochlorophyll (BChl) are modeled by a In the framework of quantum open system theory, we represent the excitation transfer as akage of the LH2 system and then calculate the transfer efficiency and average transfer time. The results show that the dimerization of the B850 BChl ring can enhance the transfer efficiency and shorten the average transfer time. © 2010 American Institute of Physics. [doi:10.1063/1.3435213]

We study t from the L. to another quasispin. the total le

INTRODUCTION

To face the present and forthcoming global energy crisis, investigations on the basic energy science for prope has received great attention and experienced imre progress based on the fundamental physics.^{1,2} In ynthetic process, the structural elegance and chemical ficiency of the natural system based on pigment mol-in transferring the energy of sunlight have stimulated ose driven investigation,³⁻¹³ finding artificial analogs phyrin-based chromophores. These artificial systems te the natural process of photosynthesis² so that the higher efficiencies could be gained than that obtained conventional solid systems.⁴ It is because one of the attractive features of photosynthesis is that the light .can be captured and transferred to the reaction center within about 100 ps and with more than 95% efficiency. In most of the plants and bacteria, the primary processes of photosynthesis are almost the same. Light is absorbed by antenna pigments composed of many chromophores, then the electronic excitations are transferred to the reaction center where photochemical reactions take place to convert the excitation energy into chemical energy. Most recent experiments have been able to determine the time scales of various transfer processes via the ultrafast laser technology.¹⁴⁻¹⁸ These great es obviously offer us a chance to quantitatively

its induced decoherence could affect on the efficiency of the primary photosynthetic event. The present paper will similarly study the influences of spatial structure on the primary processes of photosynthesis for the light-harvesting complex.

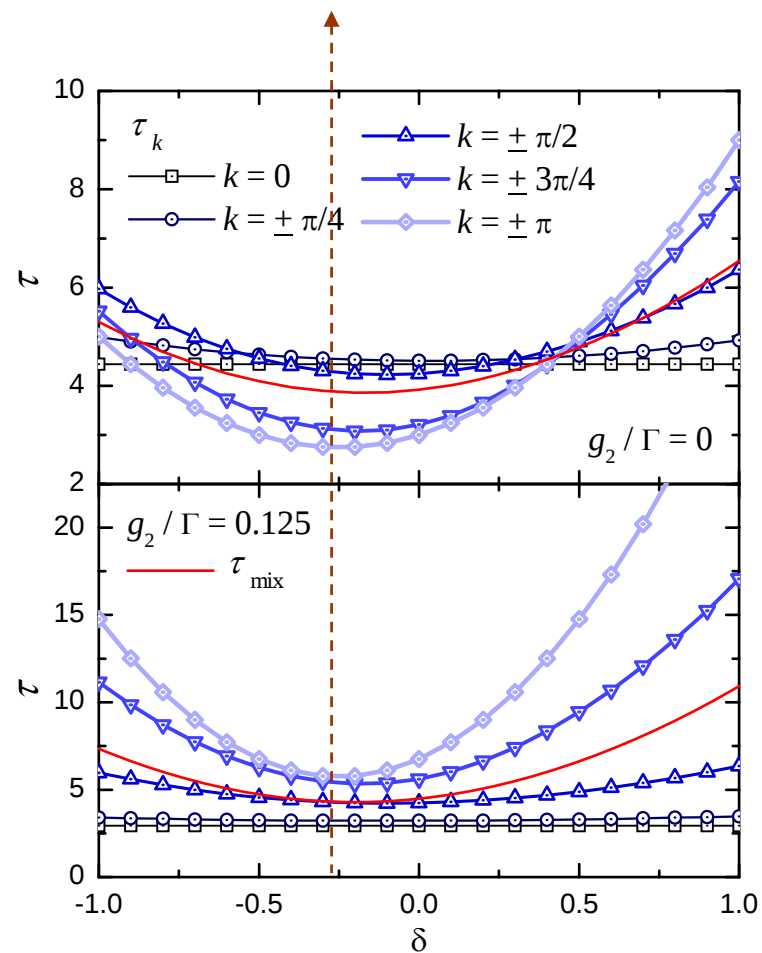
In the past, by making use of the x-ray crystallographic techniques, the structure of light-harvesting system has been elucidated.¹⁹ In the purple photosynthetic bacteria, there exist roughly two types of light-harvesting complexes, referred to as light-harvesting complex I (LH1) and LH2. In LH1, the RC is surrounded by a B875 bacteriochlorophyll (BChl) ring with maximum absorption peak at 875 nm. The LH2 complex, however, does not contain the RC but can transfer energy excitation to the RC indirectly through LH1. In the purple bacteria, LH2 is a ring-shaped aggregate built up by eight (or nine) minimal units, where each unit consists of an $\alpha\beta$ -heterodimer, three BChls, and one carotenoid. The structure of the LH2 complex is shown in Fig. 1. The BChl molecules are arranged in the LH2 complex in a double-helical structure. The top ring, including 16 (or 18) BChl molecules, is referred to as the B800 ring. The bottom ring, including 8 BChl molecules, is referred to as the B850 ring. The B800 ring is located at a distance of about 1.5 nm from the B850 ring. In each minimal unit, the carotenoid connects the B800 and B850 rings. The energy of the B800 ring is transferred to the B850 ring through the Förster mechanism,²⁰ while the electron is spatially transferred via the Marcus mechanism.²¹ Generally, it is independent of the global geometry configuration of the system.

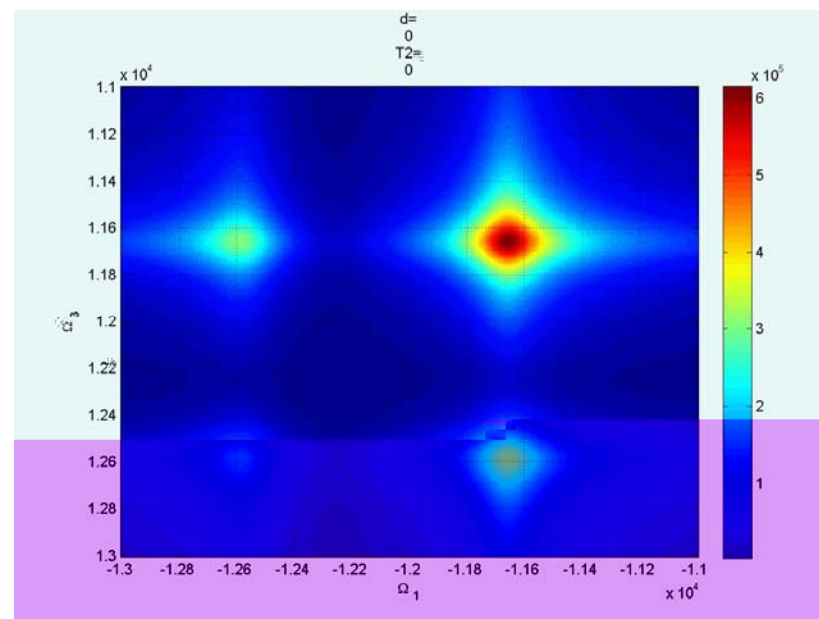
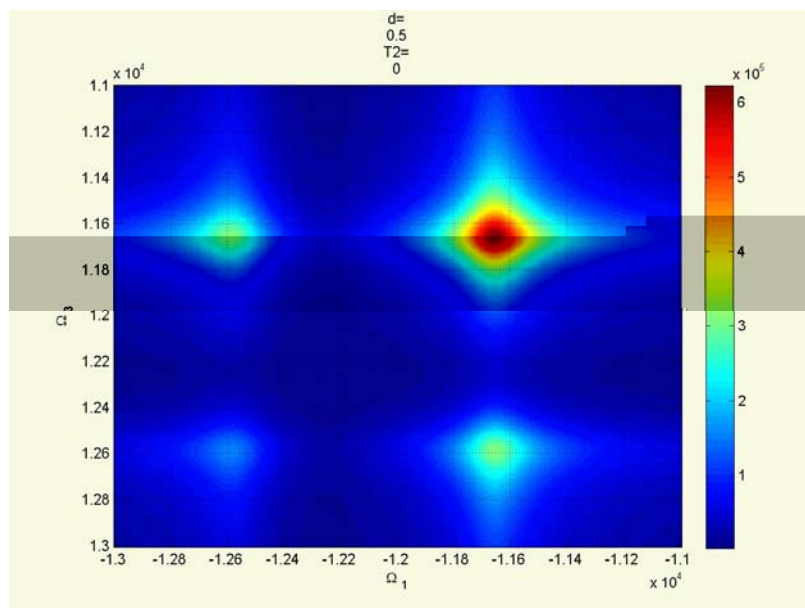
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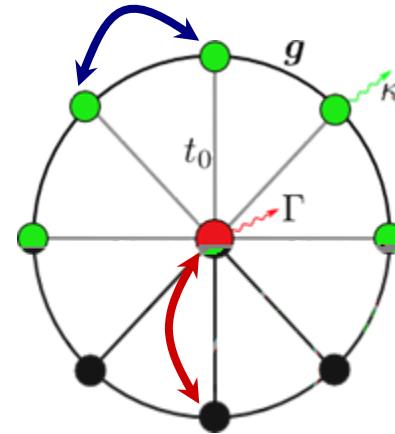
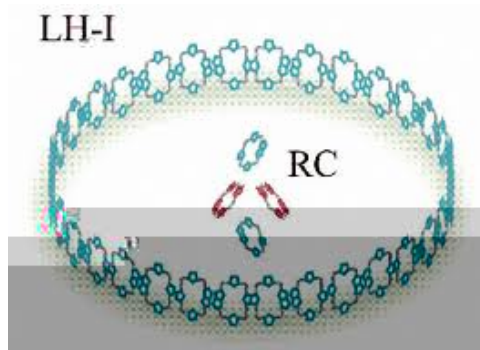
S. Yang, D. Z. Xu, Z. Song, CPS, JCP,132,234501(2010)







LH-I Model (I)



$$H_{DA} = \sum_{i=1}^N \left[\epsilon_i^\dagger \epsilon_i + \sum_{j=1}^N (t_{ij}^\dagger \epsilon_j^\dagger + t_{ij} \epsilon_j) \right]$$

$$+ \epsilon_A^\dagger A^\dagger A + \sum_{i=1}^N t_i (\epsilon_i^\dagger A^\dagger + \epsilon_i^\dagger A)$$

Γ : Usage rate of excitation

κ : Decay rate of excitation

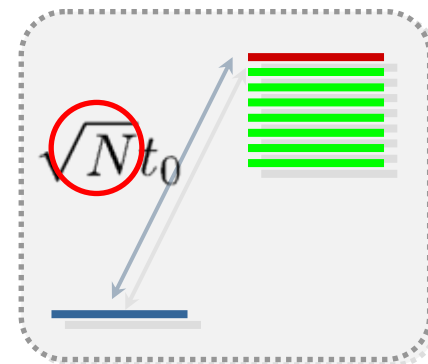
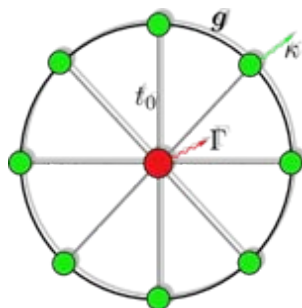
LH-I Model (II)

Fourier Transf: $\sum_k e^{ikr} \tilde{e}_k \tilde{e}_j = \frac{e^{ikr}}{\sqrt{N}}$

=0 mode: $|1_{k=0}\rangle = \frac{1}{\sqrt{N}} \sum_{i=1}^N |1_i\rangle$

Effective Hamiltonian:

$$H_{\text{eff}} = \sum_{i,j} t_{ij} \tilde{a}_i^\dagger \tilde{a}_j + \sum_i \left(\epsilon_0 + \frac{1}{\sqrt{N}} t_{i0} \right) \tilde{a}_i^\dagger \tilde{a}_i + \left(\epsilon_0 + \frac{1}{\sqrt{N}} t_{0i} \right) \tilde{a}_i^\dagger \tilde{a}_i$$



Transfer Properties

Transfer Efficiency:

$$\eta = \int_0^\infty 2\Gamma \underbrace{|v(t)|^2}_{\text{Population on Acceptor}} dt$$

Population on Acceptor

Average transfer time:

$$\tau = \frac{1}{\eta} \int_0^\infty 2\Gamma t |v(t)|^2 dt$$

Average output power:

$$\mathcal{P} = \eta/\tau$$

From Donor to Acceptor

Initial Excitation on Donor ring

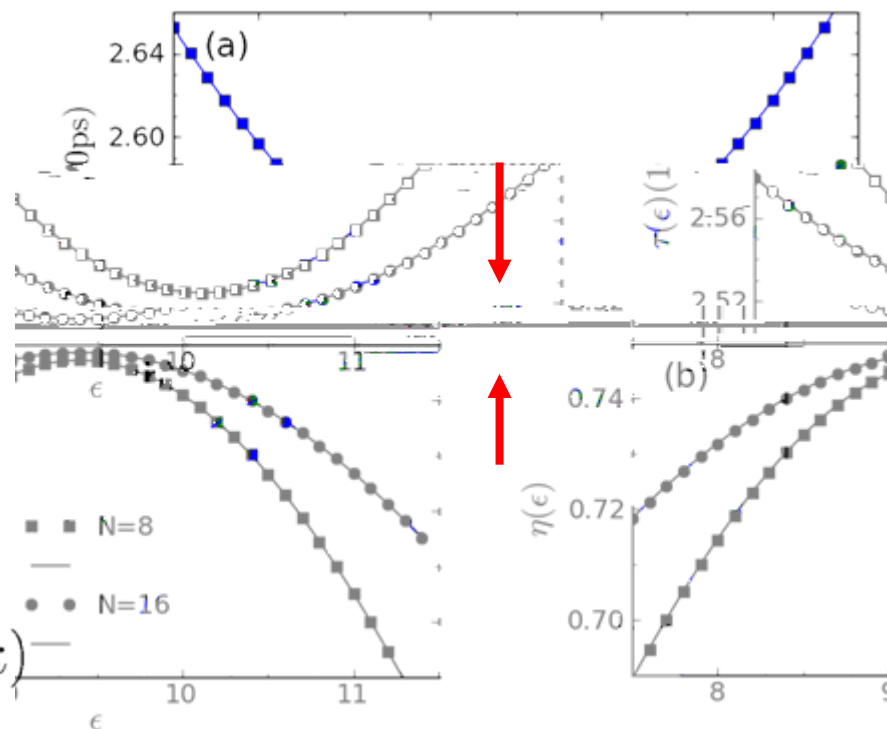
Optimal case:

$$\epsilon + 2g = \epsilon_A$$

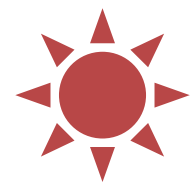
$\max \eta$

$\min \tau$

$$\eta_{\max} \leq \Gamma / (\Gamma + \kappa)$$



There should be light...



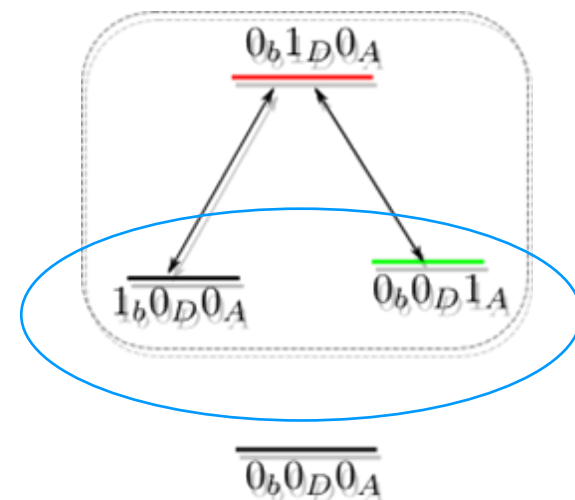
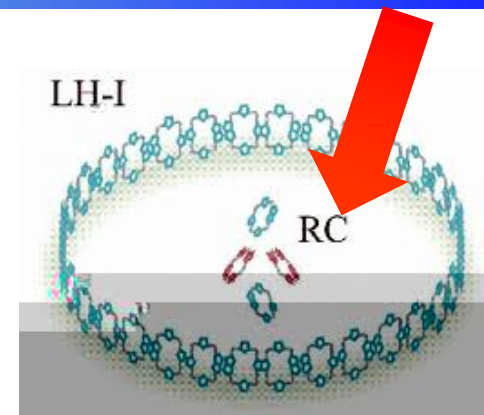
Visible light-wavelength $\sim 500\text{nm}$

LH-I radius $\sim 4-6\text{nm}$

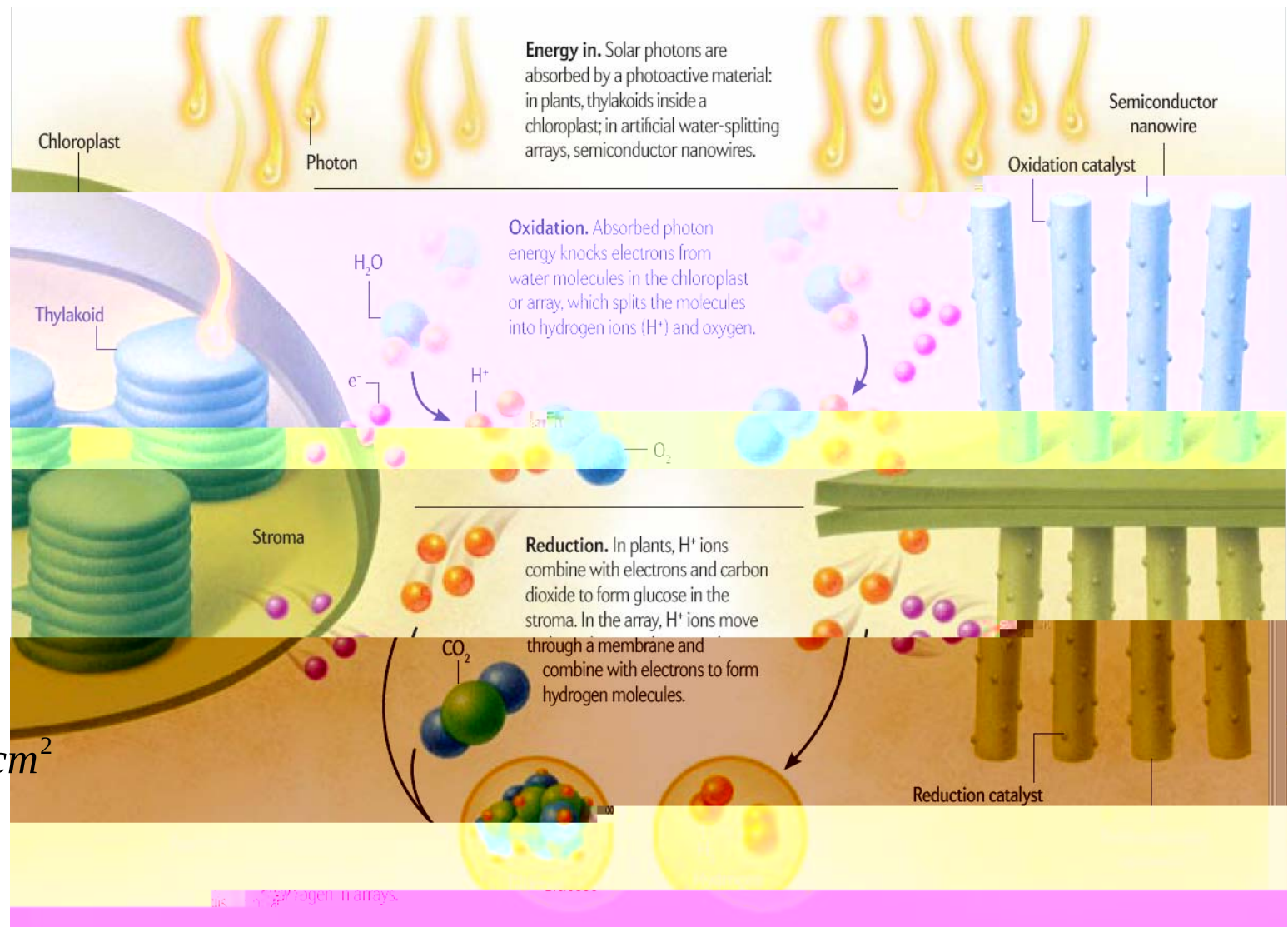
Photon-Donor interaction:

$$H_{pD} = J \sum_{i=1}^N (e_i^+ b + \text{h.c.})$$

$$H_{\text{eff}} = \omega_0 \tilde{e}_0^\dagger \tilde{e}_0 + \omega b^\dagger b + \omega_A A^\dagger A \\ + \sqrt{N} [(t_0 \tilde{e}_0^\dagger A + J \tilde{e}_0^\dagger b) + \text{h.c.}]$$



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10000 $\mu\text{sd} / \text{cm}^2$

Phys. Today **61**, July 28 (2008)



Grand challenges in basic energy sciences

Graham R. Fleming and Mark A. Ratner

Research focused in five related areas will allow unprecedented control over the microscopic world and could be the key to a sustainable future.

Thank You!